

INTEGRATION BY PARTS AND QUASI-INVARIANCE FOR THE HORIZONTAL WIENER MEASURE ON FOLIATED COMPACT MANIFOLDS

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ABSTRACT. We prove several sub-Riemannian versions of Driver's integration by parts formula which first appeared in [17]. Namely, our results are for the horizontal Wiener measure on a totally geodesic Riemannian foliation equipped with a sub-Riemannian structure. It is also shown that the horizontal Wiener measure is quasi-invariant under the action of flows generated by suitable tangent processes.

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1. INTRODUCTION

1.1. **Background.** In this paper we study quasi-invariance properties and related integration by parts formulas for the horizontal Wiener measure on a foliated Riemannian manifold equipped with a sub-Riemannian structure. These are most closely related to the well-known results by B. Driver [17] who established such properties for the Wiener measure on a path space over a compact Riemannian manifold. Quasi-invariance in such settings can be viewed as a curved version of the classical Cameron-Martin theorem for the Euclidean space. While the techniques developed for path spaces over Riemannian manifolds are not easily adaptable to the sub-Riemannian case we consider, we take advantage of the recent advances in this field. The geometric and stochastic analysis of sub-Riemannian structures on foliated manifolds has attracted a lot of attention in the past few years (see for instance [2, 7, 21, 29–32, 48]).

In particular, we make use of the tools such as Weitzenböck formulas for the sub-Laplacian extending results by J.-M. Bismut, B. Driver et al to foliated Riemannian manifolds. More precisely, the first progress in developing geometric techniques in the sub-Riemannian setting has been made in [5], where a version of Bochner’s formula for the sub-Laplacian was established and generalized curvature-dimension conditions have been studied. This Bochner-Weitzenböck formula was then used in [2] to develop a sub-Riemannian stochastic calculus. One of the difficulties in this case is that, a priori, there is no canonical connection on such manifolds such as the Levi-Civita connection in the Riemannian case. However, [2] introduces a one-parameter family of metric connections associated with Bochner’s formula proved in [5] and shows that the derivative of the sub-Riemannian heat semigroup can be expressed in terms of a damped stochastic parallel transport.

It should be noted that these connections do not preserve the geometry of the foliation in general. In particular, the corresponding parallel transport does not necessarily transform a horizontal vector into a horizontal vector, that is, these connections in general are not horizontal. As a consequence, establishing an integration by parts formulas for directional derivatives on the path space of the horizontal Brownian motion, similarly to Driver’s integration by parts formula in [17] for the Riemannian Brownian motion, is not straightforward. As a result, the integration by parts formula we prove in the current paper can not be simply deduced from the derivative formula

for the corresponding semigroup by applying the standard techniques of covariant stochastic analysis on manifolds as presented for instance in [23, Section 4], in particular [23, Theorems 4.1.1, 4.1.2]. A different approach to proving quasi-invariance in an infinite-dimensional sub-Riemannian setting has been used in [6].

Analysis on path and loop spaces has been developed over several decades, and we will not be able to refer to all the relevant publications, but we mention some which are closer to the subject and techniques of this paper. In particular, J.-M. Bismut's book [12] contains an integration by parts formula on a path space over a compact Riemannian manifold. His methods were based on the Malliavin calculus and Bismut's motivation was to deal with a hypoelliptic setting as described in [12, Section 5]. A breakthrough has been achieved by B. Driver [17], who established quasi-invariance properties of the Wiener measure over a compact Riemannian manifold, and as a consequence an integration by parts formula. This work has been simplified and extended by E. Hsu [34], and also approached by O. B. Enchev and D. W. Stroock in [25]. A review of these techniques can be found in [36]. In [37, 39] the noncompact case has been studied. Let us observe here that B. Driver in [17] and later E. Hsu in [34] have considered connections on a Riemannian manifold which are metric-compatible, but not necessarily torsion-free. This is very relevant in our setting of a foliated Riemannian manifold equipped with a sub-Riemannian structure, because on sub-Riemannian manifolds the natural connections are not torsion-free. A different approach to analysis on Riemannian path space can be found in [15], where tangent processes, Markovian connections, structure equations and other elements of what the authors call the renormalized differential geometry on the path space have been introduced.

1.2. Main results and organization of the paper. We now explain in more detail our main results without the technical details, and describe how the paper is organized. Section 4 studies quasi-invariance properties for the horizontal Wiener measure of a Riemannian foliation, and in Section 5 we prove integration by parts formulas. Although quasi-invariance properties and integration by parts formulas are intimately related and actually equivalent in many settings (see [9, 17, 18]), we use very different techniques and approaches in these two sections. To prove quasi-invariance, we develop a stochastic calculus of variations for the horizontal Brownian motion on a foliation in the spirit of [15, 17, 34], whereas to prove integration by parts formulas, we shall make use of Markovian techniques and martingale methods as presented for instance in [23, Section 4].

Let $(\mathbb{M}, g, \mathcal{F})$ be a smooth connected compact Riemannian manifold of dimension $n + m$ equipped with a totally geodesic and bundle-like foliation \mathcal{F} by m -dimensional leaves as described in Section 2. On such manifolds, one can define a horizontal Laplacian L according to [1, Section 2.2, Section 2.3]), allowing to define a horizontal Brownian motion as the diffusion on \mathbb{M}

with generator $\frac{1}{2}L$, as we describe in Section 4.1.1. The distribution of the horizontal Brownian motion is called the *horizontal Wiener measure*.

Recall that for a Riemannian manifold (\mathbb{M}, g) , for a given metric connection one can construct a development map $B \mapsto W$, where B is a Brownian motion in \mathbb{R}^{n+m} and W the Brownian motion on the manifold (\mathbb{M}, g) , see for instance [17, Theorem 3.4]. We construct development maps in the setting of a totally geodesic Riemannian foliation even though we do not have a Levi-Civita connection in this sub-Riemannian setting.

The foliation structure on \mathbb{M} induces a natural splitting of the tangent bundle into a vertical and horizontal subbundles \mathcal{V} and \mathcal{H} as described in Section 2.2. This allows us to define horizontal Brownian motion with respect to this structure. In Section 4.1.2 we show that there exist metric connections on \mathbb{M} which are compatible with the foliation \mathcal{F} in such a way that the above development map sends $(\beta, 0)$ to a horizontal Brownian motion of the foliation, where β is a Brownian motion in \mathbb{R}^n . In particular, the horizontal Brownian motion W on \mathbb{M} constructed in this way is a semimartingale on \mathbb{M} and it becomes possible to develop a horizontal stochastic calculus of variations. In this paper, the map $\beta \mapsto W$ is referred to as the *horizontal stochastic development map*. The main result of Section 3 is Theorem 1 that characterizes the variations of horizontal paths (i.e. paths transverse to the leaves of the foliation).

Before we can formulate our first main result, we need to describe some of the notation used. For details the reader is referred to Section 3. Let D be a metric connection on $(\mathbb{M}, g, \mathcal{F})$ adapted to the foliation structure as described by Assumption 1. An example of such a connection is the Bott connection introduced in Section 2.4.

The first observation is that the connection D allows us to define vector fields on the space $W_0^\infty(\mathbb{M})$ of smooth \mathbb{M} -valued paths on the interval $[0, 1]$ as follows. For $v \in W_0^\infty(\mathbb{R}^{n+m})$, the space of smooth \mathbb{R}^{n+m} -valued paths, we denote by \mathbf{D}_v the vector field on the space of smooth paths $[0, 1] \rightarrow \mathbb{M}$ defined by

$$\mathbf{D}_v(\gamma)_s = u_s(\gamma)v_s,$$

where u is the D -lift of γ to the orthonormal frame in the orthonormal frame bundle $\mathcal{O}(\mathbb{M})$. In addition, we can use the connection D to introduce the corresponding *development map* $\phi : W_0^\infty(\mathbb{R}^{n+m}) \rightarrow W_0^\infty(\mathbb{M})$ as defined in Definition 3.5. The inverse map $\phi^{-1} : W_0^\infty(\mathbb{M}) \rightarrow W_0^\infty(\mathbb{R}^{n+m})$ is referred to as an *anti-development map*. We also define horizontal development and anti-development maps in Definition 3.10.

In addition the connection D can be used to lift vector fields on \mathbb{M} to vector fields on $\mathcal{O}(\mathbb{M})$ consistent with the foliation structure as explained in Notation 3.3. We denote by A, V the fundamental vector fields on $\mathcal{O}(\mathbb{M})$ associated with this D -lift. For details we refer the reader to Notation 3.3. As motivation for the semimartingale version, we start with a theorem combining the results in the smooth setting.

Theorem 1 (Theorem 3.11, Theorem 3.15). Let D be a metric connection on $(\mathbb{M}, g, \mathcal{F})$ adapted to the foliation structure as described by Assumption 1. For a smooth path v on \mathbb{R}^{n+m} , we let $\{\zeta_t^v, t \in \mathbb{R}\}$ be the flow generated by the vector field \mathbf{D}_v on $W_0^\infty(\mathbb{M})$. Then for a smooth horizontal path γ on \mathbb{M}

$$\left. \frac{d}{dt} \right|_{t=0} \phi^{-1}(\zeta_t^v \gamma)_s \in \mathbb{R}^n, \quad s \in [0, 1]$$

if and only if the path

$$(1.1) \quad v(s) - \int_0^s T_{u_r}(A(d\omega_r^{\mathcal{H}}), Av(r)), \quad s \in [0, 1] \text{ takes values in } \mathbb{R}^n,$$

that is, it is horizontal. Here $\omega^{\mathcal{H}}$ is the horizontal anti-development of the horizontal path γ , and T is the torsion of the Bott connection.

Moreover, if (1.1) is satisfied, then

$$\left. \frac{d}{dt} \right|_{t=0} \phi^{-1}(\zeta_t^v \gamma)_s = p_v(\omega^{\mathcal{H}})_s,$$

where

$$p_v(\omega^{\mathcal{H}})_s = v(s) - \int_0^s T_{u_r}^D(A(d\omega_r^{\mathcal{H}}), Av(r) + Vv(r)) - \int_0^s \left(\int_0^r \Omega_{u_\tau}^D(A(d\omega_\tau^{\mathcal{H}}), Av(\tau) + Vv(\tau)) \right) d\omega_r^{\mathcal{H}}.$$

Here T^D is the torsion form of the connection D and Ω^D its curvature form.

If (1.1) is satisfied, we will say that the path v is *tangent* to the horizontal path γ . We stress that in (1.1) we use the torsion of the Bott connection, not the torsion of the connection D . Thus Theorem 1 shows that the notion of a tangent path is independent of the connection D , as long as it satisfies Assumption 1.

Given a horizontal path, it is easy to construct tangent paths to this path. Indeed, we show in Lemma 3.17 that if $\omega^{\mathcal{H}}$ is a smooth path in \mathbb{R}^n then for every smooth path h in \mathbb{R}^n

$$(1.2) \quad \tau_h(\omega^{\mathcal{H}})_s := h(s) + \int_0^s T_{u_r}(A(d\omega_r^{\mathcal{H}}), Ah(r))$$

is a tangent path to $\phi(\omega^{\mathcal{H}})$, where u denotes the lift of $\phi(\omega^{\mathcal{H}})$.

In Section 4.3 we use Malliavin's *principe de transfert* ansatz (see [41, Part IV Chapter VIII]) to extend the definitions for p_v and τ_h to semimartingale paths by replacing integration against smooth paths by Stratonovich stochastic integrals with respect to semimartingales. More precisely, we work on the probability space $(W_0(\mathbb{R}^n), \mathcal{B}, \mu_{\mathcal{H}})$, where \mathcal{B} is the Borel σ -algebra on the path space $W_0(\mathbb{R}^n)$ of continuous paths $[0, 1] \rightarrow \mathbb{R}^n$ starting at 0, $\mu_{\mathcal{H}}$ is the Wiener measure. The measure μ_W can be also described as the distribution of the horizontal Brownian motion on \mathbb{M} .

Given a deterministic Cameron-Martin path $h : [0, 1] \rightarrow \mathbb{R}^n$, one can then consider the \mathbb{R}^{n+m} -valued semimartingale

$$\tau_h(\omega^{\mathcal{H}})_s := h(s) + \int_0^s T_{u_r}(A \circ d\omega_r^{\mathcal{H}}, Ah(r)),$$

where $\omega^{\mathcal{H}}$ is the coordinate process and $\circ d\omega^{\mathcal{H}}$ denotes the Stratonovich integral. Note that τ_h is defined $\mu_{\mathcal{H}}$ -a.s. One can then think of $\tau_h(\omega^{\mathcal{H}})$ as a tangent process to the horizontal Brownian motion of the foliation. We will view $\tau_h : W_0(\mathbb{R}^n) \rightarrow W_0(\mathbb{R}^{n+m})$ as an adapted process with respect to the natural filtration $\{\mathcal{B}_s, 0 \leq s \leq 1\}$ generated by the horizontal Brownian motion in \mathbb{R}^{n+m} . Notice that τ_h is really an equivalence class of processes with two processes being equivalent if they are equal $\mu_{\mathcal{H}}$ -a.s. similarly to [34, p. 425]. Thus when we say that a map is defined $\mu_{\mathcal{H}}$ -a.s. we mean that we are actually working with equivalence classes of maps. It will be an important part of our results that the flows and the compositions we consider preserve the equivalence classes we are working with, but for simplicity of the presentation, those considerations will remain in the background in our discussions similarly to [34]. This aspect is discussed more thoroughly in [17].

Similarly, given an \mathbb{R}^{n+m} -valued semimartingale v , one can define the semimartingale

$$p_v(\omega^{\mathcal{H}})_s = v(s) - \int_0^s T_{u_r}^D(A \circ d\omega_r^{\mathcal{H}}, Av(r) + Vv(r)) - \int_0^s \left(\int_0^r \Omega_{u_\tau}^D(A \circ d\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) \circ d\omega_r^{\mathcal{H}}.$$

The main results of Section 4.3 include Theorem 4.21 and Theorem 4.23 which can be summarized as follows. Here we use the notion of stochastic horizontal development and anti-development $\phi_{\mathcal{H}}$ and $\phi_{\mathcal{H}}^{-1}$ as defined in Definition 4.9.

Theorem 2 (Theorem 4.21, Theorem 4.23). There exists a family of semimartingales $\{\nu_t^h, t \in \mathbb{R}\}$ such that for each fixed t the random variable $\nu_t^h : W_0(\mathbb{R}^n) \rightarrow W_0(\mathbb{R}^n)$ can be regarded as a $\mu_{\mathcal{H}}$ -a. s. defined map from the path space to itself. In particular, $s \rightarrow \nu_t^h(s)$ is a \mathbb{R}^n -valued semimartingale over the probability space $W_0(\mathbb{R}^n)$. In addition ν_t^h has the group property. Thus we can regard $\{\nu_t^h, t \in \mathbb{R}\}$ as the flow on $W_0(\mathbb{R}^n)$ generated by $p_{\tau_h} : W_0(\mathbb{R}^n) \rightarrow W_0(\mathbb{R}^n)$ which is defined $\mu_{\mathcal{H}}$ -a.s.

Moreover, the measure $\mu_{\mathcal{H}}$ is quasi-invariant under this flow, that is, the law μ_W of the horizontal Brownian motion on \mathbb{M} is quasi-invariant under the μ_W -a.s. defined flow $\zeta_t^h = \phi_{\mathcal{H}} \circ \nu_t^h \circ \phi_{\mathcal{H}}^{-1} : W_{x_0}(\mathbb{M}) \rightarrow W_{x_0}(\mathbb{M})$, $t \in \mathbb{R}$. Here $\phi_{\mathcal{H}}$ and $\phi_{\mathcal{H}}^{-1}$ are horizontal stochastic development and anti-development map correspondingly.

It should be noted that our argument follows relatively closely the one by B. Driver in [17] and later by E. Hsu in [34] (see also [14, 15, 25]) and therefore going from Theorem 1 to Theorem 2 is quite routine. In Section

4.3.6 we illustrate our results in the case of Riemannian submersions and explicitly compute the flow ζ_t^h associated to the Bott connection in some examples.

The goal of the second part of the paper is to establish several types of integration by parts formulas for the horizontal Brownian motion. In Section 5.1.1, we survey known geometric and stochastic results and introduce the notation and conventions used throughout Section 5. Most of this material is based on [7] for the geometric part and [2] for the stochastic part. The most relevant result that will be used later is the Weitzenböck formula given in Theorem 5.4. It asserts that for every $f \in C^\infty(\mathbb{M})$, $x \in \mathbb{M}$ and every $\varepsilon > 0$

$$(1.3) \quad dLf(x) = \square_\varepsilon df(x),$$

where \square_ε is a one-parameter family of sub-Laplacians on one-forms indexed by a parameter $\varepsilon > 0$. These sub-Laplacians on one-forms are constructed from a family of metric connections ∇^ε introduced in [2] whose adjoint connections $\widehat{\nabla}^\varepsilon$ in the sense of B. Driver in [17] are also metric. These connections satisfy Assumption 1, so that the results of Section 4 are applicable. Even though Section 5.1.1 introduces mostly preliminaries, we present a number of new results there such as Lemma 5.6.

In Section 5.2, we prove integration by parts formulas for the horizontal Wiener measure with the main result being Theorem 5.20 which includes the following result. Suppose F is a cylinder function, v is a tangent process on $T_x\mathbb{M}$ as defined in Definition 5.16, then we have

$$(1.4) \quad \mathbb{E}_x(\mathbf{D}_v F) = \mathbb{E}_x \left(F \int_0^1 \left\langle v'_\mathcal{H}(s) + \frac{1}{2} \parallel_{0,s}^{-1} \mathfrak{Ric}_\mathcal{H} \parallel_{0,s}, dB_s \right\rangle_\mathcal{H} \right),$$

where x is the starting point of the horizontal Brownian motion, $\mathbf{D}_v F$ is the directional derivative of F in the direction of v , $\parallel_{0,s}$ is the stochastic parallel transport for the Bott connection, and $\mathfrak{Ric}_\mathcal{H}$ is the horizontal Ricci curvature of the Bott connection. The Bott connection as defined in Section 2.4 corresponds to the adjoint connection $\widehat{\nabla}^\varepsilon$ as $\varepsilon \rightarrow \infty$. In the integration by parts formula (1.4), the tangent process v is a $T_x\mathbb{M}$ -valued process such that its horizontal part $v_\mathcal{H}$ is absolutely continuous and satisfies $\mathbb{E} \left(\int_0^1 \|v'_\mathcal{H}(s)\|_{T_x\mathbb{M}}^2 ds \right) < \infty$ and its vertical part is given by

$$(1.5) \quad v_\mathcal{V}(s) = \int_0^s \parallel_{0,r}^{-1} T(\parallel_{0,r} \circ dB_r, \parallel_{0,r} v_\mathcal{H}(r)),$$

where T is the torsion tensor of the Bott connection. Observe that (1.4) looks similar to the integration by parts formulas by J.-M. Bismut and B. Driver. This is not too surprising if one thinks about the special case when the foliation comes from a Riemannian submersion with totally geodesic fibers. We consider this case in Section 5.3.1, and we prove that then that the integration by parts formula in Theorem (5.20) is actually a horizontal lift of Driver's formula from the base space of the fibers to \mathbb{M} . However, in

general foliations do not come from submersions (see for instance [26] for necessary and sufficient conditions) and one therefore needs to develop an intrinsic horizontal stochastic calculus on \mathbb{M} to prove (1.4). Developing such a calculus is one of the main accomplishments of the current paper.

The proof of Theorem 5.20 proceeds in several steps. As in [2], the Weitzenböck formula (1.3) yields a stochastic representation for the derivative of the semigroup of the horizontal Brownian motion in terms of a damped stochastic parallel transport associated to the connection ∇^ε (see Lemma 5.21). By using techniques of [4], Lemma 5.21 implies an integration by parts formula for the damped Malliavin derivative as stated in Theorem 5.19. The final step is to prove Theorem 5.20 from Theorem 5.19. The main difficulty is that the connection ∇^ε is in general not horizontal. However, it turns out that the adjoint connection $\widehat{\nabla}^\varepsilon$ is not only metric but also horizontal. As a consequence, one can use the orthogonal invariance of the horizontal Brownian motion (Lemma 5.26) to filter out the redundant noise which is given by the torsion tensor of ∇^ε . It is remarkable that the integration by parts formula for the directional derivatives in Theorem (5.20) is actually independent of the choice of a particular connection and therefore is independent of ε in the one-parameter family of connections used to define the damped Malliavin derivative. While integration by parts formulas for the damped Malliavin derivative may be used to prove gradient bounds for the heat semigroup (as in [2]) and log-Sobolev inequalities on the path space (as in [4]), we prove that the integration by parts formula (1.4) comes from the quasi-invariance property of the horizontal Wiener measure proved in Section 4.3.

Remark 1.1. In the current paper, we restrict consideration to the case of compact manifolds mainly for the sake of conciseness. It is reasonable to conjecture that as in [39], our results may be extended to complete manifolds.

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2. GEOMETRIC PRELIMINARIES: RIEMANNIAN FOLIATIONS

2.1. Riemannian foliations. We start by recalling the notion of a foliation. Let \mathbb{M} be a smooth connected manifold of dimension $n + m$. Then a foliation of dimension m on \mathbb{M} is usually described as a collection \mathcal{F} of disjoint connected non-empty immersed m -dimensional submanifolds of \mathbb{M} (called the *leaves* of the foliation), whose union is \mathbb{M} , and such that in a neighborhood of each point in \mathbb{M} there exists a chart for \mathcal{F} as follows.

Before we define such Riemannian foliations, let us introduce some standard notation.

Notation 2.1. Suppose (\mathbb{M}, g) is a Riemannian manifold. By $T\mathbb{M}$ we denote the *tangent bundle* and by $T^*\mathbb{M}$ the *cotangent bundle*, and by $T_x\mathbb{M}$ ($T_x^*\mathbb{M}$) the *tangent (cotangent) space* at $x \in \mathbb{M}$. The inner product on $T\mathbb{M}$ induced by the metric g will be denoted by $g(\cdot, \cdot)$. If \mathcal{U} is a subbundle of the tangent bundle $T\mathbb{M}$, the restriction of g to \mathcal{U} will be denoted by $g_{\mathcal{U}}(\cdot, \cdot)$.

As always, for any $x \in \mathbb{M}$ we denote by $g(\cdot, \cdot)_x$ (or $\langle \cdot, \cdot \rangle_x$), $g_{\mathcal{U}}(\cdot, \cdot)_x$ (or $\langle \cdot, \cdot \rangle_{\mathcal{U}_x}$) (or $\langle \cdot, \cdot \rangle_{\mathcal{U}_x}$) the inner product on the fibers $T_x\mathbb{M}$ and \mathcal{U}_x correspondingly. The space of *smooth functions* on \mathbb{M} will be denoted by $C^\infty(\mathbb{M})$. The space of *smooth sections* of a vector bundle \mathcal{E} over \mathbb{M} will be denoted $\Gamma^\infty(\mathcal{E})$.

Definition 2.2. Let \mathbb{M} be a smooth connected $n+m$ -dimensional manifold. An m -dimensional foliation \mathcal{F} on \mathbb{M} is defined by a (maximal) collection of pairs $\{(U_\alpha, \pi_\alpha), \alpha \in I\}$ of open subsets U_α of \mathbb{M} and submersions $\pi_\alpha : U_\alpha \rightarrow U_\alpha^0$ onto open subsets of \mathbb{R}^n satisfying

- $\bigcup_{\alpha \in I} U_\alpha = \mathbb{M}$;
- If $U_\alpha \cap U_\beta \neq \emptyset$, there exists a local diffeomorphism $\Psi_{\alpha\beta}$ of \mathbb{R}^n such that $\pi_\alpha = \Psi_{\alpha\beta} \pi_\beta$ on $U_\alpha \cap U_\beta$.

In addition, we assume that the foliation \mathcal{F} on \mathbb{M} is a Riemannian foliation with a bundle-like metric g and totally geodesic m -dimensional leaves. Informally a bundle-like metric is similar to a product metric locally, and the notion has been introduced in [45]. We refer to [1, 43, 45, 49] for details about the geometry of Riemannian foliations, but for convenience of the reader we recall some basic definitions.

The maps π_α are called *disintegrating maps* of \mathcal{F} . The connected components of the sets $\pi_\alpha^{-1}(c)$, $c \in \mathbb{R}^n$, are called the *plaques* of the foliation. For each $p \in U_\alpha$, we define $\mathcal{V}_p := \text{Ker}((\pi_\alpha)_*|_p)$. The subbundle \mathcal{V} of $T\mathbb{M}$ with fibers \mathcal{V}_p is referred to as the *vertical distribution*. These are the vectors tangent to the leaves, the maximal integral sub-manifolds of \mathcal{V} .

Definition 2.3. Let \mathbb{M} be a smooth connected $n+m$ -dimensional Riemannian manifold. An m -dimensional foliation \mathcal{F} on \mathbb{M} is said to be *Riemannian with a bundle-like metric* if the disintegrating maps π_α are Riemannian submersions onto U_α^0 with its given Riemannian structure. If moreover the leaves are totally geodesic sub-manifolds of \mathbb{M} , then we say that the Riemannian foliation is *totally geodesic with a bundle-like metric*.

2.2. Horizontal and vertical subbundles of $T\mathbb{M}$ and forms. The subbundle \mathcal{H} which is normal to the vertical subbundle \mathcal{V} is referred to as the set of *horizontal directions*. Though this assumption is not strictly necessary in many parts of the paper, to simplify the presentation we always assume that \mathcal{H} is bracket-generating, that is, the Lie algebra of vector fields generated by global C^∞ -sections of \mathcal{H} has the full rank at each point in \mathbb{M} . Using Notation 2.1, we denote the restrictions of the metric g to \mathcal{H} and \mathcal{V} by $g_{\mathcal{H}}(\cdot, \cdot)$ and $g_{\mathcal{V}}(\cdot, \cdot)$ respectively.

We say that a one-form is *horizontal (resp. vertical)* if it vanishes on the vertical bundle \mathcal{V} (resp. on the horizontal bundle \mathcal{H}). Then the splitting of

the tangent space

$$T_x\mathbb{M} = \mathcal{H}_x \oplus \mathcal{V}_x$$

induces a splitting of the cotangent space

$$T_x^*\mathbb{M} = \mathcal{H}_x^* \oplus \mathcal{V}_x^*.$$

The subbundle \mathcal{H}^* of the cotangent bundle will be referred to as the co-horizontal bundle. Similarly, \mathcal{V}^* will be referred to as the covertical bundle.

2.3. Examples.

Example 2.1 (Riemannian submersions, Hopf fibrations). Let (\mathbb{M}, g) and (\mathbb{B}, j) be two smooth and connected Riemannian manifolds. A smooth surjective map $\pi : \mathbb{M} \rightarrow \mathbb{B}$ is called a *Riemannian submersion* if for every $x \in \mathbb{M}$ the differential $T_x\pi : T_x\mathbb{M} \rightarrow T_{\pi(x)}\mathbb{B}$ is an orthogonal projection, i.e. the map $T_x\pi(T_x\pi)^* : T_{\pi(x)}\mathbb{B} \rightarrow T_{\pi(x)}\mathbb{B}$ is the identity map. The foliation given by the fibers of a Riemannian submersion is then bundle-like (see [1, Section 2.3]). We refer to [11, Chapter 9, Section F, pp. 249-252] for Riemannian submersions with totally geodesic fibers.

The generalized Hopf fibrations (e.g. [11, Chapter 9, Section H], [44, Section 1.4.6]) offer a wide range of examples of Riemannian submersions whose fibers are totally geodesic. Let G be a Lie group, and H, K be two compact subgroups of G with $K \subset H$. Then, we have a natural fibration given by the coset map

$$\begin{aligned} \pi : G/K &\longrightarrow G/H \\ \alpha K &\longmapsto \alpha H, \end{aligned}$$

where the fiber is H/K . From [10], there exist G -invariant metrics on respectively G/K and G/H that make π a Riemannian submersion with totally geodesic fibers isometric to H/K . For instance with $G = SU(n+1)$, $H = S(U(1)U(n)) \simeq U(n)$ and $K = SU(n)$, one obtains the usual Hopf fibration $\pi : S^{2n+1} \rightarrow \mathbb{C}P^n$, see [11, Chapter 9, Section H, Example 9.81]. For $n = 1$, this reduces to the Hopf fibration $\pi : SU(2) \simeq S^3 \rightarrow \mathbb{C}P^1 \simeq S^2$.

Example 2.2 (K -contact manifolds). Another important example of a Riemannian foliation is obtained in the context of contact manifolds. Let (\mathbb{M}, θ) be a $2n+1$ -dimensional smooth contact manifold, where θ is a contact form. Then there is a unique smooth vector field Z on \mathbb{M} , called the *Reeb vector field*, satisfying

$$\theta(Z) = 1, \quad \mathcal{L}_Z(\theta) = 0,$$

where \mathcal{L}_Z denotes the Lie derivative with respect to Z . The Reeb vector field induces a foliation on \mathbb{M} , the Reeb foliation, whose leaves are the orbits of the vector field Z . It is known (see [46, 47]), that it is always possible

to find a Riemannian metric g and a $(1, 1)$ -tensor field J on \mathbb{M} so that for every vector fields X, Y

$$g(X, Z) = \theta(X), \quad J^2(X) = -X + \theta(X)Z, \quad g(X, JY) = (d\theta)(X, Y).$$

The triple (\mathbb{M}, θ, g) is called a *contact Riemannian manifold*. We see then that the Reeb foliation is totally geodesic with a bundle-like metric if and only if the Reeb vector field Z is a Killing field, that is,

$$\mathcal{L}_Z g = 0,$$

as is stated in [13, Proposition 6.4.8]. In this case, (\mathbb{M}, θ, g) is called a *K-contact Riemannian manifold*. Observe that the horizontal distribution \mathcal{H} is then the kernel of θ and that \mathcal{H} is bracket generating because θ is a contact form. We refer to [8, 47] for further details on this class of examples.

2.4. Bott connection. If we view (\mathbb{M}, g) as a Riemannian manifold, the Levi-Civita connection ∇^R is a natural choice for stochastic analysis on \mathbb{M} . But this connection is not adapted to the study of foliations because the horizontal and vertical bundles may not be parallel with respect to ∇^R . We will rather make use of the *Bott connection* on \mathbb{M} which is defined as follows.

$$\nabla_X Y = \begin{cases} \pi_{\mathcal{H}}(\nabla_X^R Y), & X, Y \in \Gamma^\infty(\mathcal{H}), \\ \pi_{\mathcal{H}}([X, Y]), & X \in \Gamma^\infty(\mathcal{V}), Y \in \Gamma^\infty(\mathcal{H}), \\ \pi_{\mathcal{V}}([X, Y]), & X \in \Gamma^\infty(\mathcal{H}), Y \in \Gamma^\infty(\mathcal{V}), \\ \pi_{\mathcal{V}}(\nabla_X^R Y), & X, Y \in \Gamma^\infty(\mathcal{V}), \end{cases}$$

where $\pi_{\mathcal{H}}$ (resp. $\pi_{\mathcal{V}}$) is the projection on \mathcal{H} (resp. \mathcal{V}). One can check that since the foliation is bundle-like and totally geodesic the Bott connection is metric-compatible, that is, $\nabla g = 0$, though unlike the Levi-Civita connection it is not torsion-free. The following properties of the Bott connection are standard but require tedious computations. We refer to [49, Chapter 5] for some of these, and to [42] for the details of the statements below and also point out that the Bott connection is a special case of a general class of connections introduced by R. Hladky in [33, Lemma 2.13].

Let T be the torsion of the Bott connection ∇ . Observe that for $X, Y \in \Gamma^\infty(\mathcal{H})$

$$\begin{aligned} T(X, Y) &= \nabla_X Y - \nabla_Y X - [X, Y] \\ &= \pi_{\mathcal{H}}(\nabla_X^R Y - \nabla_Y^R X) - [X, Y] \\ &= \pi_{\mathcal{H}}([X, Y]) - [X, Y] \\ &= -\pi_{\mathcal{V}}([X, Y]). \end{aligned}$$

Similarly one can check that the Bott connection satisfies the following properties that we record here for later use

$$\begin{aligned}
(2.1) \quad & \nabla_X Y \in \Gamma^\infty(\mathcal{H}) \text{ for any } X, Y \in \Gamma^\infty(\mathcal{H}), \\
& \nabla_X Y \in \Gamma^\infty(\mathcal{V}) \text{ for any } X, Y \in \Gamma^\infty(\mathcal{V}), \\
& T(X, Y) \in \Gamma^\infty(\mathcal{V}) \text{ for any } X, Y \in \Gamma^\infty(\mathcal{H}), \\
& T(U, V) = 0 \text{ for any } U, V \in \Gamma^\infty(\mathcal{V}), \\
& T(X, U) = 0 \text{ for any } X \in \Gamma^\infty(\mathcal{H}), U \in \Gamma^\infty(\mathcal{V}).
\end{aligned}$$

Example 2.3 (Example 2.1 revisited). Let $\pi : (\mathbb{M}, g) \rightarrow (\mathbb{B}, j)$ be a Riemannian submersion with totally geodesic leaves. A vector field $X \in \Gamma^\infty(T\mathbb{M})$ is said to be *projectable* if there exists a smooth vector field \bar{X} on \mathbb{B} such that for every $x \in \mathbb{M}$, $T_x\pi(X(x)) = \bar{X}(\pi(x))$. In that case, we say that X and \bar{X} are π -related. A vector field X on \mathbb{M} is called *basic* if it is projectable and horizontal. If \bar{X} is a smooth vector field on \mathbb{B} , then there exists a unique basic vector field X on \mathbb{M} which is π -related to \bar{X} . This vector is called the *lift* of \bar{X} . The Bott connection is then a lift of the Levi-Civita connection of (\mathbb{B}, j) in the following sense:

$$(2.2) \quad \nabla_{\bar{X}}^{\mathbb{B}} \bar{Y} = \overline{\nabla_X Y}, \quad \bar{X}, \bar{Y} \in \Gamma^\infty(\mathbb{B}),$$

where $\nabla^{\mathbb{B}}$ is the Levi-Civita connection on \mathbb{B} .

Example 2.4 (Example 2.2 revisited). Let (\mathbb{M}, θ, g) be a K-contact Riemannian manifold. The Bott connection coincides with Tanno's connection that was introduced in [47] and which is the unique connection that satisfies the following properties.

- (1) $\nabla\theta = 0$;
- (2) $\nabla Z = 0$;
- (3) $\nabla g = 0$;
- (4) $T(X, Y) = d\theta(X, Y)Z$ for any $X, Y \in \Gamma^\infty(\mathcal{H})$;
- (5) $T(Z, X) = 0$ for any vector field $X \in \Gamma^\infty(\mathcal{H})$.

2.5. Orthonormal frame bundle. We will use standard notation for orthonormal frame bundles. Suppose \mathbb{M} is a compact Riemannian manifold of dimension d . Note that in the setting of Riemannian foliations we have $d = n + m$. Recall that a frame at $x \in \mathbb{M}$ can be described as a linear isomorphism $u : \mathbb{R}^d \rightarrow T_x\mathbb{M}$ such that for the standard basis $\{e_i\}_{i=1}^d$ of \mathbb{R}^d the collection $\{u(e_i)\}_{i=1}^d$ is a basis (frame) for $T_x\mathbb{M}$. The collection of all such frames $\mathcal{F}(\mathbb{M}) := \bigcup_{x \in \mathbb{M}} \mathcal{F}(\mathbb{M})_x$ is called the *frame bundle* with the group $\text{GL}(\mathbb{R}, d)$ acting on the bundle. If \mathbb{M} is in addition Riemannian, we can restrict ourselves to consideration of Euclidean isometries $u : (\mathbb{R}^d, \langle \cdot, \cdot \rangle) \rightarrow (T_x\mathbb{M}, g)$ with the group $\text{O}(\mathbb{R}, d)$ acting on the bundle. The *orthonormal frame bundle* will be denoted by $\mathcal{O}(\mathbb{M})$.

Suppose that D is a connection on \mathbb{M} , then D induces a decomposition of each tangent space $T_u\mathcal{O}(\mathbb{M})$ into the direct sum of a horizontal subspace

and a vertical subspace as described in [38, Section 2.1]. Using such decomposition, one can then lift smooth maps on \mathbb{M} into smooth horizontal paths on $\mathcal{O}(\mathbb{M})$, see [34, p. 421]. Such a lift is usually called the horizontal lift to $\mathcal{O}(\mathbb{M})$. However, to avoid the confusion with the notion of horizontality given by the foliation on \mathbb{M} , in this paper it shall often simply be referred to as the D -lift to $\mathcal{O}(\mathbb{M})$.

3. HORIZONTAL CALCULUS OF VARIATIONS

To motivate the definition of the tangent processes to the horizontal Brownian motion on \mathbb{M} that we will use to prove quasi-invariance, we first present results on the horizontal calculus of variations of deterministic paths.

3.1. Adapted connections. Using the notation in Section 2.5, we consider $u \in \mathcal{O}(\mathbb{M})$. To take into account the foliation structure on \mathbb{M} , we shall be interested in a special subbundle of $\mathcal{O}(\mathbb{M})$, the *horizontal frame bundle*.

Definition 3.1. An isometry $u : (\mathbb{R}^{n+m}, \langle \cdot, \cdot \rangle) \rightarrow (T_x\mathbb{M}, g)$ will be called *horizontal* if $u(\mathbb{R}^n \times \{0\}) \subset \mathcal{H}_x$ and $u(\{0\} \times \mathbb{R}^m) \subset \mathcal{V}_x$. The *horizontal frame bundle* $\mathcal{O}_{\mathcal{H}}(\mathbb{M})$ is then defined as the set of $(x, u) \in \mathcal{O}(\mathbb{M})$ such that u is horizontal.

For notational convenience, when needed we identify \mathbb{R}^{n+m} with $\mathbb{R}^n \times \mathbb{R}^m$, hence we have embeddings of \mathbb{R}^n and \mathbb{R}^m into \mathbb{R}^{n+m} .

Assumption 1. We assume that D is a connection on \mathbb{M} satisfying the following properties.

- D is a *metric connection* on \mathbb{M} , that is, $Dg = 0$;
- D is *adapted to the foliation* \mathcal{F} in the following sense

$$\begin{aligned} D_X Y &\in \Gamma^\infty(\mathcal{H}), \text{ if } X \in \Gamma^\infty(\mathbb{M}), Y \in \Gamma^\infty(\mathcal{H}), \\ D_X Z &\in \Gamma^\infty(\mathcal{V}), \text{ if } X \in \Gamma^\infty(\mathbb{M}), Z \in \Gamma^\infty(\mathcal{V}); \end{aligned}$$

- For every $X \in \Gamma^\infty(\mathcal{H}), Y \in \Gamma^\infty(\mathbb{M})$, $D_X Y = \nabla_X Y$, where ∇ is the Bott connection.

Remark 3.2. In the case of a Riemannian submersion in Example 2.1, these assumptions imply that the connection D is a lift of the Levi-Civita connection on (\mathbb{B}, j) , namely,

$$\nabla_{\overline{X}}^{\mathbb{B}} \overline{Y} = \overline{D_X Y}, \quad \overline{X}, \overline{Y} \in \Gamma^\infty(\mathbb{B}),$$

where $\nabla^{\mathbb{B}}$ is the Levi-Civita connection on \mathbb{B} . We refer to Example 2.3 for further details.

Of course, an example of a connection D that satisfies the above assumptions is given by the Bott connection ∇ itself. However, we state the results of the section in greater generality using a connection D satisfying Assumption 1. This generality is relevant for Section 5, where we use other connections than the Bott connection (see Remark (5.2)). The main reason for using different connections is that while the Bott connection is

adapted to the foliation structure, the torsion of the Bott connection is not skew-symmetric.

The connection D allows us to lift vector fields on \mathbb{M} to vector fields on $\mathcal{O}(\mathbb{M})$ (see [34, p.421]). Let $e_1, \dots, e_n, f_1, \dots, f_m$ be the standard basis of \mathbb{R}^{n+m} .

Notation 3.3. We denote by A_i the vector field on $\mathcal{O}(\mathbb{M})$ such that $A_i(x, u)$ is the lift of $u(e_i)$, $i = 1, \dots, n$, $(x, u) \in \mathcal{O}(\mathbb{M})$, and we denote by V_j the vector field on $\mathcal{O}(\mathbb{M})$ such that $V_j(x, u)$ is the lift of $u(f_j)$, $j = 1, \dots, m$. We sometimes call A and V *fundamental vector fields* on $\mathcal{O}(\mathbb{M})$. For any $v \in \mathbb{R}^{n+m}$, we denote

$$Av := \sum_{i=1}^n v_i A_i,$$

$$Vv := \sum_{j=1}^m v_{j+n} V_j.$$

Then Av and Vv are vector fields on $\mathcal{O}(\mathbb{M})$ whose values at some $u \in \mathcal{O}(\mathbb{M})$ will be denoted respectively by $A_u v$ and $V_u v$.

Notation 3.4. Let x_0 be a fixed point in \mathbb{M} . By $W_0^\infty(\mathbb{R}^{n+m})$ we denote the space of smooth paths $v : [0, 1] \rightarrow \mathbb{R}^{n+m}$ such that $v(0) = 0$, and by $W_{x_0}^\infty(\mathbb{M})$ we denote the space of smooth paths $\gamma : [0, 1] \rightarrow \mathbb{M}$ such that $\gamma(0) = x_0$.

3.2. Development maps. Next we would like to recall the notion of a rolling map ϕ between path spaces over \mathbb{M} and \mathbb{R}^{n+m} or equivalently development and anti-development maps (see for instance [34, Section 2]). Let $\pi : \mathcal{O}(\mathbb{M}) \rightarrow \mathbb{M}$ be the bundle projection map. To define the rolling map $\phi : W_0^\infty(\mathbb{R}^{n+m}) \rightarrow W_{x_0}^\infty(\mathbb{M})$ we need the following differential equation on $\mathcal{O}(\mathbb{M})$

$$(3.1) \quad du_s = \sum_{i=1}^n A_i(u_s) d\omega_s^i + \sum_{i=1}^m V_i(u_s) d\omega_s^{n+i} = A_{u_s} d\omega_s + V_{u_s} d\omega_s,$$

where $\omega \in W_0^\infty(\mathbb{R}^{n+m})$. By compactness of \mathbb{M} and thus of $\mathcal{O}(\mathbb{M})$ this equation has a unique solution given an initial condition $u_0 \in \mathcal{O}(\mathbb{M})$. In the sequel we fix $u_0 \in \mathcal{O}(\mathbb{M})$ such that $\pi(u_0) = x_0$.

Definition 3.5.

- (1) For any $\omega \in W_0^\infty(\mathbb{R}^{n+m})$ the *development* of ω in \mathbb{M} is defined as $\gamma_s = \pi(u_s)$, where $\{u_s\}_{s \in [0,1]}$ is the solution to (3.1) with initial condition u_0 . Then we denote $\phi(\omega) := \gamma$. The map ϕ is also called the *rolling map*.
- (2) For any $\gamma \in W_{x_0}^\infty(\mathbb{M})$ the *anti-development* of γ in \mathbb{R}^{n+m} is the unique path $\omega \in W_0^\infty(\mathbb{R}^{n+m})$ such that if $\{u_s\}_{s \in [0,1]}$ is the solution to (3.1), then $\gamma_s = \pi(u_s)$. Then we denote $\phi^{-1}(\gamma) := \omega$.

This definition extends to continuous semimartingales, in which case we speak of stochastic development and stochastic anti-development (e.g. [38, Section 2.3] and [34, p. 433]).

3.3. Horizontal paths.

Definition 3.6. A smooth path $\omega : [0, 1] \rightarrow \mathbb{R}^{n+m}$ is called *horizontal* if it takes values in \mathbb{R}^n . The space of smooth horizontal paths such that $\omega(0) = 0$ will be denoted by $W_{0, \mathcal{H}}^\infty(\mathbb{R}^{n+m})$.

Definition 3.7. A smooth path $\gamma : [0, 1] \rightarrow \mathbb{M}$ is called *horizontal* if for every vertical smooth one-form θ we have $\int_\gamma \theta = 0$. The space of smooth horizontal paths such that $\gamma(0) = x_0$ will be denoted $W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$.

Remark 3.8. The space $W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$ contains only smooth paths, therefore it can be equivalently described as follows. A path γ is in $W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$ if and only if $\gamma'(s) \in \mathcal{H}_{\gamma(s)}$ for every $s \in [0, 1]$. The advantage of Definition 3.7 is that it will easily extend to non-smooth paths such as semimartingales.

The next step is to define the horizontal rolling map $\phi_{\mathcal{H}} : W_{0, \mathcal{H}}^\infty(\mathbb{R}^{n+m}) \rightarrow W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$ similarly to Definition 3.5 on the spaces of horizontal paths. For any $\omega^{\mathcal{H}} \in W_{0, \mathcal{H}}^\infty(\mathbb{R}^{n+m})$ we consider the differential equation on $\mathcal{O}(\mathbb{M})$ with initial condition u_0

$$(3.2) \quad du_s = \sum_{i=1}^n A_i(u_s) d\omega_s^{\mathcal{H}, i} = A_{u_i} d\omega_s^{\mathcal{H}}.$$

Observe that for $\gamma = \pi(u)$ we have

$$d\gamma_s = \sum_{i=1}^n d\pi(A_i(u_s)) d\omega_s^{\mathcal{H}, i},$$

and therefore γ is horizontal since $d\pi(A_i(u_s))$ is.

Lemma 3.9. Suppose $\gamma \in W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$, then there exists a unique $\omega^{\mathcal{H}} \in W_{0, \mathcal{H}}^\infty(\mathbb{R}^{n+m})$ such that if $\{u_s\}_{s \in [0, 1]}$ is the solution to (3.2), then $\gamma_s = \pi(u_s)$.

Proof. As before, let $e_1, \dots, e_n, f_1, \dots, f_m$ be the standard basis of \mathbb{R}^{n+m} . Note that any $\gamma \in W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$ can be viewed as an element in $W_{x_0}^\infty(\mathbb{M})$. Let $\omega \in W_0^\infty(\mathbb{R}^{n+m})$ be the anti-development of γ introduced in Definition 3.5. Then if $\{u_s\}_{s \in [0, 1]}$ is the solution to the differential equation (3.1) with initial condition u_0 , then $\gamma_s = \pi(u_s)$. Since γ is horizontal, then for every smooth vertical one-form θ one has

$$\int_{\gamma[0, s]} \theta = 0.$$

Therefore

$$\int_{\gamma[0, s]} \theta = \sum_{i=1}^n \int_0^s \theta(u_r e_i) d\omega_r^i + \sum_{i=1}^m \int_0^s \theta(u_r f_i) d\omega_r^{n+i} = 0.$$

The form θ being vertical, one deduces

$$\sum_{i=1}^m \int_0^t \theta(u_s f_i) d\omega_s^{n+i} = 0.$$

Since it is true for any θ , one deduces

$$\sum_{i=1}^m \int_0^s (u_r f_i) d\omega_r^{n+i} = 0.$$

Now observe that $u_r f_1, \dots, u_r f_m$ are linearly independent, thus for every r one has $d\omega_r^{n+i} = 0$. As a conclusion, ω is horizontal. \square

Definition 3.10 (Horizontal development and anti-development).

- (1) For any $\omega^{\mathcal{H}} \in W_{0, \mathcal{H}}^\infty(\mathbb{R}^{n+m})$ the *horizontal development* of ω in \mathbb{M} is $\gamma_s = \pi(u_s)$, where $\{u_s\}_{s \in [0,1]}$ is the solution to (3.2) with initial condition $u_0 \in \mathcal{O}(\mathbb{M})$. Then we denote $\phi_{\mathcal{H}}(\omega) := \gamma$. The map $\phi_{\mathcal{H}}$ is called the *horizontal rolling map*.
- (2) For any $\gamma \in W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$ the *horizontal anti-development* of γ is $\omega^{\mathcal{H}} \in W_{0, \mathcal{H}}^\infty(\mathbb{R}^{n+m})$ is the unique path such that if $\{u_s\}_{s \in [0,1]}$ is the solution to (3.2) with initial condition u_0 , then $\gamma_s = \pi(u_s)$. Then we denote $\phi_{\mathcal{H}}^{-1}(\gamma) := \omega$.

3.4. Paths tangent to horizontal paths. For any $v \in W_0^\infty(\mathbb{R}^{n+m})$ we consider the vector field \mathbf{D}_v on $W_{x_0}^\infty(\mathbb{M})$ defined by

$$\mathbf{D}_v(\gamma)_s := u_s(\gamma)v_s, \quad \gamma \in W^\infty(\mathbb{M}),$$

where u is the D -lift of γ to $\mathcal{O}(\mathbb{M})$. Let $\{\zeta_t^v, t \in \mathbb{R}\}$ be the flow generated by \mathbf{D}_v , i.e.

$$\frac{d}{dt}(\zeta_t^v \gamma)_s = \mathbf{D}_v(\zeta_t^v \gamma)_s, \quad \zeta_0^v \gamma = \gamma.$$

One can use the development and anti-development maps ϕ and ϕ^{-1} in Definition 3.5 to introduce a flow on $W_0^\infty(\mathbb{R}^{n+m})$ as follows

$$\xi_t^v := \phi^{-1} \circ \zeta_t^v \circ \phi, \quad t \in \mathbb{R}.$$

Note that \mathbf{D}_v , ζ_t^v and ξ_t^v depend on the connection D . We now recall [34, Theorem 2.1] that describes the generator of the flow ξ_t^v in the situation when a connection is metric-compatible but not necessarily torsion-free.

Theorem 3.11 (Theorem 2.1 in [34]). Suppose that $v \in W_0^\infty(\mathbb{R}^{n+m})$ and $\omega \in W_0^\infty(\mathbb{R}^{n+m})$. Then

$$\left. \frac{d}{dt} \right|_{t=0} \xi_t^v(\omega)_s = p_v(\omega)_s,$$

where

$$p_v(\omega)_s = v(s) - \int_0^s T_{u_r}^D(Ad\omega_r + Vd\omega_r, Av(r) + Vv(r)) - \int_0^s \left(\int_0^r \Omega_{u_\tau}^D(Ad\omega_\tau + Vd\omega_\tau, Av(\tau) + Vv(\tau)) \right) d\omega_r.$$

Here u is the D -lift to $\mathcal{O}(\mathbb{M})$ of the development of ω , T^D is the torsion form of the connection D and Ω^D is its curvature form.

We are interested in the variation of horizontal paths. Let us observe that for $\omega^{\mathcal{H}} \in W_{0, \mathcal{H}}^\infty(\mathbb{R}^{n+m})$

$$(3.3) \quad p_v(\omega^{\mathcal{H}})_s = v(s) - \int_0^s T_{u_r}^D(Ad\omega_r^{\mathcal{H}}, Av(r) + Vv(r)) - \int_0^s \left(\int_0^r \Omega_{u_\tau}^D(Ad\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) d\omega_r^{\mathcal{H}}.$$

Definition 3.12. We will say that $v \in W_0^\infty(\mathbb{R}^{n+m})$ is *tangent to the horizontal path* $\gamma \in W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$ if for every $s \in [0, 1]$, $\frac{d}{dt}|_{t=0} \phi^{-1}(\zeta_t^v \gamma)_s \in \mathbb{R}^n$.

Remark 3.13. From this definition, $v \in W_0^\infty(\mathbb{R}^{n+m})$ is tangent to the horizontal path γ if and only if $p_v(\omega^{\mathcal{H}})$ is horizontal, where ω is the horizontal anti-development of γ . Intuitively, v is tangent to γ if it yields a variation of γ in the horizontal directions only. More precisely, call a vector field ξ along $\gamma \in W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$ an *horizontal variation* of γ if $\xi(x_0) = 0$ and if there exists $(\sigma_t)_{t \in [-\varepsilon, \varepsilon]} \subset W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$ with $\sigma_0 = \gamma$ such that $\frac{d}{dt}|_{t=0} (\sigma_t)_s = \xi_s$ for $s \in [0, 1]$. Then, by Theorem 3.11 and Proposition 3.20, ξ is an horizontal variation of γ if and only if $u_s(\gamma)^{-1} \xi_s$ is tangent to the horizontal path γ . Let us note that the notion of horizontal variation is independent from any metric and any connection. It therefore yields an intrinsic notion of horizontal tangent path space. We are grateful to the referee for this observation.

Remark 3.14. One should note that even if $v \in W_0^\infty(\mathbb{R}^{n+m})$ is tangent to the horizontal path γ , it may not be true that for every $t \in \mathbb{R}$, $\zeta_t^v \gamma \in W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$.

One has the following characterization of tangent paths, which is the main result of the section.

Theorem 3.15. Let $\gamma \in W_{x_0, \mathcal{H}}^\infty(\mathbb{M})$. A path $v \in W_0^\infty(\mathbb{R}^{n+m})$ is tangent to the horizontal path γ if and only if the path

$$v(s) - \int_0^s T_{u_r}(Ad\omega_r^{\mathcal{H}}, Av(r))$$

is horizontal, i.e. takes values in \mathbb{R}^n , where $\omega^{\mathcal{H}}$ is the horizontal anti-development of γ , u is its D -lift to $\mathcal{O}(\mathbb{M})$, and T is the torsion of the Bott connection.

Proof. The path $v \in W_0^\infty(\mathbb{R}^{n+m})$ is tangent to the horizontal path γ if and only if the path

$$\begin{aligned} p_v(\omega^{\mathcal{H}})_s &= v(s) - \int_0^s T_{u_s}^D(Ad\omega_r^{\mathcal{H}}, Av(r) + Vv(r)) \\ &\quad - \int_0^s \left(\int_0^r \Omega_{u_\tau}^D(Ad\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) d\omega_r^{\mathcal{H}} \end{aligned}$$

is horizontal. Since D satisfies Assumption 1, the integral

$$\int_0^s \left(\int_0^r \Omega_{u_\tau}^D(Ad\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) d\omega_r^{\mathcal{H}}$$

is always horizontal. Let us now denote by J the difference between connections D and ∇ , that is, the tensor J is defined for any $X, Y \in \Gamma^\infty(\mathbb{M})$ by

$$J_X Y = D_X Y - \nabla_X Y.$$

We have then

$$\begin{aligned} T^D(X, Y) &= D_X Y - D_Y X - [X, Y] \\ &= T(X, Y) + J_X Y - J_Y X. \end{aligned}$$

Let us assume that X is horizontal. We have then $J_X = 0$, because $D_{\mathcal{H}} = \nabla_{\mathcal{H}}$. Also $J_Y X$ is horizontal, because D is adapted to the foliation \mathcal{F} . We deduce that the vertical part of

$$v(s) - \int_0^s T_{u_r}^D(Ad\omega_r^{\mathcal{H}}, Av(r) + Vv(r))$$

is the same as the vertical part of

$$v(s) - \int_0^s T_{u_r}(Ad\omega_r^{\mathcal{H}}, Av(r) + Vv(r)).$$

We conclude that the vertical part of $p_v(\omega^{\mathcal{H}})$ is zero if and only if the vertical part of

$$v(s) - \int_0^s T_{u_r}(Ad\omega_r^{\mathcal{H}}, Av(r) + Vv(r)).$$

is zero. By the properties in Equation (2.1), we have

$$\int_0^s T_{u_r}(Ad\omega_r^{\mathcal{H}}, Vv(r)) = 0,$$

which concludes the proof. \square

Remark 3.16. By Theorem 3.15, the notion of tangent path does not depend on the particular choice of the connection D as long as it satisfies Assumption 1.

3.5. Variations on the horizontal path space. In this section, we describe two types of variations on the horizontal path space that are induced by tangent paths. The first one is explicit and inspired by the approach by B. Driver in [19]. The second one is based on more classical flow constructions. The key ingredient is the following lemma.

Lemma 3.17. Let $h \in W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m})$. If $\omega^{\mathcal{H}} \in W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m})$, then

$$(3.4) \quad \tau_h(\omega^{\mathcal{H}})_s = h(s) + \int_0^s T_{u_r}(Ad\omega_r^{\mathcal{H}}, Ah(r))$$

is a tangent path to $\phi(\omega^{\mathcal{H}})$, where u denotes the D -lift of the horizontal development of $\omega^{\mathcal{H}}$.

Proof. Let

$$v(s) = h(s) + \int_0^s T_{u_r}(Ad\omega_r^{\mathcal{H}}, Ah(r)).$$

Since h is horizontal and T is a vertical tensor, one deduces that the horizontal part of v is h . Therefore,

$$v(s) - \int_0^s T_{u_r}(Ad\omega_r^{\mathcal{H}}, Av(r)) = h(s)$$

is horizontal. □

Let $v \in W_0^\infty(\mathbb{R}^{n+m})$, $\omega^{\mathcal{H}} \in W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m})$ and assume that v is tangent to the horizontal development of $\omega^{\mathcal{H}}$. Recall that

$$\begin{aligned} p_v(\omega^{\mathcal{H}})_s &= v(s) - \int_0^s T_{u_s}^D(Ad\omega_s^{\mathcal{H}}, Av(s) + Vv(s)) \\ &\quad - \int_0^s \left(\int_0^\tau \Omega_{u_\tau}^D(Ad\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) d\omega_\tau^{\mathcal{H}}. \end{aligned}$$

As before, let us now denote by J the difference between connections D and ∇ . For $X, Y \in \Gamma^\infty(\mathbb{M})$, we have thus

$$J_X Y = D_X Y - \nabla_X Y.$$

We can then write

$$\begin{aligned} p_v(\omega^{\mathcal{H}})_s &= v_{\mathcal{H}}(s) + \int_0^s (J_{Vv(r)})_{u_r}(Ad\omega_r^{\mathcal{H}}) \\ &\quad - \int_0^s \left(\int_0^\tau \Omega_{u_\tau}^D(Ad\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) d\omega_\tau^{\mathcal{H}}. \end{aligned}$$

More concisely, we have therefore

$$p_v(\omega^{\mathcal{H}})_s = v_{\mathcal{H}}(s) + \int_0^s q_v(\omega^{\mathcal{H}})_u d\omega_u^{\mathcal{H}},$$

where $q_v(\omega^{\mathcal{H}})_u \in \mathfrak{so}(n)$ is defined in such a way that

$$\begin{aligned} & \int_0^s q_v(\omega^{\mathcal{H}})_u d\omega_u^{\mathcal{H}} \\ &= \int_0^s (J_{Vv(r)})_{u_r} (Ad\omega_r^{\mathcal{H}}) - \int_0^s \left(\int_0^r \Omega_{u_\tau}^D (Ad\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) d\omega_r^{\mathcal{H}}. \end{aligned}$$

As a consequence, with the above notation, one has that for every $h \in W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m})$

$$p_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s = h(s) + \int_0^s q_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_u d\omega_u^{\mathcal{H}},$$

We are now ready to introduce two relevant variations of horizontal paths.

Notation 3.18. Let $h \in W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m})$.

(1) For $t \in \mathbb{R}$, we define a map $\rho_t^h : W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m}) \rightarrow W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m})$ by

$$(3.5) \quad (\rho_t^h \omega^{\mathcal{H}})_s := \int_0^s e^{tq_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_u} d\omega_u^{\mathcal{H}} + th(s).$$

(2) For $t \in \mathbb{R}$, we define a map $\nu_t^h : W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m}) \rightarrow W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m})$ as the flow generated by p_{τ_h}

$$\frac{d}{dt}(\nu_t^h \omega^{\mathcal{H}})_s = p_{\tau_h(\nu_t^h \omega^{\mathcal{H}})}(\nu_t^h \omega^{\mathcal{H}})_s, \quad \nu_0^h \omega^{\mathcal{H}} = \omega^{\mathcal{H}}.$$

Remark 3.19. Unless $q_{\tau_h} = 0$, the family $\{\rho_t^h, t \in \mathbb{R}\}$ is **not** a flow on $W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m})$, but it is a convenient explicit one-parameter variation, since we observe that $\rho_0^h \omega^{\mathcal{H}} = \omega^{\mathcal{H}}$ and

$$\left. \frac{d}{dt} \right|_{t=0} (\rho_t^h \omega^{\mathcal{H}})_s = p_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s.$$

We then have the following result, which is immediate in view of Theorem 3.11 since

$$\left. \frac{d}{dt} \right|_{t=0} (\rho_t^h \omega^{\mathcal{H}})_s = \left. \frac{d}{dt} \right|_{t=0} (\nu_t^h \omega^{\mathcal{H}})_s = p_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s.$$

Proposition 3.20 (Variation of horizontal paths along tangent paths).

Let $h \in W_{0,\mathcal{H}}^\infty(\mathbb{R}^{n+m})$, then for every $\gamma \in W_{\mathcal{H}}^\infty(\mathbb{M})$

$$\left. \frac{d}{dt} \right|_{t=0} \phi_{\mathcal{H}} \circ \rho_t^h \circ \phi_{\mathcal{H}}^{-1}(\gamma)_s = \left. \frac{d}{dt} \right|_{t=0} \phi_{\mathcal{H}} \circ \nu_t^h \circ \phi_{\mathcal{H}}^{-1}(\gamma)_s = u_s(\gamma) \tau_h(\omega^{\mathcal{H}})_s,$$

where u is the D -lift of γ , and $\omega^{\mathcal{H}}$ is its horizontal development.

4. QUASI-INVARIANCE OF THE HORIZONTAL WIENER MEASURE

In this part of the paper, we first describe two constructions of the horizontal Brownian motion, and then we develop horizontal stochastic calculus and prove quasi-invariance of the horizontal Wiener measure. Throughout this section we consider a smooth connected $n + m$ -dimensional Riemannian manifold \mathbb{M} equipped with the structure of an m -dimensional foliation \mathcal{F} , a bundle-like metric g and totally geodesic m -dimensional leaves. In addition, we assume that \mathbb{M} is compact.

4.1. Horizontal Brownian motion.

4.1.1. *Construction from the horizontal Dirichlet form.* We define the horizontal gradient $\nabla_{\mathcal{H}}f$ of a smooth function f as the projection of the Riemannian gradient of f on the horizontal bundle \mathcal{H} . Similarly, we define the vertical gradient $\nabla_{\mathcal{V}}f$ of a function f as the projection of the Riemannian gradient of f on the vertical bundle \mathcal{V} .

Consider the pre-Dirichlet form

$$\mathcal{E}_{\mathcal{H}}(f, h) = \int_{\mathbb{M}} g_{\mathcal{H}}(\nabla_{\mathcal{H}}f, \nabla_{\mathcal{H}}h) \, d\text{Vol}, \quad f, h \in C^{\infty}(\mathbb{M}),$$

where $d\text{Vol}$ is the Riemannian volume measure on \mathbb{M} . We note that $\mathcal{E}_{\mathcal{H}}$ is closable since it can be dominated by the Dirichlet form generated by the Laplace-Beltrami on \mathbb{M} which is closable since \mathbb{M} is compact, thus complete. Then there exists a unique diffusion operator L on \mathbb{M} such that for all $f, h \in C^{\infty}(\mathbb{M})$

$$\mathcal{E}_{\mathcal{H}}(f, h) = - \int_{\mathbb{M}} fLh \, d\text{Vol} = - \int_{\mathbb{M}} hLf \, d\text{Vol}.$$

The operator L is called the *horizontal Laplacian* of the foliation. If $\{X_i\}_{i=1}^n$ is a local orthonormal frame of horizontal vector fields, then we can write L in this frame

$$(4.1) \quad L = \sum_{i=1}^n X_i^2 + X_0,$$

where X_0 is a smooth vector field. Observe that the subbundle \mathcal{H} satisfies Hörmander's (bracket generating) condition, therefore by Hörmander's theorem the operator L is locally subelliptic (for comments on this terminology introduced by Fefferman-Phong we refer to [28], see also the survey papers [3, 40] or [20, p. 944]).

By [1, Proposition 5.1] the completeness of the Riemannian metric g implies that L is essentially self-adjoint on $C^{\infty}(\mathbb{M})$ and thus that $\mathcal{E}_{\mathcal{H}}$ is uniquely closable. Then we can define the semigroup $P_s = e^{\frac{s}{2}L}$ by using the spectral theorem. The diffusion process $\{W_s\}_{s \geq 0}$ corresponding to the semigroup $\{P_s\}_{s \geq 0}$ will be called the *horizontal Brownian motion* on the Riemannian foliation (\mathcal{F}, g) . Since \mathbb{M} is assumed to be compact, $1 \in \text{dom}(\mathcal{E}_{\mathcal{H}})$ and thus $P_s 1 = 1$. This implies that $\{W_s\}_{s \geq 0}$ is a non-explosive diffusion.

If the horizontal Laplacian can be written in the form 4.1 globally for smooth horizontal vector fields X_0, X_1, \dots, X_n , then $\{W_t\}_{t \geq 0}$ can be constructed from a stochastic differential equation on \mathbb{M} .

Even if the horizontal Laplacian can not be written in the form 4.1 globally, the horizontal Brownian motion $\{W_s\}_{s \geq 0}$ can still be constructed from a globally defined stochastic differential equation on a bundle over \mathbb{M} (see [21, Theorem 3.8] or Corollary 4.4). The following section provides an explicit description of such a construction that shall be used in the sequel.

4.1.2. *Construction from the orthonormal frame bundle.* We can write the vector fields $\{A_i\}_{i=1}^n$ locally in terms of the normal frames introduced in [7].

Lemma 4.1 (Lemma 2.2 in [7]). Let $x_0 \in \mathbb{M}$. Around x_0 , there exist a local orthonormal horizontal frame $\{X_1, \dots, X_n\}$ and a local orthonormal vertical frame $\{Z_1, \dots, Z_m\}$ such that the following structure relations hold

$$\begin{aligned} [X_i, X_j] &= \sum_{k=1}^n \omega_{ij}^k X_k + \sum_{k=1}^m \gamma_{ij}^k Z_k, \\ [X_i, Z_k] &= \sum_{j=1}^m \beta_{ik}^j Z_j, \end{aligned}$$

where $\omega_{ij}^k, \gamma_{ij}^k, \beta_{ik}^j$ are smooth functions such that

$$\beta_{ik}^j = -\beta_{ij}^k.$$

Moreover, at x_0 we have

$$\omega_{ij}^k = 0, \beta_{ij}^k = 0.$$

We will also need the fact (see [7, p. 918]) that in this frame the Christoffel symbols of the Bott connection ∇ are given by

$$\begin{aligned} \nabla_{X_i} X_j &= \frac{1}{2} \sum_{k=1}^n \left(\omega_{ij}^k + \omega_{ki}^j + \omega_{kj}^i \right) X_k, \\ \nabla_{Z_j} X_i &= 0, \\ \nabla_{X_i} Z_j &= \sum_{k=1}^m \beta_{ij}^k Z_k. \end{aligned}$$

Thus, from the assumption that $D_{\mathcal{H}} = \nabla_{\mathcal{H}}$, we have

$$\begin{aligned} D_{X_i} X_j &= \frac{1}{2} \sum_{k=1}^n \left(\omega_{ij}^k + \omega_{ki}^j + \omega_{kj}^i \right) X_k, \\ D_{X_i} Z_j &= \sum_{k=1}^m \beta_{ij}^k Z_k. \end{aligned}$$

For $x_0 \in \mathbb{M}$ we let $\{X_1, \dots, X_n, Z_1, \dots, Z_m\}$ be a normal frame around x_0 . If $u \in \mathcal{O}_{\mathcal{H}}(\mathbb{M})$ is a horizontal isometry, we can find an orthogonal matrix

$\{e_i^j\}_{i,j=1}^n$ such that $u(e_i) = \sum_{j=1}^n e_i^j X_j$, and $u(f_i) = \sum_{j=1}^m f_i^j Z_j$ for f_i^j , $i = 1, \dots, n$, $j = 1, \dots, m$. Let \bar{X}_j be the vector field on $\mathcal{O}_{\mathcal{H}}(\mathbb{M})$ defined by

$$\bar{X}_j f(x, u) = \lim_{t \rightarrow 0} \frac{f(e^{tX_j}(x), u) - f(x, u)}{t},$$

where $e^{tX_j}(x)$ is the exponential map on M .

Lemma 4.2. Let $x_0 \in \mathbb{M}$ and $(x, u) \in \mathcal{O}_{\mathcal{H}}(\mathbb{M})$, then

$$\begin{aligned} A_i(x, u) &= \sum_{j=1}^n e_i^j \bar{X}_j \\ &- \sum_{j,k,l,r=1}^n e_i^j e_r^l \langle D_{X_j} X_l, X_k \rangle \frac{\partial}{\partial e_r^k} - \sum_{j=1}^n \sum_{k,l,r=1}^m e_i^j f_r^l \langle D_{X_j} Z_l, Z_k \rangle \frac{\partial}{\partial f_r^k}. \end{aligned}$$

In particular, at x_0 we have

$$A_i(x_0, u) = \sum_{j=1}^n e_i^j \bar{X}_j.$$

Proof. Let $u : \mathbb{R}^{n+m} \rightarrow T_x \mathbb{M}$ be a horizontal isometry and $x(t)$ be a smooth curve in \mathbb{M} such that $x(0) = x$ and $x'(0) = u(e_i)$. We denote by $x^*(t) = (x(t), u(t))$ the D -lift to $\mathcal{O}(\mathbb{M})$ of $x(t)$ and by $x'_1(t), \dots, x'_n(t)$ the components of $x'(t)$ in the horizontal frame X_1, \dots, X_n . Since D is adapted to the foliation \mathcal{F} , the curve $x^*(t)$ takes its values in the horizontal frame bundle $\mathcal{O}_{\mathcal{H}}(\mathbb{M})$. By definition of A_i , one has

$$A_i = \sum_{j=1}^n x'_j(0) \bar{X}_j + \sum_{k,l=1}^n u'_{kl}(0) \frac{\partial}{\partial e_k^l} + \sum_{k,l=1}^m v'_{kl}(0) \frac{\partial}{\partial f_k^l},$$

where $u_{kl}(t) = \langle u(t)(e_k), X_l \rangle$ and $v_{kl}(t) = \langle u(t)(f_k), Z_l \rangle$. Since $u(t)(e_k)$ and $u(t)(f_k)$ are parallel along $x(t)$, one has

$$D_{x'(t)} u(t)(e_k) = 0, \quad D_{x'(t)} u(t)(f_k) = 0.$$

At $t = 0$, this yields the expected result. \square

In particular, Lemma 4.2 implies the following statement.

Proposition 4.3. Let $\pi : \mathcal{O}(\mathbb{M}) \rightarrow \mathbb{M}$ be the bundle projection map. For a smooth $f : \mathbb{M} \rightarrow \mathbb{R}$, and $(x, u) \in \mathcal{O}_{\mathcal{H}}(\mathbb{M})$,

$$\left(\sum_{i=1}^n A_i^2 \right) (f \circ \pi)(x, u) = Lf \circ \pi(x, u).$$

Proof. It is enough to prove this identity at x_0 . Using the fact that at x_0 we have $\langle D_{X_j} X_l, X_k \rangle = \langle D_{X_j} Z_l, Z_k \rangle = 0$, we see that

$$\sum_{i=1}^n A_i^2 = \sum_{j=1}^n \bar{X}_j^2.$$

The conclusion follows. \square

As a straightforward corollary, we can introduce the *horizontal* Brownian motion as follows.

Corollary 4.4. Let $(\Omega, (\mathcal{F}_s)_{s \geq 0}, \mathbb{P})$ be a filtered probability space that satisfies the usual conditions and let $\{B_s\}_{s \geq 0}$ be an adapted \mathbb{R}^n -valued Brownian motion on that space. Let $\{U_s\}_{s \geq 0}$ be a solution to the Stratonovich stochastic differential equation

$$(4.2) \quad dU_s = \sum_{i=1}^n A_i(U_s) \circ dB_s^i = A_{U_s} \circ dB_s, \quad U_0 \in \mathcal{O}_{\mathcal{H}}(\mathbb{M}),$$

then $W_s = \pi(U_s)$ is a *horizontal Brownian motion* on \mathbb{M} , that is, a Markov process with the generator $\frac{1}{2}L$. Here we used Notation 3.3 and identified the \mathbb{R}^n -valued Brownian motion $\{B_s\}_{s \geq 0}$ with an \mathbb{R}^{n+m} -valued process $(B_s, 0)$.

4.2. Horizontal semimartingales. Let $(\Omega, (\mathcal{F}_s)_{s \geq 0}, \mathbb{P})$ be a filtered probability space that satisfies the usual conditions.

Definition 4.5. An \mathbb{R}^{n+m} -valued \mathcal{F}_s -adapted continuous semimartingale $(W_s)_{s \geq 0}$ is called *horizontal* if for all $s \geq 0$

$$\mathbb{P}(W_s \in \mathbb{R}^n \times \{0\}) = 1.$$

The space of horizontal semimartingales with $W_0 = 0$ will be denoted by $SW_{\mathcal{H}}(\mathbb{R}^{n+m})$.

Definition 4.6. An \mathbb{M} -valued \mathcal{F}_s -adapted continuous semimartingale $\{M_s\}_{s \geq 0}$ is called *horizontal* if for every vertical smooth one-form θ , and every $s \geq 0$ the Stratonovich stochastic line integral $\int_{M[0,s]} \theta = 0$ almost surely. The space of horizontal semimartingales such that $M_0 = x_0$ will be denoted by $SW_{\mathcal{H}}(\mathbb{M})$.

Remark 4.7. We refer to [38, Section 2.4, Definition 2.4.1] for the definition of Stratonovich stochastic line integrals.

Then we have the following result, whose proof is essentially identical to the proof of Lemma 3.9 and thus omitted for conciseness.

Proposition 4.8. As before π is the bundle projection map $\mathcal{O}_{\mathcal{H}}(\mathbb{M}) \rightarrow \mathbb{M}$.

- (1) Let $\{W_s\}_{s \geq 0} \in SW_{\mathcal{H}}(\mathbb{R}^{n+m})$ and let $\{U_s\}_{s \geq 0}$ be the solution to the Stratonovich stochastic differential equation

$$dU_s = \sum_{i=1}^n A_i(U_s) \circ dW_s^i = A_{U_s} \circ dW_s, \quad U_0 \in \mathcal{O}_{\mathcal{H}}(\mathbb{M}),$$

then $M_s := \pi(U_s)$ is a horizontal semimartingale on \mathbb{M} .

- (2) Let $\{M_s\}_{s \geq 0} \in SW_{\mathcal{H}}(\mathbb{M})$. Then there exists a unique $\{W_s\}_{s \geq 0} \in SW_{\mathcal{H}}(\mathbb{R}^{n+m})$ such that if $\{U_s\}_{s \geq 0}$ is the solution to the Stratonovich stochastic differential equation

$$dU_s = \sum_{i=1}^n A_i(U_s) \circ dW_s^i = A_{U_s} \circ dW_s, \quad U_0 \in \mathcal{O}_{\mathcal{H}}(\mathbb{M}),$$

then $M_s = \pi(U_s)$.

Here we used Notation 3.3, where we introduced how fundamental vector fields A and V on $\mathcal{O}(\mathbb{M})$ in Notation 3.3 act on vectors in \mathbb{R}^{n+m} . Note that A acts on $\mathbb{R}^n \times \{0\}$ in \mathbb{R}^{n+m} , and so we can apply it to $\omega_s^{\mathcal{H}}$.

Proposition 4.8 allows us to introduce the following notion.

Definition 4.9. Suppose $\{W_s\}_{s \geq 0}$ and $\{M_s\}_{s \geq 0}$ are as in Proposition 4.8. Then

- (1) $\{M_s\}_{s \geq 0}$ is called the *stochastic horizontal development* of $\{W_s\}_{s \geq 0}$, and we denote $\phi_{\mathcal{H}}(W) := M$.
- (2) The path $\{W_s\}_{s \geq 0}$ is called the *stochastic horizontal anti-development* of $\{M_s\}_{s \geq 0}$, and we denote $\phi_{\mathcal{H}}^{-1}(M) := W$.

As a consequence, one deduces that the horizontal Brownian motion constructed in Corollary 4.4 is a horizontal semimartingale.

Definition 4.10. The *horizontal Itô map* (or *horizontal stochastic development map*) is the following adapted map defined $\mu_{\mathcal{H}}$ -a.s.

$$\begin{aligned} \phi_{\mathcal{H}} : W_0(\mathbb{R}^n) &\longrightarrow W_{x_0}(\mathbb{M}), \\ \omega^{\mathcal{H}} &\longmapsto W \end{aligned}$$

By using Proposition 4.8 and arguing as in [34, p. 433], one can construct an adapted map $\phi_{\mathcal{H}}^{-1} : W_{x_0}(\mathbb{M}) \rightarrow W_0(\mathbb{R}^n)$ defined μ_W -a.s. We will call $\phi_{\mathcal{H}}^{-1}$ the *stochastic horizontal anti-development map*.

We also refer to [16, Definition 2.5] for a discussion of the Itô map in the Riemannian setting and to the previous section for explicit constructions in our setting.

Remark 4.11. If one uses a Dirichlet form to construct the horizontal Brownian motion as in Section 4.1.1, then it does not straightforward to prove that one obtains a semimartingale. In particular, a standard approach such as the proof of [38, Theorem 3.2.1] does not readily extend to our setting.

4.3. Quasi-invariance of the horizontal Wiener measure. In this section we prove quasi-invariance of the law of the horizontal Brownian motion with respect to variations generated by suitable *tangent processes*. Our argument follows relatively closely the one by B. Driver [17] and then E. Hsu in [34] (see also [14, 15, 25]). More precisely, we will describe two types of variation of the horizontal Brownian motion paths with respect to which the horizontal Wiener measure is quasi-invariant. The first one is largely

inspired by Driver [19, Theorem 7.28]. It is explicit, see Equation (4.7) and readily yields the integration by parts formula in Section 4.3.5, but does not induce a flow. The second type of variation induces a flow and yields the sub-Riemannian analogue of [34, Theorem 4.1].

4.3.1. Framework. We will use the same framework and notation as before. In particular, we still consider an arbitrary connection D on \mathbb{M} that satisfies the properties in Assumption 1. In addition we now introduce notation needed to establish quasi-invariance. We will mainly follow the presentation in [16, 19, 34].

We work in the probability space $(W_0(\mathbb{R}^n), \mathcal{B}, \mu_{\mathcal{H}})$, where $W_0(\mathbb{R}^n)$ is the space of continuous functions $\omega^{\mathcal{H}} : [0, 1] \rightarrow \mathbb{R}^n$ such that $\omega^{\mathcal{H}}(0) = 0$, \mathcal{B} is the Borel σ -field on the path space $W_0(\mathbb{R}^n)$, and $\mu_{\mathcal{H}}$ is the Wiener measure. The coordinate process $(\omega_s^{\mathcal{H}})_{0 \leq s \leq 1}$ is therefore a Brownian motion in \mathbb{R}^n . The usual completion of the natural filtration generated by $\{\omega_s^{\mathcal{H}}\}_{0 \leq s \leq 1}$ will be denoted by \mathcal{B}_s .

We use the subscripts or superscripts \mathcal{H} , because, as before, \mathbb{R}^n is identified with the subspace $\mathbb{R}^n \times \{0\} \subset \mathbb{R}^{n+m}$. The \mathbb{R}^{n+m} -valued process $(\omega_s^{\mathcal{H}}, 0)$ will be referred to as a *horizontal Brownian motion*. The process $\{W_t\}_{0 \leq t \leq 1}$ constructed using Corollary 4.4 is the horizontal Brownian motion and the law μ_W of the horizontal Brownian motion on \mathbb{M} will be referred to as the *horizontal Wiener measure* on \mathbb{M} . Therefore, μ_W is a probability measure on the space $W_{x_0}(\mathbb{M})$ of continuous paths $w : [0, 1] \rightarrow \mathbb{M}$ with $w(0) = x_0$.

Remark 4.12. If the horizontal Laplacian can be written in Hörmander's form globally as in 4.1, then by [48, Corollary 5.4] the support of the horizontal Wiener measure μ_W is $W_{x_0}(\mathbb{M})$ itself.

4.3.2. Tangent processes to the horizontal Brownian motion. We now introduce the relevant class of tangent processes to the horizontal Brownian motion. To prove quasi-invariance, we consider the following class of tangent processes.

Definition 4.13. We define the *horizontal Cameron-Martin space* denoted by $\mathcal{CM}_{\mathcal{H}}(\mathbb{R}^{n+m})$ as the space of absolutely continuous \mathbb{R}^n -valued (deterministic) functions $\{h(s)\}_{0 \leq s \leq 1}$ such that $h(0) = 0$ and

$$\int_0^1 |h'(s)|_{\mathbb{R}^n}^2 ds < \infty.$$

Definition 4.14. Suppose $\{v(s)\}_{0 \leq s \leq 1}$ is a \mathcal{B}_s -adapted \mathbb{R}^{n+m} -valued continuous semimartingale such that

$$(4.3) \quad v(0) = 0 \text{ and } \mathbb{E} \left(\int_0^1 |v(s)|_{\mathbb{R}^{n+m}}^2 ds \right) < \infty.$$

The semimartingale $\{v(s)\}_{0 \leq s \leq 1}$ will be called a *tangent process* to the horizontal Brownian motion if the process

$$v(s) - \int_0^s T_{U_r}(A \circ d\omega_r^{\mathcal{H}}, Av(r))$$

is a horizontal Cameron-Martin path, where T denotes the torsion form of the Bott connection (not D). The space of tangent processes to the horizontal Brownian motion will be denoted by $TW_{\mathcal{H}}(\mathbb{M})$.

Remark 4.15. In Definition 4.14 we used the torsion T of the Bott connection. Observe that since T is a vertical tensor, a \mathcal{B}_s -adapted \mathbb{R}^{n+m} -valued continuous semimartingale $\{v(s)\}_{0 \leq s \leq 1}$ satisfying (4.3) is in $TW_{\mathcal{H}}(\mathbb{M})$ if and only if

- (1) The horizontal part $v_{\mathcal{H}}$ is in $\mathcal{CM}_{\mathcal{H}}(\mathbb{R}^{n+m})$;
- (2) The vertical part $v_{\mathcal{V}}$ is given by

$$v_{\mathcal{V}}(s) = \int_0^s T_{U_r}(A \circ d\omega_r^{\mathcal{H}}, Av_{\mathcal{H}}(r)).$$

As a consequence, for any $h \in \mathcal{CM}_{\mathcal{H}}(\mathbb{R}^{n+m})$,

$$(4.4) \quad \tau^h(\omega^{\mathcal{H}})_s = h(s) + \int_0^s T_{U_r}(A \circ d\omega_r^{\mathcal{H}}, Ah(r))$$

is a tangent process to the horizontal Brownian motion.

Notation 4.16. If $v \in TW_{\mathcal{H}}(\mathbb{M})$ is a tangent process, we denote

$$\begin{aligned} p_v(\omega^{\mathcal{H}})_s &:= v(s) - \int_0^s T_{U_r}^D(A \circ d\omega_r^{\mathcal{H}}, Av(r) + Vv(r)) \\ &\quad - \int_0^s \left(\int_0^r \Omega_{U_\tau}^D(A \circ d\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) \circ d\omega_r^{\mathcal{H}}, \end{aligned}$$

where Ω^D is the curvature form of the connection D .

This definition comes from Equation (3.3), where $d\omega^{\mathcal{H}}$ is replaced by the Stratonovich differential $\circ d\omega^{\mathcal{H}}$. Since D is a horizontal metric connection, the stochastic integral $\int_0^s \Omega_{U_\tau}^D(A \circ d\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau))$ restricts to \mathbb{R}^n as a skew-symmetric endomorphism of \mathbb{R}^n . Also, from the proof of Theorem 3.15 we have

$$\begin{aligned} &\int_0^s T_{U_r}^D(A \circ d\omega_r^{\mathcal{H}}, Av(r) + Vv(r)) \\ &= \int_0^s T_{U_r}(A \circ d\omega_r^{\mathcal{H}}, Av(r)) - \int_0^s J_{Vv(r)}(A \circ d\omega_r^{\mathcal{H}})_{U_r}, \end{aligned}$$

where $J = D - \nabla$. As a consequence, $p_v(\omega^{\mathcal{H}})_s$ is actually a horizontal process, that is, it is \mathbb{R}^n -valued.

We can rewrite $p_v(\omega^{\mathcal{H}})_s$ by using Itô's integral, and we obtain

$$\begin{aligned} p_v(\omega^{\mathcal{H}})_s &= v_{\mathcal{H}}(s) + \frac{1}{2} \int_0^s (\mathfrak{Ric}_{\mathcal{H}}^D)_{U_r} (Av(r) + Vv(r)) dr \\ &\quad + \int_0^s J_{Vv(r)}(A \circ d\omega_r^{\mathcal{H}})_{U_r} \\ &\quad - \int_0^s \left(\int_0^r \Omega_{U_\tau}^D(A \circ d\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) d\omega_r^{\mathcal{H}}, \end{aligned}$$

where $\mathfrak{Ric}_{\mathcal{H}}^D$ is the horizontal Ricci curvature of the connection D . We can further simplify this expression as follows.

$$\begin{aligned} &J_{Vv(s)}(A \circ d\omega_s^{\mathcal{H}})_{U_s} = \\ &\sum_{i=1}^n J_{Vv(s)}(A_i)_{U_s} \circ d\omega_s^i = \\ &J_{Vv(s)}(A(d\omega_s^{\mathcal{H}}))_{U_s} + \frac{1}{2} \sum_{i=1}^n A_i J_{Vv(s)}(A_i)_{U_s} ds \\ (4.5) \quad &- \frac{1}{2} \sum_{i=1}^n J_{T(A_i, Av_{\mathcal{H}}(s))}(A_i)_{U_s} ds \end{aligned}$$

As a result, we see that

$$\begin{aligned} p_v(\omega^{\mathcal{H}})_s &= v_{\mathcal{H}}(s) + \frac{1}{2} \sum_{i=1}^n \int_0^s A_i J_{Vv(r)}(A_i)_{U_r} dr \\ &\quad - \frac{1}{2} \sum_{i=1}^n \int_0^s J_{T(A_i, Av_{\mathcal{H}}(r))}(A_i)_{U_r} dr + \frac{1}{2} \int_0^s (\mathfrak{Ric}_{\mathcal{H}}^D)_{U_r} (Av(r) + Vv(r)) dr \\ &\quad + \int_0^s J_{Vv(r)}(A(d\omega_r^{\mathcal{H}}))_{U_r} - \int_0^s \left(\int_0^r \Omega_{U_\tau}^D(A \circ d\omega_\tau^{\mathcal{H}}, Av(\tau) + Vv(\tau)) \right) d\omega_r^{\mathcal{H}}. \end{aligned}$$

More concisely, one can thus write

$$(4.6) \quad p_v(\omega^{\mathcal{H}})_s = \int_0^s q_v(\omega^{\mathcal{H}})_r d\omega_r^{\mathcal{H}} + \int_0^s r_v(\omega^{\mathcal{H}})_r dr,$$

where q_v is a $\mathfrak{so}(n)$ -valued adapted process and r_v is an \mathbb{R}^n -valued adapted process such that $\int_0^s |r_v(u)|_{\mathbb{R}^n}^2 du < \infty$ a.s. The process p_v is therefore an adapted vector field on $W_0(\mathbb{R}^n)$ in the sense of [16, Definition 3.2].

4.3.3. First type of variation. We are now ready to construct the first relevant variation of the horizontal Brownian motion paths. The idea is to use the formula for the deterministic variation given by 3.5 to infer a formula for a convenient stochastic variation.

Notation 4.17. For any $h \in \mathcal{CM}_{\mathcal{H}}(\mathbb{R}^{n+m})$ and any $t \in \mathbb{R}$, we denote by $\rho_t^h : W_0(\mathbb{R}^n) \rightarrow W_0(\mathbb{R}^n)$ a map which is defined $\mu_{\mathcal{H}}$ -a.s. as follows

$$(4.7) \quad (\rho_t^h \omega_{\mathcal{H}})_s = \int_0^s e^{tq_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_u} d\omega_u^{\mathcal{H}} + t \int_0^s r_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_u du.$$

Remark 4.18. As in the deterministic case, observe that ρ^h is **not** the flow generated by p_{τ_h} on $W_0(\mathbb{R}^n)$. This variation is similar to [19, Theorem 7.28]. Let us however observe that $\mu_{\mathcal{H}}$ a.s., $\rho_0^h \omega^{\mathcal{H}} = \omega^{\mathcal{H}}$ and that from (4.6) one has

$$\frac{d}{dt} \Big|_{t=0} (\rho_t^h \omega^{\mathcal{H}})_s = p_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s.$$

We also note that for every $t \in \mathbb{R}$, $e^{tq_{\tau_h}}$ is an $\mathfrak{so}(n)$ -valued process so that for every $t \in \mathbb{R}$ the semimartingale $\left\{ \int_0^s e^{tq_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_u} d\omega_u^{\mathcal{H}} \right\}_{s \in [0,1]}$ is a horizontal Brownian motion with respect to $\mu_{\mathcal{H}}$.

One has then the following analogue of [19, Theorem 7.28] (see [16] for the details) which describes the differential of the horizontal stochastic development map and proves quasi-invariance of the horizontal Wiener measure for the variation described in (4.7).

Theorem 4.19 (Quasi-invariance I). Suppose $h \in \mathcal{CM}_{\mathcal{H}}(\mathbb{R}^{n+m})$.

(1) For every $t \in \mathbb{R}$ the law of the semimartingale $\{(\rho_t^h \omega_{\mathcal{H}})_s\}_{0 \leq s \leq 1}$ (under $\mu_{\mathcal{H}}$) is equivalent to $\mu_{\mathcal{H}}$, and the corresponding Radon-Nikodym density is given by

$$\begin{aligned} \frac{d(\rho_t^h)_* \mu_{\mathcal{H}}}{d\mu_{\mathcal{H}}}(\omega^{\mathcal{H}}) &= \exp \left(t \int_0^1 \left\langle r_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s, e^{tq_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s} d\omega_s^{\mathcal{H}} \right\rangle \right. \\ &\quad \left. - \frac{t^2}{2} \int_0^1 |r_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s|_{\mathbb{R}^n}^2 ds \right). \end{aligned}$$

(2) For every $t \in \mathbb{R}$ the law of the semimartingale $\{\phi_{\mathcal{H}}(\rho_t^h \omega_{\mathcal{H}})_s\}_{0 \leq s \leq 1}$ (under $\mu_{\mathcal{H}}$) is equivalent to μ_W and the corresponding Radon-Nikodym density is given by

$$\begin{aligned} \frac{d(\phi_{\mathcal{H}} \rho_t^h \phi_{\mathcal{H}}^{-1})_* \mu_W}{d\mu_W}(w) &= \exp \left(t \int_0^1 \left\langle r_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s, e^{tq_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s} d\omega_s^{\mathcal{H}} \right\rangle \right. \\ &\quad \left. - \frac{t^2}{2} \int_0^1 |r_{\tau_h(\omega^{\mathcal{H}})}(\omega^{\mathcal{H}})_s|_{\mathbb{R}^n}^2 ds \right), \end{aligned}$$

where $\omega^{\mathcal{H}} = \phi_{\mathcal{H}}^{-1}(w)$.

(3) There exists a version of $\phi_{\mathcal{H}}((\rho_t^h \omega_{\mathcal{H}}))_s$ which is continuous in (s, t) , differentiable in t , and such that

$$\frac{d}{dt} \Big|_{t=0} \phi_{\mathcal{H}}((\rho_t^h \omega_{\mathcal{H}}))_s = U_s \tau_h(\omega^{\mathcal{H}}) \quad \mu_{\mathcal{H}} - a.s.$$

Proof. The first part follows from Girsanov's theorem in the form of [17, Lemma 8.2]. The second part follows from Proposition 3.20 and is similar to [19, Theorem 7.28]. \square

4.3.4. *Second type of variation.* We now turn to the discussion of the stochastic flow generated by p_{τ^h} .

Notation 4.20. For a fixed $h \in \mathcal{CM}_{\mathcal{H}}(\mathbb{R}^{n+m})$ we denote by $SM_{\mathcal{H}}(h)$ the space of continuous and \mathcal{B}_s -adapted \mathbb{R}^n -valued semimartingales $\{z_s\}_{0 \leq s \leq 1}$ that can be written as

$$z_s = \int_0^s a_r dr + \int_0^s \sigma_r d\omega_r^{\mathcal{H}}, \quad 0 \leq s \leq 1,$$

where a is an \mathbb{R}^n -valued \mathcal{B}_s -adapted process such that there exists a deterministic constant C

$$(4.8) \quad |a_s|_{\mathbb{R}^n} \leq C(1 + |h'(s)|_{\mathbb{R}^n}),$$

and where σ is a \mathcal{B}_s -adapted process taking values in the space of isometries of \mathbb{R}^n .

Observe that by Girsanov's theorem in the form of [17, Lemma 8.2], the law of $z \in SM_{\mathcal{H}}(h)$ is equivalent to the law $\mu_{\mathcal{H}}$ of the horizontal Brownian motion. We are now in position to prove that p_{τ^h} generates a flow on the horizontal path space for which the horizontal Wiener measure on \mathbb{R}^{n+m} is quasi-invariant. The following statement is similar to [34, Theorem 3.1]. The proof of that theorem relied on the Picard iteration to find a solution in a space of \mathbb{R}^{n+m} -valued continuous semimartingales equipped with a suitable norm. In our setting the proof is almost identical, so we omit it for conciseness.

Theorem 4.21. For any $h \in \mathcal{CM}_{\mathcal{H}}(\mathbb{R}^{n+m})$ there exists a unique family of semimartingales $\{\nu_t^h, t \in \mathbb{R}\}$ such that

- $\nu_t^h \in SM_{\mathcal{H}}(h)$ for all $t \in \mathbb{R}$ and $\nu_0^h \omega^{\mathcal{H}} = \omega^{\mathcal{H}}$, $\mu_{\mathcal{H}}$ a.s.; hence the law of ν_t^h is equivalent to $\mu_{\mathcal{H}}$;
- For $\mu_{\mathcal{H}}$ -almost every $\omega^{\mathcal{H}}$, the function $t \mapsto \nu_t^h \omega^{\mathcal{H}}$ is a $W_0(\mathbb{R}^n)$ -valued continuous function;
- $\mu_{\mathcal{H}}$ -almost surely, $\nu_{t_1}^h \circ \nu_{t_2}^h(\omega^{\mathcal{H}}) = \nu_{t_1+t_2}^h(\omega^{\mathcal{H}})$, for every $(t_1, t_2) \in \mathbb{R} \times \mathbb{R}$;
- There exists a continuous version of $\{p_{\tau^h} \nu_t^h(\nu_t^h), t \in \mathbb{R}\}$ such that $\mu_{\mathcal{H}}$ -almost surely, $\{\nu_t^h, t \in \mathbb{R}\}$ satisfies the equation

$$(4.9) \quad \nu_t^h(\omega^{\mathcal{H}}) = \omega^{\mathcal{H}} + \int_0^t p_{\tau^h}(\nu_s^h(\omega^{\mathcal{H}}))(\nu_s^h(\omega^{\mathcal{H}})) ds.$$

Remark 4.22. In the previous theorem, the word unique is understood in the sense of [34, Proposition 3.3], that is, in the space $SM_{\mathcal{H}}(h)$.

We are now in position to prove quasi-invariance properties for the horizontal Wiener measure with respect to a suitable flow. The following statement is similar to Theorem 4.1 in [34]. We recall that the horizontal stochastic development $\phi_{\mathcal{H}}$ and its inverse $\phi_{\mathcal{H}}^{-1}$ are defined in Definition 4.10.

Theorem 4.23 (Quasi-invariance II). Let $h \in \mathcal{CM}_{\mathcal{H}}(\mathbb{R}^{n+m})$. The flow $\zeta_t^h = \phi_{\mathcal{H}} \circ \nu_t^h \circ \phi_{\mathcal{H}}^{-1} : W_{x_0}(\mathbb{M}) \rightarrow W_{x_0}(\mathbb{M})$, $t \in \mathbb{R}$, is defined μ_W -a.s. with the generator $U\tau^h\phi_{\mathcal{H}}^{-1}$, and for every $t \in \mathbb{R}$ the distribution of ζ_t^h under μ_W is equivalent to μ_W . More precisely, there exists a family of measurable maps

$$\zeta_t^h : W_{x_0}(\mathbb{M}) \rightarrow W_{x_0}(\mathbb{M}), \quad t \in \mathbb{R},$$

with the following properties.

- For every fixed $t \in \mathbb{R}$, the law $\mu_{\zeta_t^h}$ of ζ_t^h is equivalent to the horizontal Wiener measure μ_W and the Radon-Nikodym derivative is given by

$$\frac{d\mu_{\zeta_t^h}}{d\mu_X}(w) = \frac{d\mu_{\nu_t^h}}{d\mu_{\mathcal{H}}}(\phi_{\mathcal{H}}^{-1}w), \quad w \in W_{x_0}(\mathbb{M}).$$

- For μ_W -almost every $w \in W_{x_0}(\mathbb{M})$, the function $t \mapsto \zeta_t^h w$ is a $W_{x_0}(\mathbb{M})$ -valued continuous differentiable function;
- For μ_W -almost every $w \in W_{x_0}(\mathbb{M})$, there is a continuous version of $t \mapsto U_t\tau^h(\phi_{\mathcal{H}}^{-1}\zeta_t^h w)$ such that $\zeta_v^t w$ satisfies the differential equation

$$\frac{d\zeta_t^h w}{dt} = U_t\tau^h(\phi_{\mathcal{H}}^{-1}\zeta_t^h w);$$

- μ_W -almost surely,

$$\zeta_{t_1}^h \circ \zeta_{t_2}^h = \zeta_{t_1+t_2}^h, \quad \text{for all } (t_1, t_2) \in \mathbb{R} \times \mathbb{R}.$$

Proof. The result follows from Theorem 4.21. For details, we refer to the proof of [34, Theorem 4.1]. \square

4.3.5. *Towards the integration by parts formulas.* It is well known that a quasi-invariance result yields an integration by parts formula on the path space of the underlying diffusion, see B. Driver [17] and then E. Hsu [34] (see also [14, 15, 25]). Integration by parts formulas will be studied in more detail in the second part of the paper, so we only briefly comment on the immediate corollary of Theorem 4.19, which will be proved in another way (see Lemma 4.24) and then extended to cylinder functions. It is obtained from Theorem 4.19 by taking the Bott connection ∇ as the connection D , and following the arguments of the proof in [19, Theorem 7.32].

Lemma 4.24. Let $h \in \mathcal{CM}_{\mathcal{H}}(\mathbb{R}^{n+m})$, then for $f \in C^\infty(\mathbb{M})$,

$$\begin{aligned} & \mathbb{E}(\langle df(W_1), U_1\tau_h(\omega^{\mathcal{H}}) \rangle) \\ &= \mathbb{E}\left(f(W_1) \int_0^1 \left\langle h'(s) + \frac{1}{2}(\mathfrak{Ric}_{\mathcal{H}})_{U_s} h(s), d\omega_s^{\mathcal{H}} \right\rangle_{\mathbb{R}^n}\right), \end{aligned}$$

where \mathbb{E} is the expectation with respect to $\mu_{\mathcal{H}}$ and $\mathfrak{Ric}_{\mathcal{H}}$ is the horizontal Ricci curvature of the Bott connection (viewed as an operator on \mathbb{R}^n).

4.3.6. *The case of a Riemannian submersion: examples.* To finish this part of the paper, we discuss the case when the foliation on \mathbb{M} comes from a totally geodesic submersion $\pi : (\mathbb{M}, g) \rightarrow (\mathbb{B}, j)$ as described in Example 2.1. This should allow the reader to relate our quasi-invariance result to the Riemannian result by B. Driver in [17].

In the submersion case, the notion of horizontal lift of curves plays an important role.

Definition 4.25. Let $\bar{\gamma} : [0, \infty) \rightarrow \mathbb{B}$ be a C^1 -curve. Let $x \in \mathbb{M}$, such that $\pi(x) = \bar{\gamma}(0)$. Then, there exists a unique C^1 -horizontal curve $\gamma : [0, \infty) \rightarrow \mathbb{M}$ such that $\gamma(0) = x$ and $\pi(\gamma(t)) = \bar{\gamma}(t)$. The curve γ is called the horizontal lift of $\bar{\gamma}$ at x .

The notion of horizontal lift may be extended to Brownian motion paths on \mathbb{B} by using stochastic calculus. The argument is similar to the case of the stochastic lift of the Brownian motion of a Riemannian manifold to the orthonormal frame bundle, see for instance [17, Theorem 3.2]).

The submersion has totally geodesic fibers, therefore π is harmonic and the projected process

$$W_t^{\mathbb{B}} = \pi(W_t)$$

is, under $\mu_{\mathcal{H}}$, a Riemannian Brownian motion on \mathbb{B} started at $\pi(x_0)$. The submersion π induces a map $W_{x_0}(\mathbb{M}) \rightarrow W_{\pi(x_0)}(\mathbb{B})$ that we still denote by π . Let now h be a Cameron-Martin path in \mathbb{R}^n and consider the flow $\zeta_t^h : W_{x_0}(\mathbb{M}) \rightarrow W_{x_0}(\mathbb{M})$, $t \in \mathbb{R}$, which is defined μ_W -a.s. according to Theorem 4.23. By using the horizontal stochastic lift $W_{\pi(x_0)}(\mathbb{B}) \rightarrow W_x(\mathbb{M})$, one can construct a flow $\tilde{\zeta}_t^h : W_{\pi(x_0)}(\mathbb{B}) \rightarrow W_{\pi(x_0)}(\mathbb{B})$, $t \in \mathbb{R}$ which is unique $\mu_{W^{\mathbb{B}}}$ -a.s. as mentioned in Remark 4.22. Then we have the following commutative diagram

$$(4.10) \quad \begin{array}{ccc} W_{x_0}(\mathbb{M}) & \xrightarrow{\zeta_t^h} & W_{x_0}(\mathbb{M}) \\ \pi \downarrow & & \downarrow \pi \\ W_{\pi(x_0)}(\mathbb{B}) & \xrightarrow{\tilde{\zeta}_t^h} & W_{\pi(x_0)}(\mathbb{B}) \end{array}$$

By Theorem 4.23, the law of $W^{\mathbb{B}}$ is quasi-invariant under the flow $\tilde{\zeta}_t^h$. Note that the connection D projects down to the Levi-Civita connection on \mathbb{B} , therefore the flow $\tilde{\zeta}_t^h$ provides a version of the flow considered by E. Hsu in [34, Theorem 4.1]. Thus we recover Driver's quasi-invariance result [17] on the manifold \mathbb{B} . Further details on this example will be given in Section 5.3.1, where the generator $\tilde{\zeta}_t^h$ will be computed explicitly.

It may be useful to illustrate the diagram (4.10) in a very simple situation. Recall that the Heisenberg group is the set

$$\mathbb{H}^{2n+1} = \{(x, y, z), x \in \mathbb{R}^n, y \in \mathbb{R}^n, z \in \mathbb{R}\}$$

endowed with the group law

$$(x_1, y_1, z_1) \star (x_2, y_2, z_2) := (x_1 + x_2, y_1 + y_2, z_1 + z_2 + \langle x_1, y_2 \rangle_{\mathbb{R}^n} - \langle x_2, y_1 \rangle_{\mathbb{R}^n}).$$

The vector fields

$$\begin{aligned} X_i &= \frac{\partial}{\partial x_i} - y_i \frac{\partial}{\partial z}, 1 \leq i \leq n, \\ Y_i &= \frac{\partial}{\partial y_i} + x_i \frac{\partial}{\partial z}, 1 \leq i \leq n, \\ Z &= \frac{\partial}{\partial z} \end{aligned}$$

form a basis for the space of left-invariant vector fields on \mathbb{H}^{2n+1} . We choose a left-invariant Riemannian metric on \mathbb{H}^{2n+1} in such a way that $\{X_1, \dots, X_n, Y_1, \dots, Y_n, Z\}$ are orthonormal with respect to this metric. Note that these vector fields satisfy the following commutation relations

$$[X_i, Y_j] = 2\delta_{ij}Z, \quad [X_i, Z] = [Y_i, Z] = 0, \quad i = 1, \dots, n.$$

Then, the projection map

$$\begin{aligned} \pi : \mathbb{H}^{2n+1} &\longrightarrow \mathbb{R}^{2n} \\ (x, y, z) &\longmapsto (x, y) \end{aligned}$$

is a Riemannian submersion with totally geodesic fibers. In that example, the Bott connection is trivial: $\nabla X_i = \nabla Y_j = \nabla Z = 0$ and its torsion is given by

$$T(X_i, Y_j) = -2\delta_{ij}Z, \quad T(X_i, Z) = T(Y_i, Z) = 0.$$

Let now $W_0(\mathbb{R}^{2n})$ be the Wiener space of continuous functions $[0, 1] \rightarrow \mathbb{R}^{2n}$ starting at 0. We denote by $(B_t, \beta_t)_{0 \leq t \leq 1}$ the coordinate maps on $W_0(\mathbb{R}^{2n})$ and by $\mu_{\mathcal{H}}$ the Wiener measure on $W_0(\mathbb{R}^{2n})$, so that $(B_t, \beta_t)_{0 \leq t \leq 1}$ is a $2n$ -dimensional Brownian motion under $\mu_{\mathcal{H}}$. By using the submersion π , the Brownian motion $(B_t, \beta_t)_{0 \leq t \leq 1}$ can be horizontally lifted to the horizontal Brownian motion on \mathbb{H}^{2n+1} which is given explicitly by

$$W_t = \left(B_t, \beta_t, \sum_{i=1}^n \int_0^t B_t^i d\beta_t^i - \beta_t^i dB_t^i \right).$$

Let $h = (h_1, h_2)$ be a Cameron-Martin path in \mathbb{R}^{2n} and consider the Cameron-Martin flow $\tilde{\zeta}_t^h : W_0(\mathbb{R}^{2n}) \rightarrow W_0(\mathbb{R}^{2n})$, $t \in \mathbb{R}$, explicitly given by

$$\tilde{\zeta}_t^h(B, \beta) = (B, \beta) + th.$$

One has then a commutative diagram

$$(4.11) \quad \begin{array}{ccc} W_0(\mathbb{H}^{2n+1}) & \xrightarrow{\zeta_t^h} & W_0(\mathbb{H}^{2n+1}) \\ \pi \downarrow & & \downarrow \pi \\ W_0(\mathbb{R}^{2n}) & \xrightarrow{\tilde{\zeta}_t^h} & W_0(\mathbb{R}^{2n}) \end{array}$$

where ζ_t^h is the flow on $W_0(\mathbb{H}^{2n+1})$ defined $\mu_{\mathcal{H}}$ -a. s. by

$$\begin{aligned} \zeta_t^h(W) &= (B + th_1, \beta + th_2, \\ &\quad \sum_{i=1}^n \int_0^{\cdot} (B_u^i + th_1^i(u)) d(\beta_u^i + th_2^i(u)) - (\beta_u^i + th_2^i(u)) d(B_u^i + th_1^i(u))) \end{aligned}$$

One can compute the generator of this flow as

$$\begin{aligned} &\frac{d}{dt} \Big|_{t=0} \zeta_t^h(W) \\ &= \left(h_1, h_2, \sum_{i=1}^n \int_0^{\cdot} h_1^i(u) d\beta_u^i - h_2^i(u) dB_u^i + \sum_{i=1}^n \int_0^{\cdot} B_u^i dh_2^i(u) - \beta_u^i dh_1^i(u) \right) \\ &= \left(h_1, h_2, \sum_{i=1}^n B^i h_2^i - \beta^i h_1^i + 2 \sum_{i=1}^n \int_0^{\cdot} h_1^i(u) d\beta_u^i - h_2^i(u) dB_u^i \right) \\ &= \sum_{i=1}^n h_1^i X_i(W) + \sum_{i=1}^n h_2^i Y_i(W) + 2 \left(\sum_{i=1}^n \int_0^{\cdot} h_1^i(u) d\beta_u^i - h_2^i(u) dB_u^i \right) Z(W) \end{aligned}$$

As expected, we can interpret this generator in terms of the Bott connection as a straightforward computation shows that

$$\begin{aligned} &\int_0^s T \left(\sum_{i=1}^n X_i \circ dB_u^i + \sum_{i=1}^n Y_i \circ d\beta_u^i, \sum_{i=1}^n h_1^i(u) X_i + \sum_{i=1}^n h_2^i(u) Y_i \right) \\ &= \left(2 \sum_{i=1}^n \int_0^s h_1^i(u) d\beta_u^i - 2 \sum_{i=1}^n \int_0^s h_2^i(u) dB_u^i \right) Z(W) \end{aligned}$$

Therefore, we showed that

$$\begin{aligned} &\frac{d}{dt} \Big|_{t=0} \zeta_t^h(W) \\ &= \sum_{i=1}^n h_1^i X_i(W) + \sum_{i=1}^n h_2^i Y_i(W) \\ &\quad + \int_0^{\cdot} T \left(\sum_{i=1}^n X_i \circ dB_u^i + \sum_{i=1}^n Y_i \circ d\beta_u^i, \sum_{i=1}^n h_1^i(u) X_i + \sum_{i=1}^n h_2^i(u) Y_i \right) \end{aligned}$$

This is exactly Equation (4.4) written in the parallel frame $\{X_i, Y_j, Z\}_{i=1}^n$.

5. INTEGRATION BY PARTS FORMULAS

The goal of this part of the paper is to establish several types of integration by parts formulas for the horizontal Brownian motion. This part relies on very different techniques than the ones used in the first part and therefore we need to introduce more notation. Though we will consider the horizontal Brownian motion constructed from the frame bundle, in this part of the paper we will rely on the stochastic parallel transport rather than the stochastic lift to the frame bundle (although these are of course equivalent). Also, instead of working with general connections denoted by D in Section 4, we now consider connections satisfying Assumption 1 and with the additional property that the torsion satisfies B. Driver's anti-symmetry condition [17, Definition 8.1]. Throughout this part, we will work with the following probability space.

Notation 5.1. By $(\Omega, \mathcal{B}, \mu_{\mathcal{H}})$ we denote the probability space, where $\Omega := W_0(\mathbb{R}^n)$ is the space of continuous functions $\omega^{\mathcal{H}} : [0, 1] \rightarrow \mathbb{R}^n$ such that $\omega^{\mathcal{H}}(0) = 0$, \mathcal{B} is the Borel σ -field on $W_0(\mathbb{R}^n)$, and $\mu_{\mathcal{H}}$ is the Wiener measure on Ω . Then the coordinate process $\{\omega_s^{\mathcal{H}}\}_{0 \leq s \leq 1}$ is a Brownian motion with values in \mathbb{R}^n . The usual completion of the natural filtration generated by $\{\omega_s^{\mathcal{H}}\}_{0 \leq s \leq 1}$ will be denoted by \mathcal{F}_s .

Recall that for $x \in \mathbb{M}$ the horizontal Brownian motion on \mathbb{M} started at x is defined as $W_s = \pi(U_s)$, where U_s is a solution to the Stratonovich stochastic differential equation (4.2) with $U_0 = u_0 \in \mathcal{O}_{\mathcal{H}}(\mathbb{M})$ such that $\pi(u_0) = x$.

5.1. Horizontal Weitzenböck type formulas. We start by introducing a family of connections that will be of interest to us later, and we review some known results on the Weitzenböck formulas proved previously in [7].

5.1.1. Generalized Levi-Civita connections and adjoint connections. In Section 5.1.2 we aim at studying Weitzenböck-type identities for the horizontal Laplacian, and for this we need to introduce a new class of connections. The main reason why we use these connections is that we can not make use of the Bott connection since the adjoint connection to the Bott connection is not metric. We refer to [7, 22, 30, 31] and especially the books [23, 24] for a discussion on Weitzenböck-type identities and adjoint connections. Instead we make use of the family of connections first introduced in [2] and only keep the Bott connection as a reference connection.

This family of connections is constructed from a natural variation of the metric that we recall now. The Riemannian metric g can be split using horizontal and vertical subbundles described in Section 2.2

$$(5.1) \quad g = g_{\mathcal{H}} \oplus g_{\mathcal{V}}.$$

Using the splitting of the Riemannian metric g in (5.1) we can introduce the following one-parameter family of Riemannian metrics

$$g_\varepsilon = g_{\mathcal{H}} \oplus \frac{1}{\varepsilon} g_{\mathcal{V}}, \quad \varepsilon > 0.$$

One can check that for every $\varepsilon > 0$, $\nabla g_\varepsilon = 0$ where ∇ is the Bott connection. The metric g_ε then induces a metric on the cotangent bundle which we still denote by g_ε , and therefore

$$\|\eta\|_\varepsilon^2 = \|\eta\|_{\mathcal{H}}^2 + \varepsilon \|\eta\|_{\mathcal{V}}^2, \quad \text{for every } \eta \in T_x^* \mathbb{M}.$$

For each $Z \in \Gamma^\infty(\mathcal{V})$ there is a unique skew-symmetric endomorphism $J_Z : \mathcal{H}_x \rightarrow \mathcal{H}_x$, $x \in \mathbb{M}$ such that for all horizontal vector fields $X, Y \in \mathcal{H}_x$

$$(5.2) \quad g_{\mathcal{H}}(J_Z(X), Y)_x = g_{\mathcal{V}}(Z, T(X, Y))_x,$$

where T is the torsion tensor of ∇ . We then extend J_Z to be 0 on \mathcal{V}_x . Also, to ensure (5.2) holds also for $Z \in \Gamma^\infty(\mathcal{H})$, taking into account (2.1) we set $J_Z \equiv 0$.

Following [2] we introduce the following family of connections

$$\nabla_X^\varepsilon Y = \nabla_X Y - T(X, Y) + \frac{1}{\varepsilon} J_Y X, \quad X, Y \in \Gamma^\infty(\mathbb{M}).$$

It is easy to check that $\nabla^\varepsilon g_\varepsilon = 0$ and the torsion of ∇^ε is given by

$$T^\varepsilon(X, Y) = -T(X, Y) + \frac{1}{\varepsilon} J_Y X - \frac{1}{\varepsilon} J_X Y, \quad X, Y \in \Gamma^\infty(\mathbb{M}).$$

The adjoint connection to ∇^ε as described by B. Driver in [17], see also [23, Section 1.3] for a discussion about adjoint connections, is then given by

$$(5.3) \quad \widehat{\nabla}_X^\varepsilon Y := \nabla_X^\varepsilon Y - T^\varepsilon(X, Y) = \nabla_X Y + \frac{1}{\varepsilon} J_X Y,$$

thus $\widehat{\nabla}^\varepsilon$ is also a metric connection. Moreover, it preserves the horizontal and vertical bundles.

Remark 5.2. Note that the connection $\widehat{\nabla}^\varepsilon$ therefore satisfies Assumption 1 for every $\varepsilon > 0$.

For later use, we record that the torsion of $\widehat{\nabla}^\varepsilon$ is

$$(5.4) \quad \widehat{T}^\varepsilon(X, Y) = -T^\varepsilon(X, Y) = T(X, Y) - \frac{1}{\varepsilon} J_Y X + \frac{1}{\varepsilon} J_X Y.$$

The Riemannian curvature tensor of $\widehat{\nabla}^\varepsilon$ can be computed explicitly in terms of the Riemannian curvature tensor R of the Bott connection ∇ and it is given by the following lemma.

Lemma 5.3. For $X, Y, Z \in \Gamma^\infty(\mathbb{M})$

$$\begin{aligned} \widehat{R}^\varepsilon(X, Y)Z &= R(X, Y)Z + \frac{1}{\varepsilon} J_{T(X, Y)}Z + \frac{1}{\varepsilon^2} (J_X J_Y - J_Y J_X)Z + \\ &\quad \frac{1}{\varepsilon} (\nabla_X J)_Y Z - \frac{1}{\varepsilon} (\nabla_Y J)_X Z, \end{aligned}$$

where R is the curvature tensor of the Bott connection.

Proof.

$$\begin{aligned}
 \widehat{R}^\varepsilon(X, Y)Z &= \widehat{\nabla}_X^\varepsilon \widehat{\nabla}_Y^\varepsilon Z - \widehat{\nabla}_Y^\varepsilon \widehat{\nabla}_X^\varepsilon Z - \widehat{\nabla}_{[X, Y]}^\varepsilon Z \\
 &= (\nabla_X \nabla_Y + \frac{1}{\varepsilon} (\nabla_X J)_Y + \frac{1}{\varepsilon} J_X \nabla_Y + \frac{1}{\varepsilon} J_Y \nabla_X + \frac{1}{\varepsilon} J_{\nabla_X Y} + \frac{1}{\varepsilon^2} J_X J_Y)Z \\
 &\quad - (\nabla_Y \nabla_X + \frac{1}{\varepsilon} (\nabla_Y J)_X + \frac{1}{\varepsilon} J_Y \nabla_X + \frac{1}{\varepsilon} J_{\nabla_Y X} + \frac{1}{\varepsilon} J_X \nabla_Y + \frac{1}{\varepsilon^2} J_Y J_X)Z \\
 &\quad - \nabla_{[X, Y]} Z - \frac{1}{\varepsilon} J_{[X, Y]} Z \\
 &= R(X, Y)Z + \frac{1}{\varepsilon^2} (J_X J_Y - J_Y J_X)Z + \\
 &\quad \frac{1}{\varepsilon} (\nabla_X J)_Y Z - \frac{1}{\varepsilon} (\nabla_Y J)_X Z + \frac{1}{\varepsilon} J_{T(X, Y)} Z.
 \end{aligned}$$

□

We define the horizontal Ricci curvature $\mathfrak{Ric}_{\mathcal{H}}$ for the Bott connection as the fiberwise symmetric linear map on one-forms such that for all smooth functions f, g on \mathbb{M}

$$\langle \mathfrak{Ric}_{\mathcal{H}}(df), dg \rangle = \mathbf{Ric}(\nabla_{\mathcal{H}} f, \nabla_{\mathcal{H}} g) = \mathbf{Ric}_{\mathcal{H}}(\nabla f, \nabla g),$$

where \mathbf{Ric} is the Ricci curvature of the Bott connection ∇ and $\mathbf{Ric}_{\mathcal{H}}$ is its horizontal Ricci curvature (horizontal trace of the full curvature tensor R of the Bott connection). The fact that $\mathbf{Ric}_{\mathcal{H}}$ is symmetric follows from [33, Lemma 4.2].

5.1.2. *Weitzenböck formulas.* A key ingredient in studying the horizontal Brownian motion is the Weitzenböck formula that has been proven in [2, 7]. We recall here this formula. If Z_1, \dots, Z_m is a local vertical frame, then the $(1, 1)$ tensor

$$\mathbf{J}^2 := \sum_{\ell=1}^m J_{Z_\ell} J_{Z_\ell}$$

does not depend on the choice of the frame and may be defined globally.

Example 5.1 (Example 2.2 revisited). If \mathbb{M} is a K-contact manifold equipped with the Reeb foliation, then, by taking Z to be the Reeb vector field, one gets $\mathbf{J}^2 = J_Z^2 = -\mathbf{Id}_{\mathcal{H}}$.

The *horizontal divergence* of the torsion T is the $(1, 1)$ tensor which in a local horizontal frame X_1, \dots, X_n is defined by

$$(5.5) \quad \delta_{\mathcal{H}} T(X) := - \sum_{j=1}^n (\nabla_{X_j} T)(X_j, X).$$

By using the duality between the tangent and cotangent bundles with respect to the metric g , we can identify the $(1, 1)$ tensors \mathbf{J}^2 and $\delta_{\mathcal{H}} T$ with linear maps on the cotangent bundle $T^*\mathbb{M}$.

Namely, let $\sharp : T^*\mathbb{M} \rightarrow T\mathbb{M}$ be the standard musical (raising an index) isomorphism which is defined as the unique vector ω^\sharp such that for any $x \in \mathbb{M}$

$$g\left(\omega^\sharp, X\right)_x = \omega(X) \text{ for all } X \in T_x\mathbb{M},$$

while in local coordinates the isomorphism \sharp can be written as follows

$$\omega = \sum_{i=1}^{n+m} \omega_i dx^i \mapsto \omega^\sharp = \sum_{j=1}^{n+m} \omega^j \partial_j = \sum_{j=1}^{n+m} \sum_{i=1}^{n+m} g^{ij} \omega_i \partial_j.$$

The inverse of this isomorphism is the (lowering an index) isomorphism $\flat : T\mathbb{M} \rightarrow T^*\mathbb{M}$ defined by

$$X^\flat = g(X, \cdot)_x, X \in T_x\mathbb{M}$$

and in local coordinates

$$X = \sum_{i=1}^{n+m} X^i \partial_i \mapsto X^\flat = \sum_{i=1}^{n+m} X_i dx^i = \sum_{i=1}^{n+m} \sum_{j=1}^{n+m} g_{ij} X^j dx^i.$$

If η is a one-form, we define the horizontal gradient in a local adapted frame of η as the $(0, 2)$ tensor

$$\nabla_{\mathcal{H}}\eta = \sum_{i=1}^n \nabla_{X_i}\eta \otimes \theta_i,$$

where $\theta_i, i = 1, \dots, n$ is the dual to X_i .

Finally, for $\varepsilon > 0$, we consider the following operator which is defined on one-forms by

(5.6)

$$\square_\varepsilon := \sum_{i=1}^n (\nabla_{X_i} - \mathfrak{T}_{X_i}^\varepsilon)^2 - (\nabla_{\nabla_{X_i} X_i} - \mathfrak{T}_{\nabla_{X_i} X_i}^\varepsilon) - \frac{1}{\varepsilon} \mathbf{J}^2 + \frac{1}{\varepsilon} \delta_{\mathcal{H}} T - \mathfrak{Ric}_{\mathcal{H}},$$

where \mathfrak{T}^ε is the $(1, 1)$ tensor defined by

$$\mathfrak{T}_X^\varepsilon Y = -T(X, Y) + \frac{1}{\varepsilon} J_Y X, \quad X, Y \in \Gamma^\infty(\mathbb{M}).$$

Similarly as before, we will use the notation

$$\mathfrak{T}_{\mathcal{H}}^\varepsilon \eta := \sum_{i=1}^n \mathfrak{T}_{X_i}^\varepsilon \eta \otimes \theta_i.$$

The expression in (5.6) does not depend on the choice of the local horizontal frame and thus \square_ε may be globally defined. Formally, we have

$$(5.7) \quad \square_\varepsilon = -(\nabla_{\mathcal{H}} - \mathfrak{T}_{\mathcal{H}}^\varepsilon)^*(\nabla_{\mathcal{H}} - \mathfrak{T}_{\mathcal{H}}^\varepsilon) - \frac{1}{\varepsilon} \mathbf{J}^2 + \frac{1}{\varepsilon} \delta_{\mathcal{H}} T - \mathfrak{Ric}_{\mathcal{H}},$$

where the adjoint is understood with respect to the $L^2(\mathbb{M}, g_\varepsilon, \mu)$ inner product on sections, i.e. $\int_{\mathbb{M}} \langle \cdot, \cdot \rangle_\varepsilon d\mu$ (see [1, Lemma 5.3] for more detail). The

main result in [7] is the following. Here the Laplacian L is defined by Equation (4.1) in Section 4.1.1

Theorem 5.4 (Lemma 3.3, Theorem 3.1 in [7]). Let $f \in C^\infty(\mathbb{M})$, $x \in \mathbb{M}$ and $\varepsilon > 0$, then

$$(5.8) \quad dLf(x) = \square_\varepsilon df(x),$$

where L is defined by Equation (4.1).

Remark 5.5. Using [7, Lemma 3.4], we see that for $\varepsilon_1, \varepsilon_2 > 0$, the operator $\square_{\varepsilon_1} - \square_{\varepsilon_2}$ vanishes on exact one-forms. It is therefore no surprise that the left hand side of (5.8) does not depend of ε .

To conclude this section we remark, and this is not a coincidence, that the potential term in the Weitzenböck identity can be identified with the horizontal Ricci curvature of the adjoint connection $\widehat{\nabla}^\varepsilon$.

Lemma 5.6. The horizontal Ricci curvature of the adjoint connection $\widehat{\nabla}^\varepsilon$ is given by

$$\mathfrak{Ric}_{\mathcal{H}}^\varepsilon = \mathfrak{Ric}_{\mathcal{H}} - \frac{1}{\varepsilon} \delta_{\mathcal{H}}^* T + \frac{1}{\varepsilon} \mathbf{J}^2,$$

where $\delta_{\mathcal{H}}^* T$ denotes the adjoint of $\delta_{\mathcal{H}} T$ with respect to the metric g .

Proof. Let $X, Y \in \Gamma^\infty(TM)$ and X_1, \dots, X_n be a local horizontal orthonormal frame. By the definition of the horizontal Ricci curvature and Lemma 5.3 we have

$$\begin{aligned} & \mathbf{Ric}_{\mathcal{H}}^\varepsilon(X, Y) \\ &= \sum_{i=1}^n g_{\mathcal{H}}(\widehat{R}^\varepsilon(X_i, X)Y, X_i) \\ &= \sum_{i=1}^n g_{\mathcal{H}}(R(X_i, X)Y, X_i) + \sum_{i=1}^n g_{\mathcal{H}}\left(\frac{1}{\varepsilon} J_{T(X_i, X)} Y, X_i\right) \\ & \quad + \sum_{i=1}^n g_{\mathcal{H}}\left(\frac{1}{\varepsilon} (\nabla_{X_i} J)_X Y - \frac{1}{\varepsilon} (\nabla_X J)_{X_i} Y, X_i\right). \end{aligned}$$

For the first term, we have

$$\sum_{i=1}^n g_{\mathcal{H}}(R(X_i, X)Y, X_i) = \mathbf{Ric}_{\mathcal{H}}(X, Y).$$

For the second term, we easily see that

$$\begin{aligned} \sum_{i=1}^n g_{\mathcal{H}}(J_{T(X_i, X)} Y, X_i) &= - \sum_{i=1}^n g_{\mathcal{V}}(T(X, X_i), T(Y, X_i)) \\ &= g_{\mathcal{H}}(\mathbf{J}^2 X, Y). \end{aligned}$$

For the third term, we first observe that $g_{\mathcal{H}}((\nabla_X J)_{X_i} Y, X_i) = 0$. Then, we have

$$\begin{aligned} \sum_{i=1}^n g_{\mathcal{H}}((\nabla_{X_i} J)_X Y, X_i) &= - \sum_{i=1}^n g_{\mathcal{H}}((\nabla_{X_i} J)_X X_i, Y) \\ &= - \sum_{i=1}^n g_{\mathcal{V}}((\nabla_{X_i} T)(X_i, Y), X) \\ &= g_{\mathcal{V}}(\delta_{\mathcal{H}} T(Y), X). \end{aligned}$$

□

5.2. Integration by parts formula on the horizontal path space. We fix $\varepsilon > 0$ throughout the section. Our goal in this section is to prove integration by parts formulas on the path space of the horizontal Brownian motion. Some of the integration by parts formulas for the damped Malliavin derivative have been already announced in a less general and slightly different setting in [4]. The integration by part formulas for the intrinsic Malliavin derivative are new. We point out a significant difference of our techniques from what have been used in [1, 2, 4]. Namely, we shall mostly make use of the adjoint connection $\widehat{\nabla}^\varepsilon$ instead of the Bott connection. Below we summarize important properties of the connection $\widehat{\nabla}^\varepsilon$ which will be used extensively in the sequel.

Remark 5.7 (Properties of the adjoint connection). Let $\widehat{\nabla}^\varepsilon$ be the adjoint connection defined by Equation 5.3. Then it satisfies the following properties.

- The adjoint connection is *metric*, that is, $\widehat{\nabla}^\varepsilon g_\varepsilon = 0$;
- The adjoint connection is *horizontal*, that is, if $X \in \Gamma^\infty(\mathcal{H})$ and $Y \in \Gamma^\infty(\mathbb{M})$ then $\widehat{\nabla}_Y^\varepsilon X \in \Gamma^\infty(\mathcal{H})$;
- The torsion tensor \widehat{T}^ε of $\widehat{\nabla}^\varepsilon$ is *skew-symmetric*, that is, it satisfies B. Driver's skew-symmetry condition (see [17, Definition 8.1]) as follows. For $X, Y, Z \in \Gamma^\infty(\mathbb{M})$

$$\langle \widehat{T}^\varepsilon(X, Y), Z \rangle_\varepsilon = -\langle \widehat{T}^\varepsilon(X, Z), Y \rangle_\varepsilon.$$

The latter can be seen from Equation (5.4)

$$\widehat{T}^\varepsilon(X, Y) = T(X, Y) - \frac{1}{\varepsilon} J_Y X + \frac{1}{\varepsilon} J_X Y$$

and the definition of J . In addition to [17] we refer to [23] for more details on adjoint connections and connections with skew-symmetric torsion, including examples.

Next recall that a stochastic parallel transport on forms can be defined following [38, p. 50].

Notation 5.8. Let $\tilde{\nabla}$ be a general connection on \mathbb{M} , and $\{M_s\}_{0 \leq s \leq 1}$ be a semimartingale on \mathbb{M} . We denote by

$$\tilde{\llbracket}_{0,s} : T_{M_0}\mathbb{M} \rightarrow T_{M_s}\mathbb{M}$$

the *stochastic parallel transport* of vector fields along the paths of $\{M_s\}_{0 \leq s \leq 1}$. Then by duality we can define the stochastic parallel transport on one-forms as follows. We have

$$\tilde{\llbracket}_{0,s}^* : T_{M_s}^*\mathbb{M} \rightarrow T_{M_0}^*\mathbb{M}$$

such that for $\alpha \in T_{M_s}^*\mathbb{M}$

$$(5.9) \quad \langle \tilde{\llbracket}_{0,s}^* \alpha, v \rangle = \langle \alpha, \tilde{\llbracket}_{0,s} v \rangle, \quad v \in T_{M_0}\mathbb{M}.$$

In particular, the stochastic parallel transport for the adjoint connection $\widehat{\nabla}^\varepsilon = \nabla + \frac{1}{\varepsilon}J$ along the paths of the horizontal Brownian motion $\{W_s\}_{0 \leq s \leq 1}$ will be denoted by $\widehat{\Theta}_s^\varepsilon$. Since the adjoint connection $\widehat{\nabla}^\varepsilon$ is horizontal, the map $\widehat{\Theta}_s^\varepsilon : T_x\mathbb{M} \rightarrow T_{W_s}\mathbb{M}$ is an isometry that preserves the horizontal bundle, that is, if $u \in \mathcal{H}_x$, then $\widehat{\Theta}_s^\varepsilon u \in \mathcal{H}_{W_t}$. We see then that the anti-development of $\{W_s\}_{0 \leq s \leq 1}$ defined as

$$B_s := \int_0^s (\widehat{\Theta}_r^\varepsilon)^{-1} \circ dW_r,$$

is a Brownian motion in the horizontal space \mathcal{H}_x .

Remark 5.9. Observe that on one-forms the process $\widehat{\Theta}_s^\varepsilon : T_{W_s}^*\mathbb{M} \rightarrow T_x^*\mathbb{M}$ is a solution to the following covariant Stratonovich stochastic differential equation

$$d[\widehat{\Theta}_s^\varepsilon \alpha(W_s)] = \widehat{\Theta}_s^\varepsilon \widehat{\nabla}_{odW_s}^\varepsilon \alpha(W_s),$$

where α is any smooth one-form. Since $\widehat{\nabla}_{odW_s}^\varepsilon = \nabla_{odW_s} + \frac{1}{\varepsilon}J_{odW_s} = \nabla_{odW_s}$, we deduce that $\widehat{\Theta}^\varepsilon$ is actually independent of ε and is therefore also the stochastic parallel transport for the Bott connection. As a consequence, the Brownian motion $\{B_s\}_{0 \leq s \leq 1}$ and its filtration are also independent of the particular choice of ε .

We define a *damped parallel transport* $\tau_s^\varepsilon : T_{W_s}^*\mathbb{M} \rightarrow T_x^*\mathbb{M}$ by the formula

$$(5.10) \quad \tau_s^\varepsilon = \mathcal{M}_s^\varepsilon \Theta_s^\varepsilon,$$

where the process $\Theta_s^\varepsilon : T_{W_s}^*\mathbb{M} \rightarrow T_x^*\mathbb{M}$ is the stochastic parallel transport of one-forms with respect to the connection $\nabla^\varepsilon = \nabla - \mathfrak{I}^\varepsilon$ along the paths of $\{W_s\}_{0 \leq s \leq 1}$. The multiplicative functional $\mathcal{M}_s^\varepsilon : T_x^*\mathbb{M} \rightarrow T_x^*\mathbb{M}$, $s \geq 0$, is defined as the solution to the following ordinary differential equation

$$(5.11) \quad \begin{aligned} \frac{d\mathcal{M}_s^\varepsilon}{ds} &= -\frac{1}{2}\mathcal{M}_s^\varepsilon \Theta_s^\varepsilon \left(\frac{1}{\varepsilon} \mathbf{J}^2 - \frac{1}{\varepsilon} \delta_{\mathcal{H}} T + \mathfrak{Ric}_{\mathcal{H}} \right) (\Theta_s^\varepsilon)^{-1}, \\ \mathcal{M}_0^\varepsilon &= \mathbf{Id}. \end{aligned}$$

Observe that the process $\tau_s^\varepsilon : T_{W_s}^* \mathbb{M} \rightarrow T_x^* \mathbb{M}$ is a solution of the following covariant Stratonovich stochastic differential equation

$$(5.12) \quad \begin{aligned} d[\tau_s^\varepsilon \alpha(W_s)] &= \tau_s^\varepsilon \left(\nabla_{\circ dW_s} - \mathfrak{I}_{\circ dW_t}^\varepsilon - \frac{1}{2} \left(\frac{1}{\varepsilon} \mathbf{J}^2 - \frac{1}{\varepsilon} \delta_{\mathcal{H}} T + \mathfrak{Ric}_{\mathcal{H}} \right) ds \right) \alpha(W_s), \\ \tau_0 &= \mathbf{Id}, \end{aligned}$$

where α is any smooth one-form.

Also observe that $\mathcal{M}_s^\varepsilon$ is invertible and that its inverse is the solution of the following ordinary differential equation

$$(5.13) \quad \frac{d(\mathcal{M}_s^\varepsilon)^{-1}}{ds} = \frac{1}{2} \Theta_s^\varepsilon \left(\frac{1}{\varepsilon} \mathbf{J}^2 - \frac{1}{\varepsilon} \delta_{\mathcal{H}} T + \mathfrak{Ric}_{\mathcal{H}} \right) (\Theta_s^\varepsilon)^{-1} (\mathcal{M}_s^\varepsilon)^{-1}.$$

In particular, it implies that τ_s^ε is invertible.

5.2.1. Malliavin and directional derivatives. We recall that the horizontal Wiener measure on $W_{x_0}(\mathbb{M})$ is defined as the distribution of the horizontal Brownian motion. The coordinate process on $W_{x_0}(\mathbb{M})$ as before is denoted by $\{w_s\}_{0 \leq s \leq 1}$.

Definition 5.10. A function $F : W_{x_0}(\mathbb{M}) \rightarrow \mathbb{R}$ is called a C^k -cylinder function if there exists a partition

$$\pi := \{0 = s_0 < s_1 < s_2 < \dots < s_n \leq 1\}$$

of the interval $[0, 1]$ and $f \in C^k(\mathbb{M}^n)$ such that

$$(5.14) \quad F(w) = f(w_{s_1}, \dots, w_{s_n}) \text{ for all } w \in W_{x_0}(\mathbb{M}).$$

The function F is called a *smooth cylinder function* on $W_{x_0}(\mathbb{M})$, if there exists a partition π and $f \in C^\infty(\mathbb{M}^n)$ such that (5.14) holds.

We denote by $\mathcal{FC}^k(W_{x_0}(\mathbb{M}))$ the space of C^k -cylinder functions, and by $\mathcal{FC}^\infty(W_{x_0}(\mathbb{M}))$ the space of C^∞ -cylinder functions.

Remark 5.11. Note that the representation (5.14) of a cylinder function is not unique. However, let $F \in \mathcal{FC}^\infty(W_{x_0}(\mathbb{M}))$ and $n \geq 0$ be the minimal n such that there exists a partition

$$\pi := \{0 = s_0 < s_1 < s_2 < \dots < s_n \leq 1\}$$

of the interval $[0, 1]$ and $f \in C^k(\mathbb{M}^n)$ such that

$$(5.15) \quad F(w) = f(w_{s_1}, \dots, w_{s_n}) \text{ for all } w \in W_{x_0}(\mathbb{M}).$$

In that case, if

$$\tilde{\pi} = \{0 = \tilde{s}_0 < \tilde{s}_1 < \tilde{s}_2 < \dots < \tilde{s}_n \leq 1\}$$

is another partition of the interval $[0, 1]$ and $\tilde{f} \in C^k(\mathbb{M}^n)$ is such that

$$F(w) = \tilde{f}(w_{\tilde{s}_1}, \dots, w_{\tilde{s}_n}) \text{ for all } w \in W_x(\mathbb{M}),$$

then $\pi = \tilde{\pi}$ and $f = \tilde{f}$. Indeed, since

$$f(w_{s_1}, \dots, w_{s_n}) = \tilde{f}(w_{\tilde{s}_1}, \dots, w_{\tilde{s}_n})$$

we first deduce that $s_1 = \tilde{s}_1$. Otherwise $d_1 f = 0$ or $d_1 \tilde{f} = 0$, where d_1 denotes the differential with respect to the first component. This contradicts the fact that n is minimal. Similarly, $s_2 = \tilde{s}_2$ and more generally $s_k = \tilde{s}_k$. The representation (5.11) will be referred to as the *minimal representation* of F .

We now turn to the definition of directional derivative on the horizontal path space.

Definition 5.12. Let $F = f(w_{s_1}, \dots, w_{s_n}) \in \mathcal{F}C^\infty(W_x(\mathbb{M}))$. For an \mathcal{F} -adapted and $T_x\mathbb{M}$ -valued semimartingale $(v(s))_{0 \leq s \leq 1}$ such that $v(0) = 0$, we define the directional derivative

$$\mathbf{D}_v F = \sum_{i=1}^n \left\langle d_i f(W_{s_1}, \dots, W_{s_n}), \widehat{\Theta}_{s_i}^\varepsilon v(s_i) \right\rangle$$

Definition 5.13. For $F = f(w_{s_1}, \dots, w_{s_n}) \in \mathcal{F}C^\infty(W_x(\mathbb{M}))$ we define the *damped Malliavin derivative* by

$$\widetilde{D}_s^\varepsilon F := \sum_{i=1}^n \mathbf{1}_{[0, s_i]}(s) (\tau_s^\varepsilon)^{-1} \tau_{s_i}^\varepsilon d_i f(W_{s_1}, \dots, W_{s_n}), \quad 0 \leq s \leq 1.$$

Observe that from this definition $\widetilde{D}_s^\varepsilon F \in T_{W_s}^* \mathbb{M}$.

Remark 5.14. Note that the directional derivative \mathbf{D} is independent of ε , but the damped Malliavin derivative depends on ε . In addition, both the directional derivatives and damped Malliavin derivatives are independent of the representation of F . Indeed, let $F = f(w_{s_1}, \dots, w_{s_n})$ be the minimal representation of F . If $\tilde{f}(w_{\tilde{s}_1}, \dots, w_{\tilde{s}_N})$ is another representation of F , then for every $1 \leq j \leq N$, we have either that there exists i such that $s_i = \tilde{s}_j$ in which case $d_i f = d_j \tilde{f}$, or for all i , $s_i \neq \tilde{s}_j$ in which case $d_j \tilde{f} = 0$.

Before we can formulate the main result, we need to define an analog of the Cameron-Martin subspace.

Definition 5.15. An \mathcal{F}_s -adapted absolutely continuous \mathcal{H}_x -valued process $\{\gamma(s)\}_{0 \leq s \leq 1}$ such that $\gamma(0) = 0$ and $\mathbb{E}_x \left(\int_0^1 \|\gamma'(s)\|_{\mathcal{H}}^2 ds \right) < \infty$ will be called a *horizontal Cameron-Martin process*.

Definition 5.16. Suppose $\{v(s)\}_{0 \leq s \leq 1}$ is an \mathcal{F}_s -adapted $T_x\mathbb{M}$ -valued continuous semimartingale such that $v(0) = 0$ and $\mathbb{E}_x \left(\int_0^1 \|v(s)\|^2 ds \right) < \infty$. We call $\{v(s)\}_{0 \leq s \leq 1}$ a *tangent process* if the process

$$v(s) - \int_0^s (\widehat{\Theta}_r^\varepsilon)^{-1} T(\widehat{\Theta}_s r^\varepsilon \circ dB_r, \widehat{\Theta}_r^\varepsilon v(r))$$

is a horizontal Cameron-Martin process.

Remark 5.17. By Remark 5.9 the stochastic parallel transport $\widehat{\Theta}_s^\varepsilon$ is independent of ε , therefore the notion of a tangent process is itself independent of ε as well.

Remark 5.18. As the torsion T is a vertical tensor, then an \mathcal{F}_s -adapted $T_x\mathbb{M}$ -valued continuous semimartingale $\{v(s)\}_{0 \leq s \leq 1}$ such that

$$\mathbb{E}_x \left(\int_0^1 \|v(s)\|^2 ds \right) < \infty, \quad v(0) = 0$$

is in $TW_{\mathcal{H}}(\mathbb{M})$ if and only if

- (1) The horizontal part $v_{\mathcal{H}}$ is a horizontal Cameron-Martin process;
- (2) The vertical part $v_{\mathcal{V}}$ is given by

$$v_{\mathcal{V}}(s) = \int_0^s (\widehat{\Theta}_r^\varepsilon)^{-1} T(\widehat{\Theta}_r^\varepsilon \circ dB_r, \widehat{\Theta}_r^\varepsilon v_{\mathcal{H}}(r)).$$

The main results of this section are the following two theorems.

Theorem 5.19 (Integration by parts for the damped Malliavin derivative). Suppose $F \in \mathcal{FC}^\infty(W_x(\mathbb{M}))$ and γ is a tangent process, then

$$(5.16) \quad \mathbb{E}_x \left(\int_0^1 \langle \widetilde{D}_s^\varepsilon F, \widehat{\Theta}_s^\varepsilon \gamma'(s) \rangle ds \right) = \mathbb{E}_x \left(F \int_0^1 \langle \gamma'(s), dB_s \rangle_{\mathcal{H}} \right).$$

Theorem 5.20 (Integration by parts for the directional derivatives). Suppose $F \in \mathcal{FC}^\infty(W_x(\mathbb{M}))$ and v is a tangent process, then

$$\mathbb{E}_x(\mathbf{D}_v F) = \mathbb{E}_x \left(F \int_0^1 \left\langle v'_{\mathcal{H}}(s) + \frac{1}{2} (\widehat{\Theta}_s^\varepsilon)^{-1} \mathfrak{Ric}_{\mathcal{H}} \widehat{\Theta}_s^\varepsilon v_{\mathcal{H}}(s), dB_t \right\rangle_{\mathcal{H}} \right).$$

Even though these two integration by parts formulas seem similar, they are quite different in nature. The damped derivative is used to derive gradient bounds and functional inequalities on the path space (e.g. [2, 4]). The directional derivative, however, is more related to quasi-invariance properties such as in Section 4.3, and the expression

$$\int_0^1 \left\langle v'_{\mathcal{H}}(s) + \frac{1}{2} (\widehat{\Theta}_s^\varepsilon)^{-1} \mathfrak{Ric}_{\mathcal{H}} \widehat{\Theta}_s^\varepsilon v_{\mathcal{H}}(s), dB_s \right\rangle_{\mathcal{H}}$$

can be viewed as a horizontal divergence on the path space.

The remainder of the section is devoted to proving Theorem 5.19 and Theorem 5.20. We adapt the techniques from the Markovian stochastic calculus developed by Fang-Malliavin [27] and E. Hsu [35] in the Riemannian case to our setting.

5.2.2. *Gradient formula.* In this preliminary section we recall the gradient formula for the semigroup P_s . In the case the Yang-Mills condition is satisfied, that is, the horizontal divergence $\delta_{\mathcal{H}}T = 0$, the operator \square_ε is essentially self-adjoint on $L^2(\mathbb{M}, g_\varepsilon, \mu)$ equipped with inner product on sections, i.e. $\int_{\mathbb{M}} \langle \cdot, \cdot \rangle_\varepsilon d\mu$, and the gradient representation was first proved in [2].

Lemma 5.21 (Theorem 4.6 and Corollary 4.7 in [2], Theorem 2.7 in [32]). For $f \in C^\infty(\mathbb{M})$, the process

$$(5.17) \quad N_s = \tau_s^\varepsilon(dP_{1-s}f)(W_s), \quad 0 \leq s \leq 1,$$

is a martingale, where $dP_{1-s}f$ denotes the exterior derivative of the function $P_{1-s}f$. As a consequence, for every $0 \leq s \leq 1$,

$$(5.18) \quad dP_s f(x) = \mathbb{E}_x(\tau_s^\varepsilon df(W_s)).$$

Proof. From Itô's formula and the definition of τ^ε , we have

$$\begin{aligned} dN_s = & \tau_s^\varepsilon \left(\nabla_{\circ dW_s} - \mathfrak{T}_{\circ dW_s}^\varepsilon - \frac{1}{2} \left(\frac{1}{\varepsilon} \mathbf{J}^2 - \frac{1}{\varepsilon} \delta_{\mathcal{H}}T + \mathfrak{Ric}_{\mathcal{H}} \right) ds \right) (dP_{1-s}f)(W_s) \\ & + \tau_s^\varepsilon \frac{d}{ds} (dP_{1-s}f)(W_s) ds. \end{aligned}$$

We now see that

$$\frac{d}{ds} (dP_{1-s}f) = -\frac{1}{2} dP_{1-s} Lf = -\frac{1}{2} dLP_{1-s}f = -\frac{1}{2} \square_\varepsilon dP_{1-s}f,$$

where we used Theorem 5.4. Observe that the bounded variation part of

$$\tau_s^\varepsilon \left(\nabla_{\circ dW_s} - \mathfrak{T}_{\circ dW_s}^\varepsilon - \frac{1}{2} \left(\frac{1}{\varepsilon} \mathbf{J}^2 - \frac{1}{\varepsilon} \delta_{\mathcal{H}}T + \mathfrak{Ric}_{\mathcal{H}} \right) ds \right) (dP_{1-s}f)(W_s)$$

is given by $\frac{1}{2} \tau_s^\varepsilon \square_\varepsilon dP_{1-s}f(W_s) ds$ which cancels out with the expression

$$\tau_s^\varepsilon \frac{d}{ds} (dP_{1-s}f)(W_s) ds$$

in the first equation. The martingale property follows from a bound similar to [2, Lemma 4.3] or [30, 31, Theorem 2.7]. \square

5.2.3. *Integration by parts formula for the damped Malliavin derivative.* We prove Theorem 5.19 in this section. Some of the key arguments may be found in [2, 4], however since our framework is more general here (for example, we do not assume the Yang-Mills condition that the horizontal divergence $\delta_{\mathcal{H}}T = 0$) and we now use the adjoint connection $\widehat{\nabla}^\varepsilon$ instead of the Bott connection, for the sake of self-containment, we give a complete proof.

Lemma 5.22. For $f \in C^\infty(\mathbb{M})$, and γ horizontal Cameron-Martin process

$$\begin{aligned} & \mathbb{E}_x \left(f(W_1) \int_0^1 \langle \gamma'(s), dB_s \rangle_{\mathcal{H}} \right) = \\ & \mathbb{E}_x \left(\left\langle \tau_1^\varepsilon df(W_1), \int_0^1 (\tau_s^{\varepsilon,*})^{-1} \widehat{\Theta}_s^\varepsilon \gamma'(s) ds \right\rangle \right). \end{aligned}$$

Proof. Consider again the martingale process N_s defined by (5.17). We have then for $f \in C^\infty(\mathbb{M})$

$$\begin{aligned} & \mathbb{E}_x \left(f(W_s) \int_0^s \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \right) = \mathbb{E}_x \left(f(W_s) \int_0^s \langle \widehat{\Theta}_r^\varepsilon \gamma'(r), \widehat{\Theta}_r^\varepsilon dB_r \rangle_{\mathcal{H}} \right) \\ & = \mathbb{E}_x \left((f(W_s) - \mathbb{E}_x(f(W_s))) \int_0^s \langle \widehat{\Theta}_r^\varepsilon \gamma'(r), \widehat{\Theta}_r^\varepsilon dB_r \rangle_{\mathcal{H}} \right) \\ & = \mathbb{E}_x \left(\int_0^s \langle dP_{s-r} f(W_r), \widehat{\Theta}_r^\varepsilon dB_r \rangle \int_0^s \langle \widehat{\Theta}_r^\varepsilon \gamma'(r), \widehat{\Theta}_r^\varepsilon dB_r \rangle_{\mathcal{H}} \right) \\ & = \mathbb{E}_x \left(\int_0^s \langle dP_{s-r} f(W_r), \widehat{\Theta}_r^\varepsilon \gamma'(r) \rangle dr \right) \\ & = \mathbb{E}_x \left(\int_0^s \langle \tau_r^\varepsilon dP_{s-r} f(W_r), (\tau_r^{\varepsilon,*})^{-1} \widehat{\Theta}_r^\varepsilon \gamma'(r) \rangle dr \right) \\ & = \mathbb{E}_x \left(\int_0^s \langle N_r, (\tau_r^{\varepsilon,*})^{-1} \widehat{\Theta}_r^\varepsilon \gamma'(r) \rangle dr \right) \\ & = \mathbb{E}_x \left(\left\langle N_s, \int_0^s (\tau_r^{\varepsilon,*})^{-1} \widehat{\Theta}_r^\varepsilon \gamma'(r) dr \right\rangle \right), \end{aligned}$$

where we integrated by parts in the last equality. \square

Remark 5.23. A similar proof as above actually yields that for $f \in C^\infty(\mathbb{M})$, γ horizontal Cameron-Martin process and $0 \leq s \leq 1$,

$$\begin{aligned} & \mathbb{E}_x \left(f(W_1) \int_s^1 \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_s \right) = \\ & \mathbb{E}_x \left(\left\langle \tau_1^\varepsilon df(W_1), \int_s^1 (\tau_r^{\varepsilon,*})^{-1} \widehat{\Theta}_r^\varepsilon \gamma'(r) dr \right\rangle \mid \mathcal{F}_s \right). \end{aligned}$$

Lemma 5.22 shows that integration by parts formula (5.16) holds for cylinder functions of the type $F = f(W_s)$. We now turn to the proof of Theorem 5.19 by using induction on n in a representation of a cylinder function F . To run the induction argument we need the following fact.

Proposition 5.24. Let $F = f(W_{s_1}, \dots, W_{s_n}) \in \mathcal{FC}^\infty(W_x(\mathbb{M}))$. We have

$$d\mathbb{E}_x(F) = \mathbb{E}_x \left(\sum_{i=1}^n \tau_{s_i}^\varepsilon d_i f(W_{s_1}, \dots, W_{s_n}) \right).$$

Proof. We will proceed by induction on n . Consider a cylinder function $F = f(W_{s_1}, \dots, W_{s_n})$. For $n = 1$ the statement follows from Lemma 5.21,

which implies that

$$d\mathbb{E}_x(f(W_{s_1})) = dP_{s_1}f(x) = \mathbb{E}_x(\tau_{s_1}^\varepsilon df(W_{s_1})).$$

Now we assume that the claim holds for any cylinder function of the form $F = f(W_{s_1}, \dots, W_{s_k})$ for any $k \leq n - 1$. By the Markov property we have

$$\mathbb{E}_x(F) = \mathbb{E}_x(\mathbb{E}(F | \mathcal{F}_{s_1})) = \mathbb{E}_x(g(W_{s_1})),$$

where $g(y) = \mathbb{E}_y(f(y, W_{s_2-s_1}, \dots, W_{s_n-s_1}))$. Therefore

$$d\mathbb{E}_x(F) = \mathbb{E}(\tau_{s_1}^\varepsilon dg(W_{s_1})).$$

By using the induction hypothesis, we obtain

$$\begin{aligned} dg(y) &= \mathbb{E}_y(d_1 f(y, W_{s_2-s_1}, \dots, W_{s_n-s_1})) + \\ &\quad \mathbb{E}_y \left(\sum_{i=2}^n \tau_{s_i-s_1}^\varepsilon d_i f(y, W_{s_2-s_1}, \dots, W_{s_n-s_1}) \right) \\ &= \mathbb{E}_y \left(\sum_{i=1}^n \tau_{s_i-s_1}^\varepsilon d_i f(y, W_{s_2-s_1}, \dots, W_{s_n-s_1}) \right). \end{aligned}$$

By the multiplicative property of τ^ε and the Markov property of W we have

$$\begin{aligned} \mathbb{E}_{W_{s_1}}(\tau_{s_i-s_1}^\varepsilon d_i f(y, W_{s_2-s_1}, \dots, W_{s_n-s_1})) &= \\ (\tau_{s_1}^\varepsilon)^{-1} \mathbb{E}(\tau_{s_i}^\varepsilon d_i f(W_{s_1}, \dots, W_{s_n}) | \mathcal{F}_{s_1}). \end{aligned}$$

Therefore we conclude

$$d\mathbb{E}_x(F) = \mathbb{E}_x \left(\sum_{i=1}^n \tau_{s_i}^\varepsilon d_i f(W_{s_1}, \dots, W_{s_n}) \right).$$

□

Remark 5.25. As expected, the expression

$$\mathbb{E}_x \left(\sum_{i=1}^n \tau_{s_i}^\varepsilon d_i f(W_{s_1}, \dots, W_{s_n}) \right)$$

is independent of the choice of the representation of the cylinder function F , as follows from Remark 5.14.

Proof of Theorem 5.19. We use induction on n in a representation of the cylinder function F . More precisely, we would like to show that for any $F = f(W_{s_1}, \dots, W_{s_n}) \in \mathcal{FC}^\infty(W_x(\mathbb{M}))$ and $s \leq s_1$ we have

$$\begin{aligned} (5.19) \quad \mathbb{E}_x \left(F \int_s^{s_n} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \Big| \mathcal{F}_s \right) &= \\ \mathbb{E}_x \left(\sum_{i=1}^n \langle d_i f(W_{s_1}, \dots, W_{s_n}), \tau_{s_i}^{\varepsilon,*} \int_s^{s_i} (\tau_r^{\varepsilon,*})^{-1} \gamma'(r) dr \rangle \Big| \mathcal{F}_s \right). \end{aligned}$$

The case $n = 1$ is Lemma 5.22 and Remark 5.23. Assume that (5.19) holds for any cylinder function F represented by a partition of size $n - 1$ for $n \geq 2$. Let $F = f(W_{s_1}, \dots, W_{s_n}) \in \mathcal{FC}^\infty(W_x(\mathbb{M}))$. We have for $s \leq s_1$,

$$\begin{aligned}
& \mathbb{E}_x \left(F \int_s^1 \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_s \right) \\
&= \mathbb{E}_x \left(F \int_s^{s_n} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_s \right) \\
&= \mathbb{E}_x \left(F \int_s^{s_1} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_s \right) + \mathbb{E}_x \left(F \int_{s_1}^{s_n} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_s \right) \\
&= \mathbb{E}_x \left(\mathbb{E}_x(F \mid \mathcal{F}_{s_1}) \int_s^{s_1} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_s \right) + \\
& \quad \mathbb{E}_x \left(\mathbb{E}_x \left(F \int_{s_1}^{s_n} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_{s_1} \right) \mid \mathcal{F}_s \right).
\end{aligned}$$

By the Markov property we have

$$\mathbb{E}_x(F \mid \mathcal{F}_{s_1}) = g(W_{s_1}),$$

where $g(y) = \mathbb{E}_y(f(y, W_{s_2-s_1}, \dots, W_{s_n-s_1}))$. Thus by Lemma 5.22 and Remark 5.23

$$\begin{aligned}
& \mathbb{E}_x \left(\mathbb{E}_x(F \mid \mathcal{F}_{s_1}) \int_s^{s_1} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_s \right) = \\
& \mathbb{E}_x \left(g(W_{s_1}) \int_s^{s_1} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_s \right) = \\
& \mathbb{E}_x \left(\left\langle dg(W_{s_1}), (\tau_{s_1}^\varepsilon)^* \int_s^{s_1} (\tau_r^{\varepsilon,*})^{-1} \widehat{\Theta}_r^\varepsilon \gamma'(r) dr \right\rangle \mid \mathcal{F}_s \right)
\end{aligned}$$

Now according to Proposition 5.24

$$dg(y) = \mathbb{E}_y \left(\sum_{i=1}^n \tau_{s_i-s_1}^\varepsilon d_i f(y, W_{s_2-s_1}, \dots, W_{s_n-s_1}) \right).$$

Using the fact that

$$\begin{aligned}
& \mathbb{E}_{W_{s_1}} \left(\tau_{s_i-s_1}^\varepsilon d_i f(y, W_{s_2-s_1}, \dots, W_{s_n-s_1}) \right) = \\
& (\tau_{s_1}^\varepsilon)^{-1} \mathbb{E}_x \left(\tau_{s_i}^\varepsilon d_i f(W_{s_1}, \dots, W_{s_n}) \mid \mathcal{F}_{s_1} \right),
\end{aligned}$$

we conclude

$$\begin{aligned}
& \mathbb{E}_x \left(\mathbb{E}_x(F \mid \mathcal{F}_{s_1}) \int_s^{s_1} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \mid \mathcal{F}_s \right) \\
&= \mathbb{E}_x \left(\sum_{i=1}^n \langle d_i f(W_{s_1}, \dots, W_{s_n}), \tau_{s_i}^{\varepsilon,*} \int_s^{s_1} (\tau_r^{\varepsilon,*})^{-1} \widehat{\Theta}_r^\varepsilon \gamma'(r) dr \rangle \mid \mathcal{F}_s \right).
\end{aligned}$$

Using the induction hypothesis that (5.19) holds for $n - 1$ we see that

$$\begin{aligned} & \mathbb{E}_x \left(F \int_{s_1}^{s_n} \langle \gamma'(r), dB_r \rangle_{\mathcal{H}} \middle| \mathcal{F}_{s_1} \right) = \\ & \mathbb{E}_x \left(\sum_{i=1}^n \langle d_i f(W_{s_1}, \dots, W_{s_n}), \tau_{s_i}^{\varepsilon,*} \int_{s_1}^{s_i} (\tau_r^{\varepsilon,*})^{-1} \gamma'(r) dr \rangle \middle| \mathcal{F}_{s_1} \right). \end{aligned}$$

□

5.2.4. Integration by parts formula for the directional derivatives. In this section we prove Theorem 5.20. One of the main ingredients B. Driver used in [17] in the Riemannian case was the orthogonal invariance of the Brownian motion to filter out redundant noise. As a complement to Lemma 5.22, we first prove the following result.

Lemma 5.26. Let $\{\mathcal{O}_s\}_{0 \leq s \leq 1}$ be a continuous \mathcal{F} -adapted process taking values in the space of skew-symmetric endomorphisms of \mathcal{H}_x such that $\mathbb{E} \left(\int_0^1 \|\mathcal{O}_s\|^2 ds \right) < \infty$, where $\|\mathcal{O}_s\|^2 = \mathbf{Tr}(\mathcal{O}_s^* \mathcal{O}_s)$. For $f \in C^\infty(\mathbb{M})$, we have

$$\mathbb{E}_x \left(\left\langle \tau_1^\varepsilon df(W_1), \int_0^1 (\tau_s^{\varepsilon,*})^{-1} \widehat{\Theta}_s^\varepsilon \left(\mathcal{O}_s dB_s - \frac{1}{2} T_{\mathcal{O}_s}^\varepsilon ds \right) \right\rangle \right) = 0,$$

where $T_{\mathcal{O}_s}^\varepsilon$ is the tensor given in a horizontal frame e_1, \dots, e_n by

$$T_{\mathcal{O}_s}^\varepsilon = \sum_{i=1}^n (\widehat{\Theta}_s^\varepsilon)^{-1} T^\varepsilon(e_i, \widehat{\Theta}_s^\varepsilon \mathcal{O}_s (\widehat{\Theta}_s^\varepsilon)^{-1} e_i).$$

Proof. Recall that we considered the following martingale in (5.17)

$$N_s = \tau_s^\varepsilon (dP_{1-s} f)(W_s), \quad 0 \leq s \leq 1.$$

We have then

$$\begin{aligned} & \mathbb{E}_x \left(\left\langle \tau_1^\varepsilon df(W_1), \int_0^1 (\tau_s^{\varepsilon,*})^{-1} \widehat{\Theta}_s^\varepsilon \mathcal{O}_s dB_s \right\rangle \right) = \\ & \mathbb{E}_x \left(\left\langle N_1, \int_0^1 (\tau_s^{\varepsilon,*})^{-1} \widehat{\Theta}_s^\varepsilon \mathcal{O}_s dB_s \right\rangle \right). \end{aligned}$$

From the proof of Lemma 5.21, we have

$$\begin{aligned} & dN_s \\ &= \tau_s^\varepsilon \left(\nabla_{odW_s} - \mathfrak{I}^\varepsilon_{odW_s} - \frac{1}{2} \left(\frac{1}{\varepsilon} \mathbf{J}^2 - \frac{1}{\varepsilon} \delta_{\mathcal{H}} T + \mathfrak{Ric}_{\mathcal{H}} \right) ds \right) (dP_{1-s} f)(W_s) \\ &+ \tau_s^\varepsilon \frac{d}{ds} (dP_{1-s} f)(W_s) ds \\ &= \tau_s^\varepsilon \left(\nabla_{\widehat{\Theta}_s^\varepsilon dB_s} - \mathfrak{I}^\varepsilon_{\widehat{\Theta}_s^\varepsilon dB_s} \right) (dP_{1-s} f)(W_s) = \tau_s^\varepsilon \nabla_{\widehat{\Theta}_s^\varepsilon dB_s}^\varepsilon dP_{1-s} f(W_s). \end{aligned}$$

where, as before, ∇^ε denotes the connection $\nabla - \mathfrak{I}^\varepsilon$. Let us denote by $\mathbf{Hess}^\varepsilon$ the Hessian for the connection ∇^ε . One has therefore

$$\begin{aligned} & \mathbb{E}_x \left(\left\langle N_1, \int_0^1 (\tau_s^{\varepsilon,*})^{-1} \widehat{\Theta}_s^\varepsilon \mathcal{O}_s dB_s \right\rangle \right) = \\ & \mathbb{E}_x \left(\int_0^1 \mathbf{Hess}^\varepsilon P_{1-s} f(\widehat{\Theta}_s^\varepsilon dB_s, \widehat{\Theta}_s^\varepsilon \mathcal{O}_s dB_s)(W_s) \right) \end{aligned}$$

Due to the skew symmetry of \mathcal{O} and the fact that for $h \in C^\infty(\mathbb{M})$, $X, Y \in \Gamma^\infty(\mathbb{M})$,

$$\mathbf{Hess}^\varepsilon h(X, Y) - \mathbf{Hess}^\varepsilon h(Y, X) = T^\varepsilon(X, Y)h,$$

we deduce

$$\begin{aligned} & \mathbb{E}_x \left(\left\langle N_1, \int_0^1 (\tau_s^{\varepsilon,*})^{-1} \widehat{\Theta}_s^\varepsilon \mathcal{O}_s dB_s \right\rangle \right) = \\ & \frac{1}{2} \mathbb{E}_x \left(\int_0^1 \left\langle dP_{1-s} f, \widehat{\Theta}_s^\varepsilon T_{\mathcal{O}_s}^\varepsilon \right\rangle ds \right) = \\ & \frac{1}{2} \mathbb{E}_x \left(\int_0^1 \left\langle N_s, (\tau_s^{\varepsilon,*})^{-1} \widehat{\Theta}_s^\varepsilon T_{\mathcal{O}_s}^\varepsilon \right\rangle ds \right). \end{aligned}$$

Integrating by parts the right hand side yields the conclusion. \square

We are now in position to prove the integration by parts formula for cylinder functions of the type $F = f(W_s)$.

Lemma 5.27. Let v be a tangent process. For $f \in C^\infty(\mathbb{M})$,

$$\begin{aligned} & \mathbb{E}_x \left(\left\langle df(W_1), \widehat{\Theta}_1^\varepsilon v(1) \right\rangle \right) = \\ & \mathbb{E}_x \left(f(W_1) \int_0^1 \left\langle v_{\mathcal{H}}'(s) + \frac{1}{2} (\widehat{\Theta}_s^\varepsilon)^{-1} \mathfrak{Ric}_{\mathcal{H}} \widehat{\Theta}_s^\varepsilon v_{\mathcal{H}}(s), dB_s \right\rangle_{\mathcal{H}} \right). \end{aligned}$$

Proof. Let v be a tangent process. We define

$$h(s) = v(s) - \int_0^s (\widehat{\Theta}_r^\varepsilon)^{-1} T(\widehat{\Theta}_r^\varepsilon \circ dB_r, \widehat{\Theta}_r^\varepsilon v(r)).$$

By definition of tangent processes, we have that $h = v_{\mathcal{H}}$ is a horizontal Cameron-Martin process. By Equation (5.4) we have

$$\widehat{T}^\varepsilon(\circ dW_s, \widehat{\Theta}_s^\varepsilon v(s)) = T(\widehat{\Theta}_s^\varepsilon \circ dB_s, \widehat{\Theta}_s^\varepsilon v(s)) - \frac{1}{\varepsilon} J_{\widehat{\Theta}_s^\varepsilon v(s)}(\widehat{\Theta}_s^\varepsilon \circ dB_s).$$

Therefore we get

$$\begin{aligned}
 & dv(s) + (\widehat{\Theta}_s^\varepsilon)^{-1} \left(-\widehat{T}^\varepsilon(\circ dW_s, \cdot) + \frac{1}{2} \left(\frac{1}{\varepsilon} \mathbf{J}^2 - \frac{1}{\varepsilon} \delta_{\mathcal{H}}^* T + \mathfrak{Ric}_{\mathcal{H}} \right) ds \right) \widehat{\Theta}_s^\varepsilon v(s) \\
 &= dh(s) + \frac{1}{\varepsilon} (\widehat{\Theta}_s^\varepsilon)^{-1} J_{\widehat{\Theta}_s^\varepsilon v(s)} (\widehat{\Theta}_s^\varepsilon \circ dB_s) \\
 &+ \frac{1}{2} (\widehat{\Theta}_s^\varepsilon)^{-1} \left(\frac{1}{\varepsilon} \mathbf{J}^2 - \frac{1}{\varepsilon} \delta_{\mathcal{H}}^* T + \mathfrak{Ric}_{\mathcal{H}} \right) \widehat{\Theta}_s^\varepsilon h(s) ds \\
 &= dh(s) + \frac{1}{\varepsilon} (\widehat{\Theta}_s^\varepsilon)^{-1} J_{\widehat{\Theta}_s^\varepsilon v(s)} (\widehat{\Theta}_s^\varepsilon dB_s) + \frac{1}{2} (\widehat{\Theta}_s^\varepsilon)^{-1} (\mathfrak{Ric}_{\mathcal{H}}) \widehat{\Theta}_s^\varepsilon h(s) ds.
 \end{aligned}$$

In the last computation, the transformation of the Stratonovitch differential

$$(\widehat{\Theta}_s^\varepsilon)^{-1} J_{\widehat{\Theta}_s^\varepsilon v(s)} (\widehat{\Theta}_s^\varepsilon \circ dB_s)$$

into Itô's differential

$$(\widehat{\Theta}_s^\varepsilon)^{-1} J_{\widehat{\Theta}_s^\varepsilon v(s)} (\widehat{\Theta}_s^\varepsilon dB_s)$$

is similar to (4.5). It is then a consequence of Itô's formula that

$$v(s) = (\widehat{\Theta}_s^\varepsilon)^{-1} \tau_s^{\varepsilon,*} \int_0^s (\tau_r^{\varepsilon,*})^{-1} \widehat{\Theta}_r^\varepsilon \circ dM_r,$$

where

$$dM_s = dh(s) + \frac{1}{\varepsilon} (\widehat{\Theta}_s^\varepsilon)^{-1} J_{\widehat{\Theta}_s^\varepsilon v(s)} \widehat{\Theta}_s^\varepsilon dB_s + \frac{1}{2} (\widehat{\Theta}_s^\varepsilon)^{-1} (\mathfrak{Ric}_{\mathcal{H}}) \widehat{\Theta}_s^\varepsilon h(s) ds.$$

Converting the Stratonovitch integral into Itô's integral finally yields

$$\begin{aligned}
 v(s) &= (\widehat{\Theta}_s^\varepsilon)^{-1} \tau_s^{\varepsilon,*} \int_0^s (\tau_r^{\varepsilon,*})^{-1} \widehat{\Theta}_r^\varepsilon \left(dh(s) + \frac{1}{2} (\widehat{\Theta}_s^\varepsilon)^{-1} (\mathfrak{Ric}_{\mathcal{H}}) \widehat{\Theta}_s^\varepsilon h(s) ds \right. \\
 &\left. + \mathcal{O}_s dB_s - \frac{1}{2} T_{\mathcal{O}_s}^\varepsilon ds \right),
 \end{aligned}$$

with

$$\mathcal{O}_s = \frac{1}{\varepsilon} (\widehat{\Theta}_s^\varepsilon)^{-1} J_{\widehat{\Theta}_s^\varepsilon v(s)} \widehat{\Theta}_s^\varepsilon.$$

Since \mathcal{O}_s is a skew-symmetric horizontal endomorphism, one can conclude from Lemmas 5.22 and 5.26 that

$$\begin{aligned}
 & \mathbb{E}_x \left(\left\langle df(W_s), \widehat{\Theta}_t^\varepsilon v(s) \right\rangle \right) \\
 &= \mathbb{E}_x \left(f(W_s) \int_0^1 \left\langle v'_{\mathcal{H}}(s) + \frac{1}{2} (\widehat{\Theta}_s^\varepsilon)^{-1} \mathfrak{Ric}_{\mathcal{H}} \widehat{\Theta}_s^\varepsilon v_{\mathcal{H}}(s), dB_s \right\rangle_{\mathcal{H}} \right)
 \end{aligned}$$

because $h(s) = v_{\mathcal{H}}(s)$. \square

Now Theorem 5.20 can be proven using induction on n in the representation of a cylinder function F . The case $n = 1$ is Lemma 5.27, and showing the induction step is similar to how Theorem 5.19 has been proven, so for the sake of conciseness of the paper, we omit the details. As a direct corollary of Theorem 5.20, we obtain the following.

Corollary 5.28. Let $F, G \in \mathcal{FC}^\infty(W_x(\mathbb{M}))$ and v be a tangent process. We have

$$\mathbb{E}_x(F\mathbf{D}_v G) = \mathbb{E}_x(G\mathbf{D}_v^* F),$$

where

$$\mathbf{D}_v^* = -\mathbf{D}_v + \int_0^1 \left\langle v'_{\mathcal{H}}(s) + \frac{1}{2}(\widehat{\Theta}_s^\varepsilon)^{-1} \mathfrak{Ric}_{\mathcal{H}} \widehat{\Theta}_s^\varepsilon v_{\mathcal{H}}(s), dB_s \right\rangle_{\mathcal{H}}.$$

Proof. By Theorem 5.20, we have

$$\mathbb{E}_x(\mathbf{D}_v(FG)) = \mathbb{E}_x \left(FG \int_0^1 \left\langle v'_{\mathcal{H}}(s) + \frac{1}{2}(\widehat{\Theta}_s^\varepsilon)^{-1} \mathfrak{Ric}_{\mathcal{H}} \widehat{\Theta}_s^\varepsilon v_{\mathcal{H}}(s), dB_s \right\rangle_{\mathcal{H}} \right).$$

Since $\mathbf{D}_v(FG) = F\mathbf{D}_v(G) + G\mathbf{D}_v(F)$, the conclusion follows immediately. \square

5.3. Examples.

5.3.1. Riemannian submersions. In this section, we verify that the integration by parts formula we obtained for the directional derivatives is consistent with and generalizes the formulas known in the Riemannian case. Let us assume here that the foliation on \mathbb{M} comes from a totally geodesic submersion $\pi : (\mathbb{M}, g) \rightarrow (\mathbb{B}, j)$ as in Example 2.1. Since the submersion has totally geodesic fibers, π is harmonic and the projected process

$$W_s^{\mathbb{B}} = \pi(W_s)$$

is a Brownian motion on \mathbb{B} started at $\pi(x)$. Observe that from the definition of submersion, the derivative map $T_x\pi$ is an isometry from \mathcal{H}_x to $T_x\mathbb{B}$. From Example 2.3, the connection $\widehat{\nabla}^\varepsilon$ projects down to the Levi-Civita connection on \mathbb{B} . Therefore the stochastic parallel transport $\widehat{\Theta}_t^\varepsilon$ projects down to the stochastic parallel transport for the Levi-Civita connection along the paths of $\{W_s^{\mathbb{B}}\}_{0 \leq s \leq 1}$. More precisely,

$$\parallel_{0,s} = T_{W_s} \pi \circ \widehat{\Theta}_s^\varepsilon \circ (T_x \pi)^{-1},$$

where $\parallel_{0,s} : T_{\pi(x)}\mathbb{B} \rightarrow T_{W_s^{\mathbb{B}}}\mathbb{B}$ is the stochastic parallel transport for the Levi-Civita connection along the paths of $\{W_s^{\mathbb{B}}\}_{0 \leq s \leq 1}$. Consider now a Cameron-Martin process $\{h(s)\}_{0 \leq s \leq 1}$ in $T_{\pi(x)}\mathbb{B}$ and a cylinder function $F = f(W_{s_1}^{\mathbb{B}}, \dots, W_{s_n}^{\mathbb{B}})$ on \mathbb{B} . The function $F = f(\pi(W_{s_1}), \dots, \pi(W_{s_n}))$ is then in $\mathcal{FC}^\infty(W_{\mathcal{H}}(\mathbb{M}))$ (Refer to [19, Definition 7.4]). Using Theorem 5.20, one gets

$$\mathbb{E}_x(\mathbf{D}_v F) = \mathbb{E}_x \left(F \int_0^1 \left\langle v'_{\mathcal{H}}(s) + \frac{1}{2}(\widehat{\Theta}_s^\varepsilon)^{-1} \mathfrak{Ric}_{\mathcal{H}} \widehat{\Theta}_s^\varepsilon v_{\mathcal{H}}(s), dB_s \right\rangle_{\mathcal{H}} \right),$$

where $v_{\mathcal{H}}$ is the horizontal lift of h , that is, $v_{\mathcal{H}} = (T_x\pi)^{-1}h$. By definition, we have

$$\begin{aligned} \mathbf{D}_v F &= \sum_{i=1}^n \left\langle d_i f(W_{s_1}^{\mathbb{B}}, \dots, W_{s_n}^{\mathbb{B}}), (T_{W_{s_i}}\pi) \circ \widehat{\Theta}_{s_i}^{\varepsilon} v(s_i) \right\rangle \\ &= \sum_{i=1}^n \left\langle d_i f(W_{s_1}^{\mathbb{B}}, \dots, W_{s_n}^{\mathbb{B}}), \llbracket_{0, s_i} h(s_i) \right\rangle \end{aligned}$$

It is easy to check that $\mathfrak{Ric}_{\mathcal{H}}$ is the horizontal lift of the Ricci curvature $\mathfrak{Ric}^{\mathbb{B}}$ of \mathbb{B} . Therefore, the integration by parts formula for the directional derivative $\mathbf{D}_v F$ can be rewritten as follows.

$$\begin{aligned} &\mathbb{E}_x \left(\sum_{i=1}^n \left\langle d_i f(W_{s_1}^{\mathbb{B}}, \dots, W_{s_n}^{\mathbb{B}}), \llbracket_{0, s_i} h(s_i) \right\rangle \right) \\ &= \mathbb{E}_x \left(F \int_0^1 \left\langle h'(s) + \frac{1}{2} \llbracket_{0, s}^{-1} \mathfrak{Ric}^{\mathbb{B}} \llbracket_{0, s} h(s), dB_s^{\mathbb{B}} \right\rangle_{T_{\pi(x)}\mathbb{B}} \right), \end{aligned}$$

where $B^{\mathbb{B}}$ is the Brownian motion on $T_{\pi(x)}\mathbb{B}$ given by $B^{\mathbb{B}} = T_x\pi(B)$. This is exactly Driver's integration by parts formula in [17] for the Riemannian Brownian motion $X^{\mathbb{B}}$.

5.3.2. K -contact manifolds. In this section, we assume that the Riemannian foliation on \mathbb{M} is the Reeb foliation of a K -contact structure. The Reeb vector field on \mathbb{M} will be denoted by R and the almost complex structure by \mathbf{J} . The torsion of the Bott connection is then

$$T(X, Y) = \langle \mathbf{J}X, Y \rangle_{\mathcal{H}} R.$$

Therefore with the previous notation, one has

$$J_Z X = \langle Z, R \rangle \mathbf{J}X.$$

and the vertical part of a tangent process is given by

$$\begin{aligned} v_{\mathcal{V}}(s) &= - \int_0^s (\widehat{\Theta}_r^{\varepsilon})^{-1} T(\widehat{\Theta}_r^{\varepsilon} \circ dB_r, \widehat{\Theta}_r^{\varepsilon} v_{\mathcal{H}}(r)) \\ &= \int_0^s ((\widehat{\Theta}_r^{\varepsilon})^{-1} R) \langle \mathbf{J} \widehat{\Theta}_r^{\varepsilon} v_{\mathcal{H}}(r), \widehat{\Theta}_r^{\varepsilon} \circ dB_r \rangle_{\mathcal{H}}. \end{aligned}$$

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