

QUASI-INVARIANCE FOR HEAT KERNEL MEASURES ON SUB-RIEMANNIAN INFINITE-DIMENSIONAL HEISENBERG GROUPS

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ABSTRACT. We study heat kernel measures on sub-Riemannian infinite-dimensional Heisenberg-like Lie groups. In particular, we show that Cameron-Martin type quasi-invariance results hold in this subelliptic setting and give L^p -estimates for the Radon-Nikodym derivatives. The main ingredient in our proof is a generalized curvature-dimension estimate which holds on approximating finite-dimensional projection groups. Such estimates were first introduced by Baudoin and Garofalo in [4].

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1. INTRODUCTION

We prove Cameron-Martin type quasi-invariance results for subelliptic heat kernel measures on infinite-dimensional Heisenberg-like groups. These groups were first defined in [11] and quasi-invariance was proved for elliptic heat kernel measures in this setting. Quasi-invariance results are of interest, for example, in the study of smoothness of measures on infinite-dimensional spaces. In finite dimensions, one typically defines smoothness as absolute continuity with respect to some reference measure and smoothness of the associated density. In infinite dimensions, in the absence of a canonical reference measure, alternative interpretations of smoothness must be made; see for example [5, 10, 22, 23]. In particular, in an infinite-dimensional setting, it is natural to interpret quasi-invariance as a smoothness property.

Quasi-invariance of heat kernel measures in infinite dimensions has previously been the subject of much study in elliptic settings; see for example the review [8] and references therein. Typically the proofs in the elliptic case rely on lower bounds on the Ricci curvature (as was the case in [11]); of course such lower bounds are unavailable in a subelliptic setting. To the authors' knowledge, this is the first quasi-invariance result in an infinite-dimensional subelliptic setting.

1.1. Statement of Results. Let (W, H, μ) be an abstract Wiener space and let \mathbf{C} be a finite-dimensional inner product space. Let $\mathfrak{g} = W \times \mathbf{C}$ be an infinite-dimensional Heisenberg-like Lie algebra, which is constructed as an infinite-dimensional step 2 nilpotent Lie algebra with continuous Lie bracket satisfying the following condition:

$$(1.1) \quad [W, W] = \mathbf{C}.$$

Let G denote $W \times \mathbf{C}$ thought of as a group with operation

$$g_1 \cdot g_2 = g_1 + g_2 + \frac{1}{2}[g_1, g_2].$$

Then G is a Lie group with Lie algebra \mathfrak{g} , and G contains the subgroup $G_{CM} = H \times \mathbf{C}$ which has Lie algebra \mathfrak{g}_{CM} . See Section 3.2 for definitions and details.

Now let $\{B_t\}_{t \geq 0}$ be a Brownian motion on W . The solution to the stochastic differential equation

$$(1.2) \quad dg_t = g_t \circ dB_t \quad \text{with } g_0 = e$$

is a Brownian motion on G , which is defined explicitly in Proposition 5.1 and Definition 5.2. For all $t > 0$, let $\nu_t = \text{Law}(g_t)$ denote the heat kernel measure at time t . At this point, let us briefly comment that, as mentioned in [15], it may seem at first glance that the restriction $\dim(\mathbf{C}) < \infty$ implies that this subelliptic example is in some sense only finitely many steps from being elliptic. However, this is truly a subelliptic model and the topologies one must deal with in this setting significantly change the standard analysis

and introduce several non-trivial complications not present in the elliptic case. For further discussion, see Section 1.3 of [15].

The main results of the present paper are presented in Section 5.2. Namely, in Theorem 5.9, we prove that ν_t is quasi-invariant under translation by elements of G_{CM} and obtain bounds on the L^p -norms of the Radon-Nikodym derivatives. Then given the equivalence of measures and the L^p -estimates on the associated densities, we are able to immediately prove in Corollary 5.11 that the semi-group is strong Feller on G_{CM} .

To put these results in context, again recall that if W is finite-dimensional, then (1.1) implies that $\text{span}\{(\xi_i, 0), [(\xi_i, 0), (\xi_j, 0)]\} = \mathfrak{g}$, where $\{\xi_i\}_{i=1}^{\dim(W)}$ is some orthonormal basis of W , and thus we have satisfaction of Hörmander’s condition. This then implies that ν_t is a smooth measure, in the sense that ν_t is absolutely continuous with respect to Haar measure on $G = W \times \mathbb{C}$ and its density is a smooth function on G . If W and thus G is infinite-dimensional, however, the notion of smoothness is not so well-defined, and quasi-invariance may be interpreted as a first step toward proving ν_t is a “strictly positive” smooth measure.

1.2. Discussion of proofs. Functional inequalities provide a powerful tool to study the problem of the equivalence of heat kernel measures. In particular, it is a well-known fact that, on a finite-dimensional complete Riemannian manifold \mathbb{M} with non-negative Ricci curvature, the heat semi-group $\{P_t\}_{t \geq 0}$ satisfies the Harnack type inequality

$$(1.3) \quad (P_t f)^\alpha(x) \leq P_t f^\alpha(y) \exp\left(\frac{\alpha}{\alpha - 1} \frac{d^2(x, y)}{4t}\right),$$

where $x, y \in \mathbb{M}$, $f \in L^\infty(\mathbb{M})$ with $f \geq 0$, and $\alpha > 1$ (see for example [27]). Using the above inequality with indicator functions immediately implies that the heat kernels measures $p_t(x, dz)$ and $p_t(y, dz)$ are equivalent. Of course, in a finite-dimensional framework, the latter is obvious and may be seen from the positivity of the heat kernel. But the relevant fact here is that the functional inequality (1.3) is independent of the dimension of the manifold \mathbb{M} and we may therefore hope that it holds even in some infinite-dimensional settings, where equivalence of measures is a highly non-trivial problem.

Again, a lower bound on the Ricci curvature of \mathbb{M} typically plays a major role in the proof of inequalities like (1.3) and such bounds are unavailable in our subelliptic setting. However, in a recent work [4], Baudoin and Garofalo introduced a generalized curvature-dimension inequality that holds in a general class of sub-Riemannian settings. The main idea is to control sub-Riemannian curvature quantities both in the horizontal and the vertical directions. In the present paper, we prove that a uniform generalized curvature-dimension inequality holds on appropriate finite-dimensional approximation groups of G , and that, as a consequence, a uniform version of (1.3) holds on the finite-dimensional approximation groups. A by-product

of this is a Cameron-Martin type quasi-invariance result for the heat kernel measure on G .

Let us emphasize that our approach is actually quite general and does not rely on the specific nature of the infinite-dimensional Heisenberg-like groups we consider here. The main ingredient is the existence of *good* finite-dimensional approximations on which uniform generalized curvature-dimension bounds hold. As was done in [12] in the elliptic setting, one could significantly generalize the method to include other infinite-dimensional subelliptic settings.

The organization of the paper is briefly as follows. In Section 2 we review results for subelliptic heat kernels on finite-dimensional Lie groups under the assumption of generalized curvature-dimension estimates and satisfaction of certain commutation relations. In particular, in Section 2.1, we show that under these assumptions reverse Poincaré and log Sobolev estimates hold, and in Section 2.2 we show how these estimates in turn give Wang type and integrated Harnack inequalities. The ideas in this section in a finite-dimensional setting are discussed in more detail and greater generality in [3].

In Section 3 we review the definition of the infinite-dimensional Heisenberg-like groups first considered in [11]. In the present paper, we choose a topological structure that is better adapted to subellipticity, and in this way our construction parallels [15], where subelliptic heat kernel measure was also studied. In this section, we also review the properties of infinite-dimensional Heisenberg-like groups required for the sequel, as well as recalling the Cameron-Martin subgroup and the finite-dimensional projection groups which serve as approximations.

In Section 4.1 we show that the curvature-dimension estimates and commutation relations considered in Section 2 are satisfied on the infinite-dimensional Heisenberg-like groups and their finite-dimensional projections, and thus in Section 4.2 we are able to show that reverse Poincaré and log Sobolev inequalities, as well as Wang type Harnack inequalities, hold on the finite-dimensional projection groups.

Finally, in Section 5.1 we review the construction of subelliptic heat kernel measures on the infinite-dimensional Heisenberg-like groups, and in Section 5.2 we show how the Wang type Harnack inequalities proved in Section 4.2 for the finite-dimensional projection groups imply quasi-invariance of the subelliptic heat kernel measure under translation by elements of the Cameron-Martin space, as well as yielding L^p -estimates for the Radon-Nikodym derivatives. We also show how the quasi-invariance and L^p -estimates then immediately imply that P_t is strong Feller on G_{CM} .

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2. FUNCTIONAL INEQUALITIES ON FINITE-DIMENSIONAL LIE GROUPS

For the whole of this section, G will denote a real *finite-dimensional* connected unimodular Lie group with Lie algebra \mathfrak{g} and identity element e . Let dx denote bi-invariant Haar measure on G . For $x \in G$, let L_x and R_x denote left and right translation by x , respectively. For $A \in \mathfrak{g}$, let \tilde{A} denote the unique left invariant vector field on G such that $\tilde{A}(e) = A \in \mathfrak{g}$, that is, $\tilde{A}(g) = L_{x*}A$.

We will assume there is a linearly independent collection $\{X_i\}_{i=1}^n \subset \mathfrak{g}$ for which there exists some r such that

$$(2.1) \quad \mathfrak{g} = \text{span}\{X_{i_1}, [X_{i_1}, X_{i_2}], \dots, [X_{i_1}, [X_{i_2}, \dots, [X_{i_{r-1}}, X_{i_r}] \dots]] : \\ i_1, \dots, i_r = 1, \dots, n\}.$$

We will refer to $\mathcal{H} := \text{span}(\{X_i\}_{i=1}^n)$ as the *horizontal directions*, and we will suppose that \mathfrak{g} is equipped with an inner product for which $\{X_i\}_{i=1}^n$ is an orthonormal basis of \mathcal{H} . In particular, (2.1) implies that $\{\tilde{X}_i\}_{i=1}^n$ is a Hörmander set of vector fields, in the sense that the \tilde{X}_i 's and all their commutators up to order r generate $T_g G$ for all $g \in G$. Thus, by the classical Hörmander's theorem [18] the second order differential operator

$$L = \sum_{i=1}^n \tilde{X}_i^2$$

is hypoelliptic, and, for all $t > 0$, there exists a smooth kernel $p_t : G \times G \rightarrow \mathbb{R}^+$ so that

$$P_t f(x) := e^{t\bar{L}/2} f(x) = \int_G f(y) p_t(x, y) dy,$$

for all $f \in L^2(G, dy)$, where \bar{L} denotes the $L^2(G, dy)$ closure of $L|_{C_c^\infty(G)}$. The heat semi-group $\{P_t\}_{t>0}$ is a symmetric Markov semi-group, and, since p_t is smooth, the mapping $(t, x) \mapsto P_t f(x)$ is smooth on $(0, \infty) \times G$ for any $f \in L^p(dx)$, $p \in [1, \infty]$.

In a slight abuse of notation, we will write that $p_t(e, x) = p_t(x)$. In particular, the left invariance of L implies that p_t is a left convolution kernel and thus $p_t(x, y) = p_t(xy^{-1})$. We call the measure $p_t dy$ the *heat kernel measure* on G associated to L . By construction, the operator P_t commutes with left translations, and since our Lie group is unimodular the Haar measure is bi-invariant. Thus, $p_t(x, y) = p_t(e, x^{-1}y)$. Since it is additionally known that p_t is a symmetric kernel, we have the following lemma of some standard properties of the heat kernel.

Lemma 2.1. *For all $t > 0$ and $x, y \in G$,*

- (1) $p_t(x, y) = p_t(x^{-1}y) = p_t(y^{-1}x)$, and
- (2) $p_t(x^{-1}) = p_t(x)$.

Another well-known interpretation of the heat kernel measure $p_t dy$ is as the distribution in t of Brownian motion on G . For $\{b_i\}_{i=1}^n$ independent real-valued Brownian motions, $b_t = \sum_{i=1}^n b_t^i X_i$ is a Brownian motion on \mathcal{H} . Now

for $t > 0$ let g_t denote the solution to the following Stratonovitch stochastic differential equation:

$$dg_t = g_t \circ db_t := L_{g_t*} \circ db_t = \sum_{i=1}^n \tilde{X}_i(g_t) \circ db_t^i, \quad \text{with } g_0 = e.$$

Then, $\{g_t\}_{t \geq 0}$ is a Brownian motion on G started at the identity, and, for all $t > 0$, $\text{Law}(g_t) = p_t dy$.

For $f, g \in C^\infty(G)$, we define the standard differential forms

$$\begin{aligned} \Gamma(f, g) &:= \frac{1}{2}(L(fg) - fLG - gLf) = \sum_{j=1}^n \left(\tilde{X}_j f \right) \left(\tilde{X}_j g \right) \text{ and} \\ \Gamma_2(f, g) &:= \frac{1}{2} (L\Gamma(f, g) - \Gamma(f, Lg) - \Gamma(g, Lf)). \end{aligned}$$

We follow the usual notational convention that $\Gamma(f) := \Gamma(f, f)$ and $\Gamma_2(f) := \Gamma_2(f, f)$. We will also need to consider derivatives in non-horizontal, or *vertical*, directions. Let $\mathcal{V} \subset \mathfrak{g}$ denote the Lie subalgebra such that

$$\mathfrak{g} = \mathcal{H} \oplus \mathcal{V},$$

and let $\{Z_\ell\}_{\ell=1}^d$ denote an orthonormal basis of \mathcal{V} , where $d = \dim(\mathcal{V})$, and following [2] and [4] we define

$$\begin{aligned} \Gamma^Z(f, g) &:= \sum_{\ell=1}^d \left(\tilde{Z}_\ell f \right) \left(\tilde{Z}_\ell g \right) \text{ and} \\ \Gamma_2^Z(f, g) &:= \frac{1}{2} (L\Gamma^Z(f, g) - \Gamma^Z(f, Lg) - \Gamma^Z(g, Lf)). \end{aligned}$$

Again we will let $\Gamma^Z(f) := \Gamma^Z(f, f)$ and $\Gamma_2^Z(f) := \Gamma_2^Z(f, f)$.

2.1. Reverse inequalities under curvature bound assumptions. In this section we consider some implications of certain curvature bound assumptions on G . These curvature bounds first appear in [4] and serve as generalizations of the standard curvature-dimension inequalities which appear in the elliptic literature. The assumption may be stated as follows: Suppose that there exist $\alpha, \beta > 0$ such that, for any $\nu > 0$ and $f \in C^\infty(G)$,

$$(2.2) \quad \Gamma_2(f) + \nu \Gamma_2^Z(f) \geq \alpha \Gamma^Z(f) - \frac{\beta}{\nu} \Gamma(f).$$

Before proceeding with the primary results of this section, we prove the following lemma which will be helpful in the sequel.

Lemma 2.2. *Suppose $\phi : J \rightarrow \mathbb{R}$ is a smooth function on an open interval $J \subset \mathbb{R}$ and $f : G \rightarrow \mathbb{R}$ is a measurable function. Fix $T > 0$ and set*

$$\Sigma(t, x) = P_t(\phi(P_{T-t}f))(x),$$

for all $t \in [0, T]$ and $x \in G$ (assuming $P_{T-t}f(G) \subset J$). Then

$$\frac{d\Sigma}{dt} = \frac{1}{2} P_t(\phi''(P_{T-t}f)\Gamma(P_{T-t}f))(x).$$

Proof. For simplicity, set $u_t = P_{T-t}f$. Then we just compute

$$(2.3) \quad \frac{d\Sigma}{dt} = P_t \left(L\phi(u_t)/2 + \frac{d}{dt}\phi(u_t) \right) = \frac{1}{2}P_t (L\phi(u_t) - \phi'(u_t)Lu_t).$$

Note that

$$\begin{aligned} L\phi(u_t) &= \sum_{i=1}^n \tilde{X}_i^2 \phi(u_t) = \sum_{i=1}^n \tilde{X}_i \left(\phi'(u_t)(\tilde{X}_i u_t) \right) \\ &= \sum_{i=1}^n \left(\phi''(u_t)(\tilde{X}_i u_t)^2 + \phi'(u_t)(\tilde{X}_i^2 u_t) \right) = \phi''(u_t)\Gamma(u_t) + \phi'(u_t)Lu_t. \end{aligned}$$

Then combining this with (2.3) yields the desired result. \square

We now prove that, assuming the curvature bound stated above, a reverse Poincaré inequality holds on G .

Notation 2.3. Let \mathcal{C} denote the set of functions $f : G \rightarrow \mathbb{R}$ such that $f \in C^\infty(G) \cap L^\infty(G)$ and $f, \sqrt{\Gamma(f)}, \sqrt{\Gamma^Z(f)} \in L^2(G)$. Note for example that \mathcal{C} includes all smooth functions with compact support. We will also let \mathcal{C}^+ denote functions f such that $f = g + \varepsilon$ for some $g \in \mathcal{C}$ with $g \geq 0$ and $\varepsilon > 0$.

Remark 2.4. It is shown in [4] that these function spaces are stable under P_t ; that is, if $f \in \mathcal{C}$ then $P_t f \in \mathcal{C}$ for all $t > 0$, and similarly for \mathcal{C}^+ .

Proposition 2.5. Assume that a curvature bound is satisfied as stated above in (2.2). Then, for all $T > 0$ and $f \in \mathcal{C}$,

$$\Gamma(P_T f) + \alpha T \Gamma^Z(P_T f) \leq \frac{1 + \frac{2\beta}{\alpha}}{T} (P_T(f^2) - (P_T f)^2).$$

In particular, this implies that

$$\Gamma(P_T f) \leq \frac{1 + \frac{2\beta}{\alpha}}{T} (P_T(f^2) - (P_T f)^2).$$

Proof. For $t \in [0, T]$, define the functional

$$\Phi(t) = a(t)P_t(\Gamma(P_{T-t}f)) + b(t)P_t(\Gamma^Z(P_{T-t}f)),$$

where $a, b : [0, T] \rightarrow [0, \infty)$ are control functions yet to be chosen. A straightforward computation shows that

$$\begin{aligned} \Phi'(t) &= a'(t)P_t(\Gamma(P_{T-t}f)) + b'(t)P_t(\Gamma^Z(P_{T-t}f)) \\ &\quad + 2a(t)P_t(\Gamma_2(P_{T-t}f)) + 2b(t)P_t(\Gamma_2^Z(P_{T-t}f)), \end{aligned}$$

and the inequality (2.2) implies that

$$\Gamma_2(P_{T-t}f) + \frac{b(t)}{a(t)}\Gamma_2^Z(P_{T-t}f) \geq -\frac{\beta a(t)}{b(t)}\Gamma(P_{T-t}f) + \alpha\Gamma^Z(P_{T-t}f).$$

Thus,

$$(2.4) \quad \Phi' \geq \left(a' - 2\beta \frac{a^2}{b} \right) P_t(\Gamma(P_{T-t}f)) + (b' + 2\alpha a) P_t(\Gamma^Z(P_{T-t}f)).$$

We now choose the functions a and b so that

$$b' + 2\alpha a = 0$$

and

$$a' - 2\beta \frac{a^2}{b} = C,$$

where C is a constant independent of t . This leads to the candidates

$$a(t) = \frac{1}{\alpha}(T - t)$$

and

$$b(t) = (T - t)^2.$$

For this choice of a and b , the inequality (2.4) becomes

$$(2.5) \quad \Phi'(t) \geq -\frac{1}{\alpha} \left(1 + \frac{2\beta}{\alpha} \right) P_t(\Gamma(P_{T-t}f)).$$

By Lemma 2.2 with $\phi(x) = x^2$, we have that

$$\frac{d}{dt} P_t(P_{T-t}f)^2 = P_t(\Gamma(P_{T-t}f)),$$

and thus integrating (2.5) from 0 to T yields the desired result. \square

Under an additional assumption, the curvature bound (2.2) also implies a reverse log Sobolev type inequality which is much stronger than the previous reverse Poincaré inequality. The additional required assumption here is the following commutation: for any $f \in C^\infty(G)$,

$$(2.6) \quad \Gamma(f, \Gamma^Z(f)) = \Gamma^Z(f, \Gamma(f)).$$

First we prove the following lemma given this assumption.

Lemma 2.6. *For fixed $T > 0$, $x \in G$, and $f \in \mathcal{C}$, define the entropy functionals*

$$\Phi_1(t) = P_t((P_{T-t}f)\Gamma(\ln P_{T-t}f))(x)$$

and

$$\Phi_2(t) = P_t((P_{T-t}f)\Gamma^Z(\ln P_{T-t}f))(x),$$

for $t \in [0, T]$. Then, assuming (2.6) holds,

$$\Phi_1'(t) = 2P_t((P_{T-t}f)\Gamma_2(\ln P_{T-t}f))(x)$$

and

$$\Phi_2'(t) = 2P_t((P_{T-t}f)\Gamma_2^Z(\ln P_{T-t}f))(x).$$

Proof. For $t \in [0, T]$ and $x \in G$, consider the functionals

$$\phi_1(t, x) = (P_{T-t}f)(x)\Gamma(\ln P_{T-t}f)(x)$$

and

$$\phi_2(t, x) = (P_{T-t}f)(x)\Gamma^Z(\ln P_{T-t}f)(x).$$

If we prove that

$$L\phi_1 + \frac{\partial\phi_1}{\partial t} = 2(P_{T-t}f)\Gamma_2(\ln P_{T-t}f)$$

and

$$L\phi_2 + \frac{\partial\phi_2}{\partial t} = 2(P_{T-t}f)\Gamma_2^Z(\ln P_{T-t}f),$$

then the result follows almost immediately.

Again for simplicity, take $u(t, x) = P_{T-t}f(x)$ and here let $u_t = \frac{\partial u}{\partial t}$. Then a simple computation gives

$$\frac{\partial\phi_1}{\partial t} = u_t\Gamma(\ln u) + 2u\Gamma\left(\ln u, \frac{u_t}{u}\right).$$

On the other hand,

$$L\phi_1 = Lu\Gamma(\ln u) + uL\Gamma(\ln u) + 2\Gamma(u, \Gamma(\ln u)).$$

Combining these equations we obtain

$$L\phi_1 + \frac{\partial\phi_1}{\partial t} = uL\Gamma(\ln u) + 2\Gamma(u, \Gamma(\ln u)) + 2u\Gamma\left(\ln u, \frac{u_t}{u}\right).$$

We now see that

$$\begin{aligned} 2u\Gamma_2(\ln u) &= u(L\Gamma(\ln u) - 2\Gamma(\ln u, L(\ln u))) \\ &= uL\Gamma(\ln u) - 2u\Gamma(\ln u, L(\ln u)). \end{aligned}$$

Observing that

$$L(\ln u) = -\frac{\Gamma(u)}{u^2} - \frac{u_t}{u},$$

we may conclude

$$L\phi_1 + \frac{\partial\phi_1}{\partial t} = 2(P_{T-t}f)\Gamma_2(\ln P_{T-t}f).$$

In the same vein, we obtain

$$L\phi_2 + \frac{\partial\phi_2}{\partial t} = uL\Gamma^Z(\ln u) + 2\Gamma(u, \Gamma^Z(\ln u)) + 2u\Gamma^Z\left(\ln u, \frac{u_t}{u}\right).$$

This time using the definition of Γ_2^Z , we find

$$\begin{aligned} 2u\Gamma_2^Z(\ln u) &= u(L\Gamma^Z(\ln u) - 2\Gamma^Z(\ln u, L(\ln u))) \\ &= uL\Gamma^Z(\ln u) + 2u\Gamma^Z\left(\ln u, \frac{\Gamma(u)}{u^2}\right) + 2u\Gamma^Z\left(\ln u, \frac{u_t}{u}\right). \end{aligned}$$

From this last equation it is now clear that under the assumption (2.6) we have

$$L\phi_2 + \frac{\partial\phi_2}{\partial t} = 2u\Gamma_2^Z(\ln u),$$

and this concludes the proof. \square

Given this lemma, we may now prove the following reverse log Sobolev inequalities holds.

Theorem 2.7. *Suppose that (2.2) and (2.6) are satisfied. Then, for any $T > 0$ and $f \in \mathcal{C}^+$,*

$$\Gamma(\ln P_T f) + \alpha T \Gamma^Z(\ln P_T f) \leq \frac{1 + \frac{2\beta}{\alpha}}{T} \left(\frac{P_T(f \ln f)}{P_T f} - \ln P_T f \right).$$

In particular, the following reverse log Sobolev inequality holds

$$(2.7) \quad \Gamma(\ln P_T f) \leq \frac{1 + \frac{2\beta}{\alpha}}{T} \left(\frac{P_T(f \ln f)}{P_T f} - \ln P_T f \right).$$

Proof. For $t \in [0, T]$, we define the functional

$$\Psi(t) = a(t)\Phi_1(t) + b(t)\Phi_2(t),$$

where Φ_1 and Φ_2 are as defined in Lemma 2.6 and a, b are non-negative control functions to be chosen later. Since we assume (2.6) holds, Lemma 2.6 implies that

$$\begin{aligned} \Psi'(t) &= a'(t)\Phi_1(t) + b'(t)\Phi_2(t) \\ &\quad + 2a(t)P_t((P_{T-t}f)\Gamma_2(\ln P_{T-t}f)) + 2b(t)P_t((P_{T-t}f)\Gamma_2^Z(\ln P_{T-t}f)). \end{aligned}$$

Thus, given the curvature bound and working exactly as in Proposition 2.5, we are lead to the same choices

$$a(t) = \frac{1}{\alpha}(T-t) \quad \text{and} \quad b(t) = (T-t)^2.$$

For this choice of a and b , we have the inequality

$$(2.8) \quad \Psi'(t) \geq -\frac{1}{\alpha} \left(1 + \frac{2\beta}{\alpha} \right) \Phi_1(t).$$

Taking $\phi(x) = x \ln x$ in Lemma 2.2 implies that

$$\begin{aligned} \frac{d}{dt} P_t((P_{T-t}f)(\ln P_{T-t}f)) &= P_t \left(\frac{\Gamma(P_{T-t}f)}{P_{T-t}f} \right) \\ &= P_t((P_{T-t}f)\Gamma(\ln P_{T-t}f)) = \Phi_1(t), \end{aligned}$$

and thus integrating (2.8) from 0 to T then yields the claimed result. \square

Remark 2.8. Other choices of control functions could be made to satisfy the desired criteria in the proofs of Proposition 2.5 and Theorem 2.7. In particular, we could have taken

$$a(t) = \frac{1}{\alpha}((1+\delta)T-t)$$

and

$$b(t) = ((1+\delta)T-t)^2,$$

for any $\delta \geq 0$. This choice of a and b would give the following generalized estimates. The following would generalize Proposition 2.5: for all $T > 0$ and $f \in L^\infty(G)$,

$$\begin{aligned} & \Gamma(P_T f) + \alpha(1 + \delta)T\Gamma^Z(P_T f) \\ & \leq \frac{1 + \frac{2\beta}{\alpha}}{(1 + \delta)T}(P_T(f^2) - (P_T f)^2) + \frac{\delta}{1 + \delta}P_T(\Gamma(f)) + \frac{\alpha\delta^2}{1 + \delta}TP_T(\Gamma^Z(f)). \end{aligned}$$

Also, the following statement would generalize Theorem 2.7: for all $T > 0$ and $f \in \mathcal{C}$,

$$\begin{aligned} & P_T f \Gamma(\ln P_T f) + \alpha(1 + \delta)TP_T f \Gamma^Z(\ln P_T f) \\ & \leq \frac{1 + \frac{2\beta}{\alpha}}{(1 + \delta)T}(P_T(f \ln f) - (P_T f) \ln P_T f) \\ & \quad + \frac{\delta}{1 + \delta}P_T(f \Gamma(\ln f)) + \frac{\alpha\delta^2}{1 + \delta}TP_T(f \Gamma^Z(\ln f)). \end{aligned}$$

2.2. Wang type and integrated Harnack inequalities. A reverse log Sobolev inequality such as in Theorem 2.7 is sufficient to prove an analogue of Wang's dimension-free Harnack inequality. Estimates of this type were first proved by Wang in a Riemannian setting under the assumption of a lower bound on the Ricci curvature [27]. Before stating the estimate, we must make the following definition.

Notation 2.9. (*Horizontal distance*)

(1) The length of a C^1 -path $\sigma : [a, b] \rightarrow G$ is defined as

$$\ell(\sigma) := \int_a^b |L_{\sigma^{-1}(s)*}\dot{\sigma}(s)|_{\mathfrak{g}} ds.$$

(2) A C^1 -path $\sigma : [a, b] \rightarrow G$ is horizontal if $L_{\sigma(t)^{-1}*}\dot{\sigma}(t) \in \mathcal{H} \times \{0\}$ for a.e. t . Let $C^{1,h}$ denote the set of horizontal paths $\sigma : [0, 1] \rightarrow G$.

(3) The horizontal distance between $x, y \in G$ is defined by

$$d(x, y) := \inf\{\ell(\sigma) : \sigma \in C^{1,h} \text{ such that } \sigma(0) = x \text{ and } \sigma(1) = y\}.$$

Proposition 2.10. Suppose there exists a constant $C < \infty$ such that, for all $T > 0$ and $f \in \mathcal{C}^+$,

$$(2.9) \quad \Gamma(\ln P_T f) \leq \frac{C}{T} \left(\frac{P_T(f \ln f)}{P_T f} - \ln P_T f \right).$$

Then, for all $T > 0$, $x, y \in G$, $f \in L^\infty(G)$ with $f \geq 0$, and $p \in (1, \infty)$,

$$(2.10) \quad (P_T f)^p(x) \leq P_T f^p(y) \exp \left(C \frac{d^2(x, y)}{4(p-1)T} \right).$$

Proof. First take $f \in \mathcal{C}^+$. Let $b(s) = 1 + (p-1)s$ for $s \in [0, 1]$ and $\sigma : [0, 1] \rightarrow G$ be an arbitrary horizontal path such that $\sigma(0) = x$ and $\sigma(1) = y$. Define the functional

$$\phi(s) = \frac{p}{b(s)} \ln P_T f^{b(s)}(\sigma(s)), \quad \text{for } s \in [0, 1].$$

Since

$$\frac{d}{ds} f^{b(s)}(\sigma(s)) = f^{b(s)}(\sigma(s)) \left((p-1) \ln f(\sigma(s)) + b(s) \langle d(\ln f)(\sigma(s)), \sigma'(s) \rangle \right),$$

differentiating ϕ with respect to s and applying (2.9) yields

$$\begin{aligned} \phi'(s) &= -\frac{p(p-1)}{b(s)^2} \ln P_T f^{b(s)} \\ &\quad + \frac{p}{b(s)} \left(\frac{(p-1)P_T(f^{b(s)} \ln f) + b(s)P_T(f^{b(s)} \langle d(\ln f), \sigma' \rangle)}{P_T f^{b(s)}} \right) \\ &= \frac{p(p-1)}{b(s)^2} \left(\frac{P_T(f^{b(s)} \ln f^{b(s)})}{P_T f^{b(s)}} - \ln P_T f^{b(s)} \right) \\ &\quad + \frac{p}{b(s)} \langle d(\ln P_T f^{b(s)}), \sigma'(s) \rangle \\ &\geq \frac{p(p-1)T}{b(s)^2 C} \Gamma(\ln P_T f^{b(s)}) + \frac{p}{b(s)} \langle d(\ln P_T f^{b(s)}), \sigma'(s) \rangle. \end{aligned}$$

Now, for every $\lambda > 0$,

$$\langle d(\ln P_T f^{b(s)}), \sigma'(s) \rangle \geq -\frac{1}{2\lambda} \Gamma(\ln P_T f^{b(s)}) - \frac{\lambda}{2} |\sigma'(s)|^2,$$

since σ horizontal implies that $\sigma'(s) \in \text{span}\{\tilde{X}_i(\sigma(s))\}_{i=1}^n$. In particular, choosing

$$\lambda = \frac{C}{2p(p-1)T} b(s)^2$$

gives

$$\phi'(s) \geq -\frac{C}{4p(p-1)T} (1 + (p-1)s)^2 |\sigma'(s)|^2.$$

Integrating this inequality from 0 to 1 yields

$$\ln P_T f^p(y) - \ln(P_T f)^p(x) \geq -\frac{C}{4p(p-1)T} \int_0^1 (1 + (p-1)s)^2 |\sigma'(s)|^2 ds.$$

Maximizing $\int_0^1 (1 + (p-1)s)^2 |\sigma'(s)|^2 ds$ over the set of horizontal paths such that $\sigma(0) = x$ and $\sigma(1) = y$ shows that (2.10) holds for $f \in \mathcal{C} \cap C^\infty(G)$.

To prove the estimate for general $f \geq 0$, let $C^\infty(G) \ni h_n \geq 0$ be an increasing sequence of functions with compact support such that $h_n \uparrow 1$. Then (2.10) holds for $g = h_n P_\tau f + \varepsilon \in \mathcal{C} \cap C^\infty(G)$ for all $n, \tau > 0$, and $\varepsilon > 0$. Then letting $\varepsilon \rightarrow 0, \tau \rightarrow 0$, and $n \rightarrow \infty$ in the inequality completes the proof of (2.10) for $f \geq 0$. \square

Wang type Harnack inequalities in turn are equivalent to so-called “integrated Harnack inequalities” as in the following lemma. Here we follow the proof of 2.4 in [29]. An alternative form and proof of the following equivalence can be found in Lemma D.1 of [12].

Lemma 2.11. *Let $T > 0$, $x, y \in G$, $p \in (1, \infty)$, and $C \in (0, \infty]$. Then*

$$(2.11) \quad (P_T f)^p(x) \leq C P_T f^p(y), \quad \text{for all } f \in L^\infty(G) \text{ with } f \geq 0,$$

if and only if

$$(2.12) \quad \left(\int_G \left[\frac{p_T(x, z)}{p_T(y, z)} \right]^{1/(p-1)} p_T(y, z) dz \right)^{p-1} \leq C.$$

Proof. Set $J_{x,y}(z) = \frac{p_T(x,z)}{p_T(y,z)}$ and $f_n := (n \wedge J_{x,y})^{1/(p-1)}$ for $n \geq 1$. Then applying (2.11) to f_n yields

$$\begin{aligned} (P_T f_n)^p(x) &\leq C P_T f_n^p(y) = C \int_G (n \wedge J_{x,y}(z))^{p/(p-1)} p_T(y, z) dz \\ &\leq C \int_G \left(n \wedge \frac{p_T(x, z)}{p_T(y, z)} \right)^{1/(p-1)} p_T(x, z) dz = C P_T f_n(x). \end{aligned}$$

Thus,

$$P_T J_{x,y}^{1/(p-1)}(x) = \lim_{n \rightarrow \infty} P_T f_n(x) \leq C^{1/(p-1)},$$

which yields (2.12).

For the converse, we have by Hölder’s inequality that

$$\begin{aligned} P_T f(x) &= \int_G f(z) \frac{p_T(x, z)}{p_T(y, z)} p_T(y, z) dz \\ &\leq (P_T f^p)^{1/p}(y) \left(\int_G \frac{p_T(x, z)^{p/(p-1)}}{p_T(y, z)} p_T(y, z) dz \right)^{(p-1)/p} \\ &= (P_T f^p)^{1/p}(y) \left(\int_G \frac{p_T(x, z)^{1/(p-1)}}{p_T(y, z)} p_T(x, z) dz \right)^{(p-1)/p} \\ &\leq (P_T f^p)^{1/p}(y) C^{1/p}, \end{aligned}$$

which completes the proof. \square

3. INFINITE-DIMENSIONAL HEISENBERG-LIKE GROUPS

In this section, we review the definitions for infinite-dimensional Heisenberg-like groups, which are infinite-dimensional Lie groups based on an abstract Wiener space. Much of the material in this section also appears in [15].

3.1. Abstract Wiener spaces. For the reader's convenience, we summarize several well-known properties of Gaussian measures and abstract Wiener spaces that are required for the sequel. These results as well as more details on abstract Wiener spaces and some particular examples may be found in [6, 20].

Suppose that W is a real separable Banach space and \mathcal{B}_W is the Borel σ -algebra on W .

Definition 3.1. A measure μ on (W, \mathcal{B}_W) is called a (mean zero, non-degenerate) *Gaussian measure* provided that its characteristic functional is given by

$$(3.1) \quad \hat{\mu}(u) := \int_W e^{iu(x)} d\mu(x) = e^{-\frac{1}{2}q(u,u)}, \quad \text{for all } u \in W^*,$$

for $q = q_\mu : W^* \times W^* \rightarrow \mathbb{R}$ a symmetric, positive definite quadratic form. That is, q is a real inner product on W^* .

Lemma 3.2. *If $u, v \in W^*$, then*

$$\int_W u(w)v(w) d\mu(w) = q(u, v).$$

Proof. Let $u_*\mu := \mu \circ u^{-1}$ denote the measure on \mathbb{R} which is the push forward of μ under u . Then by equation (3.1) $u_*\mu$ is normal with mean 0 and variance $q(u, u)$. Thus,

$$\int_W u^2(w) d\mu(w) = q(u, u).$$

Polarizing this identity gives the desired result. \square

A proof of the following standard theorem may be found for example in Appendix A of [11].

Theorem 3.3. *Let μ be a Gaussian measure on a real separable Banach space W . For $p \in [1, \infty)$, let*

$$(3.2) \quad C_p := \int_W \|w\|_W^p d\mu(w).$$

For $w \in W$, let

$$\|w\|_H := \sup_{u \in W^* \setminus \{0\}} \frac{|u(w)|}{\sqrt{q(u, u)}}$$

and define the Cameron-Martin subspace $H \subset W$ by

$$H := \{h \in W : \|h\|_H < \infty\}.$$

Then

- (1) For all $p \in [1, \infty)$, $C_p < \infty$.
- (2) H is a dense subspace of W .
- (3) There exists a unique inner product $\langle \cdot, \cdot \rangle_H$ on H such that $\|h\|_H^2 = \langle h, h \rangle_H$ for all $h \in H$, and H is a separable Hilbert space with respect to this inner product.

- (4) For any $h \in H$, $\|h\|_W \leq \sqrt{C_2} \|h\|_H$.
 (5) If $\{e_j\}_{j=1}^\infty$ is an orthonormal basis for H , then for any $u, v \in H^*$

$$q(u, v) = \langle u, v \rangle_{H^*} = \sum_{j=1}^{\infty} u(e_j)v(e_j).$$

It follows from item (4) that any $u \in W^*$ restricted to H is in H^* . Therefore, by item (5) and Lemma 3.2,

$$(3.3) \quad \int_W u^2(w) d\mu(w) = q(u, u) = \|u\|_{H^*}^2 = \sum_{j=1}^{\infty} |u(e_j)|^2.$$

More generally we have the following lemma.

Lemma 3.4. *Let K be a real Hilbert space and $\varphi : W \rightarrow K$ be a linear map. Then*

$$\|\varphi\|_{H^* \otimes K}^2 = \sum_{j=1}^{\infty} \|\varphi(e_j)\|_K^2 = \int_W \|\varphi(w)\|_K^2 d\mu(w).$$

Proof. Let $\{f_\ell\}_{\ell=1}^{\dim(K)}$ be an orthonormal basis of K . Then by equation (3.3)

$$\begin{aligned} \int_W \|\varphi(w)\|_K^2 d\mu(w) &= \int_W \sum_{\ell=1}^{\dim(K)} |\langle \varphi(w), f_\ell \rangle_K|^2 d\mu(w) \\ &= \sum_{\ell=1}^{\dim(K)} \|\langle \varphi(\cdot), f_\ell \rangle_K\|_{H^*}^2 = \sum_{\ell=1}^{\dim(K)} \sum_{j=1}^{\infty} |\langle \varphi(e_j), f_\ell \rangle_K|^2 \\ &= \sum_{j=1}^{\infty} \|\varphi(e_j)\|_K^2 = \|\varphi\|_{H^* \otimes K}^2. \end{aligned}$$

□

This leads to the following facts for linear maps on W . First we set the some notation.

Notation 3.5. *Let K be a real Hilbert space, and suppose $\alpha : H^{\otimes m} \rightarrow K$ is a multi-linear map. Then the Hilbert-Schmidt norm of α is defined by*

$$\begin{aligned} \|\alpha\|_2^2 &:= \|\alpha\|_{(H^*)^{\otimes m} \otimes K}^2 = \sum_{j_1, \dots, j_m=1}^{\infty} \|\alpha(e_{j_1}, \dots, e_{j_m})\|_K^2 \\ &= \sum_{j_1, \dots, j_m=1}^{\infty} \sum_{\ell=1}^{\dim(K)} \langle \alpha(e_{j_1}, \dots, e_{j_m}), f_\ell \rangle_K^2, \end{aligned}$$

where $\{e_j\}_{j=1}^\infty$ and $\{f_\ell\}_{\ell=1}^{\dim(K)}$ are orthonormal bases of H and K , respectively.

One may verify directly that these norms are independent of the chosen bases.

Lemma 3.6. *Suppose K is a Hilbert space and $\varphi : W \rightarrow K$ is a continuous linear map. Then $\varphi : H \rightarrow K$ is Hilbert-Schmidt, that is, $\|\varphi\|_2 < \infty$.*

Proof. By Lemma 3.4,

$$\begin{aligned} \|\varphi\|_2^2 &= \|\varphi\|_{H^* \otimes K}^2 = \int_W \|\varphi(w)\|_K^2 d\mu(w) \\ &\leq \|\varphi\|_0^2 \int_W \|w\|_W^2 d\mu(w) = C_2 \|\varphi\|_0^2, \end{aligned}$$

where $C_2 < \infty$ is as defined in (3.2) and

$$\|\varphi\|_0 := \sup\{\|\varphi(w)\|_K : \|w\|_W = 1\} < \infty$$

by the continuity of φ . \square

Similarly, we may prove the following.

Lemma 3.7. *Suppose K is a Hilbert space and $\rho : W \times W \rightarrow K$ is a continuous bilinear map. Then $\rho : H \times H \rightarrow K$ is Hilbert-Schmidt.*

Proof. Note first that, for each $w \in W$, $\varphi = \rho(w, \cdot)$ is a continuous linear operator and thus, by the proof Lemma 3.6,

$$\|\rho(w, \cdot)\|_2^2 = \|\rho(w, \cdot)\|_{H^* \otimes K}^2 \leq C_2 \|\rho(w, \cdot)\|_0^2 \leq C_2 \|\rho\|_0 \|w\|_W^2,$$

where

$$\|\rho\|_0 := \sup\{\|\rho(w, w')\|_K : \|w\|_W = \|w'\|_W = 1\} < \infty.$$

Then viewing $w \mapsto \rho(w, \cdot)$ as a continuous linear map from W to the Hilbert space $H^* \otimes K$, Lemma 3.4 implies that

$$\begin{aligned} \|\rho\|_2^2 &= \|\rho(w, \cdot)\|_{H^* \otimes (H^* \otimes K)}^2 = \int_W \|\rho(w, \cdot)\|_{H^* \otimes K}^2 d\mu(w) \\ &\leq \int_W C_2 \|\rho\|_0^2 \|w\|_W^2 d\mu(w) = C_2^2 \|\rho\|_0^2 < \infty. \end{aligned}$$

\square

3.2. Infinite-dimensional Heisenberg-like groups. We revisit the definition of the infinite-dimensional Heisenberg-like groups that were first considered in [11]. Note that since we are interested in subelliptic heat kernel measures on these groups, there are some necessary modifications to the topology as was done in [15]. First we set the following notation which will hold for the rest of the paper.

Notation 3.8. *Let (W, H, μ) be a real abstract Wiener space. Let \mathbf{C} be a real Hilbert space with inner product $\langle \cdot, \cdot \rangle_{\mathbf{C}}$ and $\dim(\mathbf{C}) = d < \infty$. Let $\omega : W \times W \rightarrow \mathbf{C}$ be a continuous skew-symmetric bilinear form on W . We will also trivially assume that ω is surjective (otherwise, we just restrict to a linear subspace of \mathbf{C}).*

Definition 3.9. Let \mathfrak{g} denote $W \times \mathbf{C}$ when thought of as a Lie algebra with the Lie bracket given by

$$(3.4) \quad [(X_1, V_1), (X_2, V_2)] := (0, \omega(X_1, X_2)).$$

Let G denote $W \times \mathbf{C}$ when thought of as a group with multiplication given by

$$(3.5) \quad g_1 g_2 := g_1 + g_2 + \frac{1}{2}[g_1, g_2],$$

where g_1 and g_2 are viewed as elements of \mathfrak{g} . For $g_i = (w_i, c_i)$, this may be written equivalently as

$$(3.6) \quad (w_1, c_1) \cdot (w_2, c_2) = \left(w_1 + w_2, c_1 + c_2 + \frac{1}{2}\omega(w_1, w_2) \right).$$

We will call G constructed in this way a *Heisenberg-like group*.

It is easy to verify that, given this bracket and multiplication, \mathfrak{g} is indeed a Lie algebra and G is a group. Note that $g^{-1} = -g$ and the identity $e = (0, 0)$.

Notation 3.10. Let \mathfrak{g}_{CM} denote $H \times \mathbf{C}$ when thought of as a Lie subalgebra of \mathfrak{g} , and we will refer to \mathfrak{g}_{CM} as the Cameron-Martin subalgebra of \mathfrak{g} . Similarly, let G_{CM} denote $H \times \mathbf{C}$ when thought of as a subgroup of G , and we will refer to G_{CM} as the Cameron-Martin subgroup of G .

We will equip $\mathfrak{g} = G$ with the homogeneous norm

$$\|(w, c)\|_{\mathfrak{g}} := \sqrt{\|w\|_W^2 + \|c\|_{\mathbf{C}}},$$

and analogously on $\mathfrak{g}_{CM} = G_{CM}$ we define

$$\|(A, a)\|_{\mathfrak{g}_{CM}} := \sqrt{\|A\|_H^2 + \|a\|_{\mathbf{C}}}.$$

One may easily see that G and G_{CM} are topological groups with respect to the topologies induced by the homogeneous norms, see for example Lemma 2.9 of [15].

Before proceeding, let us give the basic motivating examples for the construction of these infinite-dimensional Heisenberg-like groups. In what follows, if X is a complex vector space, let X_{Re} denote X thought of as a real vector space. If $(H, \langle \cdot, \cdot \rangle_H)$ is a complex Hilbert space, let $\langle \cdot, \cdot \rangle_{H_{\text{Re}}} := \text{Re} \langle \cdot, \cdot \rangle_H$, in which case $(H_{\text{Re}}, \langle \cdot, \cdot \rangle_{H_{\text{Re}}})$ becomes a real Hilbert space.

Example 3.1 (Finite-dimensional Heisenberg group). Let $W = H = (\mathbb{C}^n)_{\text{Re}} \cong \mathbb{R}^{2n}$ and μ be standard Gaussian measure on \mathbb{R}^{2n} . Then (W, H, μ) is an abstract Wiener space. Let $\mathbf{C} = \mathbb{R}$ and $\omega(w, z) := \text{Im} \langle w, z \rangle$, where $\langle w, z \rangle = w \cdot \bar{z}$ is the usual inner product on \mathbb{C}^n . Then $G = \mathbb{R}^{2n} \times \mathbb{R}$ equipped with a group operation as defined in (3.6) is a finite-dimensional Heisenberg group.

Example 3.2 (Heisenberg group of a symplectic vector space). Let $(K, \langle \cdot, \cdot \rangle)$ be a complex Hilbert space and Q be a strictly positive trace class operator on K . For $h, k \in K$, let $\langle h, k \rangle_Q := \langle h, Qk \rangle$ and $\|h\|_Q := \sqrt{\langle h, h \rangle_Q}$, and let $(K_Q, \langle \cdot, \cdot \rangle_Q)$ denote the Hilbert space completion of $(K, \|\cdot\|_Q)$. Then $W = (K_Q)_{\text{Re}}$ and $H = K_{\text{Re}}$ determines an abstract Wiener space (see, for example, exercise 17 on p.59 of [20]). Letting $\mathbf{C} = \mathbb{R}$ and

$$\omega(w, z) := \text{Im}\langle w, z \rangle_Q,$$

then $G = (K_Q)_{\text{Re}} \times \mathbb{R}$ equipped with a group operation as defined in (3.6) is an infinite-dimensional Heisenberg-like group.

3.3. Finite-dimensional projection groups. The finite-dimensional projections of G defined in this section will be important in the sequel. Note that the construction of these projections is quite natural in the sense that they come from the usual projections of the abstract Wiener space; however, the projections defined here are not group homomorphisms, which is a complicating factor in the analysis.

As usual, let (W, H, μ) denote an abstract Wiener space. Let $i : H \rightarrow W$ be the inclusion map, and $i^* : W^* \rightarrow H^*$ be its transpose so that $i^*\ell := \ell \circ i$ for all $\ell \in W^*$. Also, let

$$H_* := \{h \in H : \langle \cdot, h \rangle_H \in \text{Range}(i^*) \subset H^*\}.$$

That is, for $h \in H$, $h \in H_*$ if and only if $\langle \cdot, h \rangle_H \in H^*$ extends to a continuous linear functional on W , which we will continue to denote by $\langle \cdot, h \rangle_H$. Because H is a dense subspace of W , i^* is injective and thus has a dense range. Since $H \ni h \mapsto \langle \cdot, h \rangle_H \in H^*$ is a linear isometric isomorphism, it follows that $H_* \ni h \mapsto \langle \cdot, h \rangle_H \in W^*$ is a linear isomorphism also, and so H_* is a dense subspace of H .

Suppose that $P : H \rightarrow H$ is a finite rank orthogonal projection such that $PH \subset H_*$. Let $\{e_j\}_{j=1}^n$ be an orthonormal basis for PH . Then we may extend P to a (unique) continuous operator from $W \rightarrow H$ (still denoted by P) by letting

$$(3.7) \quad Pw := \sum_{j=1}^n \langle w, e_j \rangle_H e_j$$

for all $w \in W$.

Notation 3.11. Let $\text{Proj}(W)$ denote the collection of finite rank projections on W such that

- (1) $PW \subset H_*$,
- (2) $P|_H : H \rightarrow H$ is an orthogonal projection (that is, P has the form given in equation (3.7)), and
- (3) PW is sufficiently large to satisfy Hörmander's condition (that is, $\{\omega(A, B) : A, B \in PW\} = \mathbf{C}$).

For each $P \in \text{Proj}(W)$, we define $G_P := PW \times \mathbf{C} \subset H_* \times \mathbf{C}$ and a corresponding projection $\pi_P : G \rightarrow G_P$

$$\pi_P(w, x) := (Pw, x).$$

We will also let $\mathfrak{g}_P = \text{Lie}(G_P) = PW \times \mathbf{C}$. In the context of Section 2, note that, for each $P \in \text{Proj}(W)$, G_P is a finite-dimensional connected unimodular Lie group (in fact, \mathfrak{g}_P is step 2 stratified) with $\mathcal{H} = PH$ and $\mathcal{V} = \mathbf{C}$.

3.4. Derivatives and differential forms on G . For $x \in G$, again let $L_x : G \rightarrow G$ and $R_x : G \rightarrow G$ denote left and right multiplication by x , respectively. As G is a vector space, to each $x \in G$ we can associate the tangent space $T_x G$ to G at x , which is naturally isomorphic to G .

Notation 3.12 (Linear and group derivatives). *Let $f : G \rightarrow \mathbf{C}$ denote a Fréchet smooth function for G considered as a Banach space with respect to the norm*

$$|(w, c)|_G := \sqrt{\|w\|_W^2 + \|c\|_{\mathbf{C}}^2}.$$

Then, for $x \in G$, and $h, k \in \mathfrak{g}$, let

$$f'(x)h := \partial_h f(x) = \left. \frac{d}{dt} \right|_0 f(x + th)$$

and

$$f''(x)(h \otimes k) := \partial_h \partial_k f(x).$$

For $v, x \in G$, let $v_x \in T_x G$ denote the tangent vector satisfying $v_x f = f'(x)v$. If $x(t)$ is any smooth curve in G such that $x(0) = x$ and $\dot{x}(0) = v$ (for example, $x(t) = x + tv$), then

$$L_{g*} v_x = \left. \frac{d}{dt} \right|_0 g \cdot x(t).$$

In particular, for $x = e$ and $v_e = h \in \mathfrak{g}$, again we let $\tilde{h}(g) := L_{g*} h$, so that \tilde{h} is the unique left invariant vector field on G such that $\tilde{h}(e) = h$. As usual we view \tilde{h} as a first order differential operator acting on smooth functions by

$$(\tilde{h}f)(x) = \left. \frac{d}{dt} \right|_0 f(x \cdot \sigma(t)),$$

where $\sigma(t)$ is a smooth curve in G such that $\sigma(0) = e$ and $\dot{\sigma}(0) = h$ (for example, $\sigma(t) = th$).

Proposition 3.13. *Let $f : G \rightarrow \mathbb{R}$ be a smooth function, $h = (A, a) \in \mathfrak{g}$ and $x = (w, c) \in G$. Then*

$$\tilde{h}(x) := l_{x*} h = \left(A, a + \frac{1}{2} \omega(w, A) \right)_x \quad \text{for all } x = (w, c) \in G$$

and in particular

$$(3.8) \quad \widetilde{(A, a)}f(x) = f'(x) \left(A, a + \frac{1}{2}\omega(w, A) \right).$$

Furthermore, if $h = (A, a)$, $k = (B, b)$, then

$$(3.9) \quad (\tilde{h}\tilde{k}f - \tilde{k}\tilde{h}f) = \widetilde{[h, k]}f.$$

That is, the Lie algebra structure on \mathfrak{g} induced by the Lie algebra structure on the left invariant vector fields on G is the same as the Lie algebra structure defined in equation (3.4).

Proof. Since $th = t(A, a)$ is a curve in G passing through the identity at $t = 0$, we have

$$\begin{aligned} \tilde{h}(x) &= \frac{d}{dt} \Big|_0 [x \cdot (th)] = \frac{d}{dt} \Big|_0 [(w, c) \cdot t(A, a)] \\ &= \frac{d}{dt} \Big|_0 \left[\left(w + tA, c + ta + \frac{t}{2}\omega(w, A) \right) \right] \\ &= \left(A, a + \frac{1}{2}\omega(w, A) \right). \end{aligned}$$

So by the chain rule, $(\tilde{h}f)(x) = f'(x)\tilde{h}(x)$ and hence

$$(3.10) \quad \begin{aligned} (\tilde{h}\tilde{k}f)(x) &= \frac{d}{dt} \Big|_0 [f'(x \cdot th) \tilde{k}(x \cdot th)] \\ &= f''(x) \left(\tilde{h}(x) \otimes \tilde{k}(x) \right) + f''(x) \frac{d}{dt} \Big|_0 \tilde{k}(x \cdot th), \end{aligned}$$

where

$$\frac{d}{dt} \Big|_0 \tilde{k}(x \cdot th) = \frac{d}{dt} \Big|_0 \left(B, a + \frac{1}{2}\omega(w + tA, B) \right) = \left(0, \frac{1}{2}\omega(A, B) \right).$$

Since $f''(x)$ is symmetric, it now follows by subtracting equation (3.10) from itself with h and k interchanged that

$$\left(\tilde{h}\tilde{k}f - \tilde{k}\tilde{h}f \right) (x) = f'(x) (0, \omega(A, B)) = f'(x) [h, k] = \left(\widetilde{[h, k]}f \right) (x)$$

as desired. \square

Definition 3.14. A function $f : G \rightarrow \mathbb{C}$ is a (smooth) cylinder function if it may be written as $f = F \circ \pi_P$, for some $P \in \text{Proj}(W)$ and (smooth) $F : G_P \rightarrow \mathbb{C}$.

Lemma 3.15. Suppose that $\ell : H \rightarrow \mathbf{C}$ is a continuous linear map. Then for any orthonormal basis $\{e_j\}_{j=1}^\infty$ of H the series

$$(3.11) \quad \sum_{j=1}^{\infty} \ell(e_j) \otimes \ell(e_j) \in \mathbf{C} \otimes \mathbf{C}$$

and

$$(3.12) \quad \sum_{j=1}^{\infty} \ell(e_j) \otimes e_j \in \mathbf{C} \otimes H$$

are convergent and independent of the basis.

Proof. First note that by Lemma 3.6, ℓ is Hilbert-Schmidt, and thus

$$\sum_{j=1}^{\infty} \|\ell(e_j) \otimes \ell(e_j)\|_{\mathbf{C} \otimes \mathbf{C}} = \sum_{j=1}^{\infty} \|\ell(e_j)\|_{\mathbf{C}}^2 < \infty,$$

which shows that the sum in equation (3.11) is absolutely convergent. Similarly, since $\{\ell(e_j) \otimes e_j\}_{j=1}^{\infty}$ is an orthogonal set in $\mathbf{C} \otimes H$ and

$$\sum_{j=1}^{\infty} \|\ell(e_j) \otimes e_j\|_{\mathbf{C} \otimes H}^2 = \sum_{j=1}^{\infty} \|\ell(e_j)\|_{\mathbf{C}}^2 < \infty,$$

the sum in equation (3.12) is convergent as well.

Now recall that if H and K are two real Hilbert spaces then the Hilbert space tensor product $H \otimes K$ is unitarily equivalent to the space of Hilbert-Schmidt operators $HS(H, K)$ from H to K . Under this identification, $h \otimes k \in H \otimes K$ corresponds to the operator (still denoted by $h \otimes k$) in $HS(H, K)$ defined by

$$H \ni h' \mapsto (h \otimes k)h' = (h, h')_H k \in K.$$

Using this identification we have that, for all $c \in \mathbf{C}$,

$$\begin{aligned} \left(\sum_{j=1}^{\infty} \ell(e_j) \otimes \ell(e_j) \right) c &= \sum_{j=1}^{\infty} \ell(e_j) \langle \ell(e_j), c \rangle_{\mathbf{C}} = \sum_{j=1}^{\infty} \ell(e_j) \langle e_j, \ell^* c \rangle_{\mathbf{C}} \\ &= \ell \left(\sum_{j=1}^{\infty} \langle e_j, \ell^* c \rangle_{\mathbf{C}} e_j \right) = \ell \ell^* c \end{aligned}$$

and

$$\left(\sum_{j=1}^{\infty} \ell(e_j) \otimes e_j \right) c = \sum_{j=1}^{\infty} e_j \langle \ell(e_j), c \rangle_{\mathbf{C}} = \sum_{j=1}^{\infty} e_j \langle e_j, \ell^* c \rangle_{\mathbf{C}} = \ell^* c,$$

which clearly shows that equations (3.11) and (3.12) are basis-independent. \square

Notation 3.16. For $x = (w, c) \in G$, let $\gamma(x)$ and $\chi(x)$ be the elements of $\mathfrak{g}_{CM} \otimes \mathfrak{g}_{CM}$ defined by

$$\begin{aligned}\gamma(x) &:= \sum_{j=1}^{\infty} (0, \omega(w, e_j)) \otimes (e_j, 0) \text{ and} \\ \chi(x) &:= \sum_{j=1}^{\infty} (0, \omega(w, e_j)) \otimes (0, \omega(w, e_j))\end{aligned}$$

where $\{e_j\}_{j=1}^{\infty}$ is any orthonormal basis for H . Both γ and χ are well defined by Lemma 3.15 taking $\ell = \omega(x, \cdot)$.

The following proposition is proved in Proposition 3.29 of [11], although the statement given there is for the elliptic case. We reproduce the short proof here for the reader's convenience.

Proposition 3.17. Let $\{e_j\}_{j=1}^{\infty}$ be an orthonormal basis for H . Then, for any smooth cylinder function $f : G \rightarrow \mathbb{R}$,

$$Lf(x) := \sum_{j=1}^{\infty} \left[\widetilde{(e_j, 0)}^2 f \right] (x)$$

is well-defined and independent of basis. In particular, if $f = F \circ \pi_P$, $x = (w, c) \in G$, ∂_h is as in Notation 3.12 for all $h \in \mathfrak{g}_{CM}$, and

$$\Delta_H f(x) := \sum_{j=1}^{\infty} \partial_{(e_j, 0)}^2 f(x) = (\Delta_{PH} F)(Pw, c),$$

then

$$(3.13) \quad Lf(x) = (\Delta_H f)(x) + f''(x) \left(\gamma(x) + \frac{1}{4} \chi(x) \right),$$

where γ and χ are as defined in Notation 3.16.

Proof. Recall from equation (3.8) that

$$\widetilde{(e_j, 0)} f(x) = f'(x) \left(e_j, \frac{1}{2} \omega(w, e_j) \right).$$

Applying $\widetilde{(e_j, 0)}$ to both sides of this equation then gives

$$\begin{aligned}\widetilde{(e_j, 0)}^2 f(x) &= f''(x) \left(\left(e_j, \frac{1}{2} \omega(w, e_j) \right) \otimes \left(e_j, \frac{1}{2} \omega(w, e_j) \right) \right) \\ &= f''(x) ((e_j, 0) \otimes (e_j, 0)) + f''(x) ((0, \omega(w, e_j)) \otimes (e_j, 0)) \\ &\quad + \frac{1}{4} f''(x) ((0, \omega(w, e_j)) \otimes (0, \omega(w, e_j))),\end{aligned}$$

wherein we have used that

$$\partial_{e_j} \omega(\cdot, e_j) = \omega(e_j, e_j) = 0$$

and the fact that $f''(x)$ is symmetric. Summing on j then shows that

$$\begin{aligned} \sum_{j=1}^{\infty} \left[\widetilde{(e_j, 0)}^2 f \right] (x) &= \sum_{j=1}^{\infty} f''(x) ((e_j, 0) \otimes (e_j, 0)) + f''(x) \left(\gamma(x) + \frac{1}{4} \chi(x) \right) \\ &= \sum_{j=1}^{\infty} \partial_{(e_j, 0)}^2 f(x) + f''(x) \left(\gamma(x) + \frac{1}{4} \chi(x) \right), \end{aligned}$$

which verifies equation (3.13) and thus shows that Lf is independent of the choice of orthonormal basis for H . \square

Similarly, we may prove the following proposition.

Proposition 3.18. *Let $\{e_j\}_{j=1}^{\infty}$ be an orthonormal basis for H . Then, for any smooth cylinder functions $f, g : G \rightarrow \mathbb{R}$,*

$$\Gamma(f, g)(x) := \sum_{j=1}^{\infty} \left(\widetilde{(e_j, 0)} f \right) (x) \left(\widetilde{(e_j, 0)} g \right) (x)$$

is well-defined and independent of basis. In particular, if $f = F \circ \pi_P$ and $g = G \circ \pi_Q$ for $P, Q \in \text{Proj}(W)$, $x = (w, c) \in G$, ∂_h is as in Notation 3.12 for all $h \in \mathfrak{g}_{CM}$, and

$$\nabla_H f(x) := \sum_{j=1}^{\infty} (\partial_{(e_j, 0)} f(x)) e_j = \nabla_{PH} F(Pw, c),$$

then

$$\begin{aligned} \Gamma(f, g)(x) &= \langle \nabla_H f(x), \nabla_H g(x) \rangle_H + \frac{1}{4} (f'(x) \otimes g'(x)) \chi(x) \\ &\quad + \frac{1}{2} (f'(x) \otimes g'(x) + g'(x) \otimes f'(x)) \gamma(x), \end{aligned}$$

where γ and χ are as defined in Notation 3.16.

Proof. Recall again from equation (3.8) that

$$\begin{aligned} \widetilde{(e_j, 0)} f(x) &= f'(x) \left(e_j, \frac{1}{2} \omega(w, e_j) \right) = f'(x) \left((e_j, 0) + \frac{1}{2} (0, \omega(w, e_j)) \right) \\ &= \partial_{(e_j, 0)} f(x) + \frac{1}{2} f'(x)(0, \omega(w, e_j)). \end{aligned}$$

Thus,

$$\begin{aligned} \Gamma(f, g)(x) &= \sum_{j=1}^{\infty} \left\{ \partial_{(e_j, 0)} f(x) \partial_{(e_j, 0)} g(x) + \frac{1}{2} f'(x)(0, e_j) g'(x)(0, \omega(w, e_j)) \right. \\ &\quad \left. + \frac{1}{2} f'(x)(0, \omega(w, e_j)) g'(x)(0, e_j) + \frac{1}{4} f'(x)(0, \omega(w, e_j)) g'(x)(0, \omega(w, e_j)) \right\}. \end{aligned}$$

Note for example that

$$\begin{aligned}
& \sum_{j=1}^{\infty} f'(x)(0, \omega(w, e_j)) g'(x)(e_j, 0) \\
&= \sum_{j=1}^{\infty} \langle f'(x) \otimes g'(x), (0, \omega(w, e_j)) \otimes (e_j, 0) \rangle \\
&= \left\langle f'(x) \otimes g'(x), \sum_{j=1}^{\infty} (0, \omega(w, e_j)) \otimes (e_j, 0) \right\rangle \\
&= \langle f'(x) \otimes g'(x), \gamma(x) \rangle.
\end{aligned}$$

Similarly,

$$\begin{aligned}
& \sum_{j=1}^{\infty} f'(x)(0, \omega(w, e_j)) g'(x)(0, \omega(w, e_j)) \\
&= \left\langle f'(x) \otimes g'(x), \sum_{j=1}^{\infty} (0, \omega(w, e_j)) \otimes (0, \omega(w, e_j)) \right\rangle \\
&= \langle f'(x) \otimes g'(x), \chi(x) \rangle.
\end{aligned}$$

□

Thus, along with L and Γ , we are able to consider the following differential forms which are well-defined for smooth cylinder functions f, g on G

$$\begin{aligned}
\Gamma_2(f, g) &:= \frac{1}{2} (L\Gamma(f, g) - \Gamma(f, Lg) - \Gamma(g, Lf)), \\
\Gamma^Z(f, g) &:= \sum_{\ell=1}^d \left(\widetilde{(0, f_\ell)} f \right) \left(\widetilde{(0, f_\ell)} g \right), \text{ and} \\
\Gamma_2^Z(f, g) &:= \frac{1}{2} (L\Gamma^Z(f, g) - \Gamma^Z(f, Lg) - \Gamma^Z(g, Lf)).
\end{aligned}$$

Of course, for the finite-dimensional groups G_P we may define the same forms for $f, g \in C^\infty(G_P)$ as was done for more general finite-dimensional groups in Section 2. These will be denoted by $L_P, \Gamma_P, \Gamma_{2,P}$, and $\Gamma_{2,P}^Z$. In particular, if $\{e_i\}_{i=1}^n$ is an orthonormal basis of PH , then

$$L_P f = \sum_{j=1}^n \widetilde{(e_j, 0)}^2 f \quad \text{and} \quad \Gamma_P(f, g) = \sum_{j=1}^n \left(\widetilde{(e_j, 0)} f \right) \left(\widetilde{(e_j, 0)} g \right).$$

3.5. Distances on G . We define the sub-Riemannian distance on G_{CM} analogously to how it was done in finite dimensions in Section 2. We recall its relevant properties, including the fact that the topology induced by this metric is equivalent to the topology induced by $\|\cdot\|_{\mathfrak{g}_{CM}}$.

Notation 3.19. (*Horizontal distance on G_{CM}*)

(1) For $x = (A, a) \in G_{CM}$, let

$$|x|_{\mathfrak{g}_{CM}}^2 := \|A\|_H^2 + \|a\|_{\mathbf{C}}^2.$$

The length of a C^1 -path $\sigma : [a, b] \rightarrow G_{CM}$ is defined as

$$\ell(\sigma) := \int_a^b |L_{\sigma^{-1}(s)*}\dot{\sigma}(s)|_{\mathfrak{g}_{CM}} ds.$$

(2) A C^1 -path $\sigma : [a, b] \rightarrow G_{CM}$ is horizontal if $L_{\sigma(t)^{-1}*}\dot{\sigma}(t) \in H \times \{0\}$ for a.e. t . Let $C_{CM}^{1,h}$ denote the set of horizontal paths $\sigma : [0, 1] \rightarrow G_{CM}$.

(3) The horizontal distance between $x, y \in G_{CM}$ is defined by

$$d(x, y) := \inf\{\ell(\sigma) : \sigma \in C_{CM}^{1,h} \text{ such that } \sigma(0) = x \text{ and } \sigma(1) = y\}.$$

The horizontal distance is defined analogously on G_P and will be denoted by d_P . In particular, for a sequence $\{P_n\}_{n=1}^\infty \subset \text{Proj}(W)$, we will let $d_n := d_{P_n}$.

Remark 3.20. Note that if $\sigma(t) = (A(t), a(t))$ is a horizontal path, then

$$L_{\sigma(t)^{-1}*}\dot{\sigma}(t) = \left(\dot{A}(t), \dot{a}(t) - \frac{1}{2}\omega(A(t), \dot{A}(t)) \right) \in H \times \{0\}$$

implies that σ must satisfy

$$a(t) = a(0) + \frac{1}{2} \int_0^t \omega(A(s), \dot{A}(s)) ds,$$

and the length of σ is given by

$$\ell(\sigma) = \int_0^1 |L_{\sigma^{-1}(s)*}\dot{\sigma}(s)|_{\mathfrak{g}_{CM}} ds = \int_0^1 \|\dot{A}(s)\|_H ds.$$

The following proposition is Propositions 2.17 and 2.18 of [15]. We refer the reader to that paper for the proof.

Proposition 3.21. *If $\{\omega(A, B) : A, B \in H\} = \mathbf{C}$, then there exist finite constants $K_1 = K_1(\omega)$ and $K_2 = K_2(d, \omega)$ such that*

$$(3.14) \quad K_1(\|A\|_H + \sqrt{\|a\|_{\mathbf{C}}}) \leq d(e, (A, a)) \leq K_2(\|A\|_H + \sqrt{\|a\|_{\mathbf{C}}}),$$

for all $(A, a) \in \mathfrak{g}_{CM}$. In particular, this is sufficient to imply that the topologies induced by d and $\|\cdot\|_{\mathfrak{g}_{CM}}$ are equivalent.

Remark 3.22. The equivalence of the homogeneous norm and horizontal distance topologies is a standard result in finite dimensions. However, the usual proof of this result relies on compactness arguments that must be avoided in infinite dimensions. Thus, the proof for Proposition 3.21 included in [15] necessarily relies on different methods particular to the structure of the present groups. Using these methods, we are currently unable to remove the dependence on $d = \dim(\mathbf{C})$ from the coefficient in the upper bound. The reader is referred to [15] for further details.

The following fact is not required for the sequel. However, it is natural to expect and thus we include it for completeness. Also, a similar limiting argument will be employed in the proof of quasi-invariance in Theorem 5.9.

Lemma 3.23. *For any $m \in \mathbb{N}$ and $x \in G_m$,*

$$d_n(e, x) \rightarrow d(e, x), \text{ as } n \rightarrow \infty.$$

Proof. First it is clear that, for any n and $x, y \in G_n$, $d_n(x, y) \geq d(x, y)$. In particular, if $x, y \in G_m$ for some m , then $x, y \in G_n$ for all $n \geq m$ and $d_n(x, y)$ is decreasing as $n \rightarrow \infty$. Now let $x = (w, c) \in G_m$, and consider an arbitrary horizontal path $\sigma : [0, 1] \rightarrow G_{CM}$ such that $\sigma(0) = e$ and $\sigma(1) = g$. Recall that, by Remark 3.20, σ must have the form

$$\sigma(t) = \left(A(t), \frac{1}{2} \int_0^t \omega(A(s), \dot{A}(s)) ds \right).$$

For $n \geq m$, consider the “projected” horizontal paths $\sigma_n : [0, 1] \rightarrow G_n$ given by

$$\sigma_n(t) = (A_n(t), a_n(t)) := \left(P_n A(t), \frac{1}{2} \int_0^t \omega(P_n A(s), P_n \dot{A}(s)) ds \right).$$

Note that $A_n(1) = P_n A(1) = P_n w = w$, and let

$$\varepsilon_n := c - a_n(1) = c - \frac{1}{2} \int_0^1 \omega(P_n A(s), P_n \dot{A}(s)) ds \in \mathbf{C}.$$

Then, for d_n the horizontal distance in G_n ,

$$d_n(e, x) = d_n(e, (w, c)) = d_n(e, (w, a_n(1) + \varepsilon_n)) = d_n(e, (w, a_n(1)) \cdot (0, \varepsilon_n)).$$

Now, for any left-invariant metric d , we have that

$$d(e, xy) \leq d(e, x) + d(x, xy) = d(e, x) + d(e, y).$$

Thus,

$$(3.15) \quad d_n(e, x) \leq d_n(e, (w, a_n(1))) + d_n(e, (0, \varepsilon_n)) \leq \ell(\sigma_n) + C\sqrt{\|\varepsilon_n\|_{\mathbf{C}}},$$

where the second inequality holds by (3.14) with constant $C = C(d, \omega)$ not depending on n . (Of course, the estimate (3.14) is stated for the horizontal distance d on G_{CM} and not d_n . However, it is clear from the proof in [15] that the same estimate holds for each d_P with common coefficients K_1 and K_2 for all sufficiently large $P \in \text{Proj}(W)$. See the proof of Proposition 2.17 of [15] for details.)

Now, for any $k \geq n$, it is clear that $\ell(\sigma_n) \leq \ell(\sigma_k)$, since

$$\ell(\sigma_n) = \int_0^1 \|P_n \dot{A}(s)\|_H ds = \int_0^1 \sqrt{\sum_{j=1}^n |\langle \dot{A}(s), e_j \rangle_H|^2} ds,$$

where $\{e_j\}_{j=1}^n$ is an orthonormal basis of $P_n H$. Thus, for all $k \geq n$,

$$(3.16) \quad d_n(e, x) \leq \ell(\sigma_k) + C\sqrt{\|\varepsilon_n\|_{\mathbf{C}}}.$$

Dominated convergence implies that

$$\lim_{k \rightarrow \infty} \ell(\sigma_k) = \lim_{k \rightarrow \infty} \int_0^1 \|P_k \dot{A}(s)\| ds = \int_0^1 \|\dot{A}(s)\| ds = \ell(\sigma),$$

and thus allowing $k \rightarrow \infty$ in (3.16) gives

$$d_n(e, x) \leq \ell(\sigma) + C\sqrt{\|\varepsilon_n\|_{\mathbf{C}}}.$$

Now taking the infimum over all horizontal paths in G_{CM} such that $\sigma(0) = e$ and $\sigma(1) = g$ implies that

$$d_n(e, x) \leq d(e, x) + C\sqrt{\|\varepsilon_n\|_{\mathbf{C}}}.$$

One may also show via dominated convergence that

$$\lim_{n \rightarrow \infty} \|\varepsilon_n\|_{\mathbf{C}} = \lim_{n \rightarrow \infty} \left\| \frac{1}{2} \int_0^1 \omega(A(s), \dot{A}(s)) - \omega(P_n A(s), P_n \dot{A}(s)) ds \right\|_{\mathbf{C}} = 0.$$

Thus, given an arbitrary $\varepsilon > 0$, for all sufficiently large n ,

$$d(e, x) \leq d_n(e, x) \leq d(e, x) + \varepsilon.$$

□

4. INFINITE-DIMENSIONAL COMPUTATIONS

Now given the structure of the infinite-dimensional Heisenberg-like groups and their finite-dimensional projections defined in the previous section, we wish to consider estimates like the ones discussed in Section 2. In the first subsection, we will show that the desired curvature-dimension estimate (2.2) and commutation formula (2.6) hold for the differential forms defined on G and G_P . In the second section, we then record the reverse inequalities and Harnack estimates that follow as a result.

4.1. Curvature-dimension bounds and commutation relations. In this section, $\{e_i\}_{i=1}^{\infty}$ and $\{f_\ell\}_{\ell=1}^d$ will denote orthonormal bases for H and \mathbf{C} respectively, where $d = \dim \mathbf{C}$. Let $\|\omega\|_2$ denote the Hilbert-Schmidt norm of $\omega : H \times H \rightarrow \mathbf{C}$ as defined in Notation 3.5. That is,

$$\|\omega\|_2^2 := \|\omega\|_{H^* \otimes H^* \otimes \mathbf{C}}^2 := \sum_{i,j=1}^{\infty} \|\omega(e_i, e_j)\|_{\mathbf{C}}^2 = \sum_{i,j=1}^{\infty} \sum_{\ell=1}^d \langle \omega(e_i, e_j), f_\ell \rangle_{\mathbf{C}}^2.$$

As ω is a continuous bilinear operator on W , Lemma 3.7 implies that ω is Hilbert-Schmidt and so $\|\omega\|_2 < \infty$.

For the rest of this section, we also fix $P \in \text{Proj}(W)$ and let $\{e_i\}_{i=1}^n$ denote an orthonormal basis of PH . Let $\|\omega\|_{2,P}$ denote the Hilbert-Schmidt norm of ω restricted to PH , that is,

$$\|\omega\|_{2,P}^2 := \sum_{i,j=1}^n \|\omega(e_i, e_j)\|_{\mathbf{C}}^2 = \sum_{i,j=1}^n \sum_{\ell=1}^d \langle \omega(e_i, e_j), f_\ell \rangle_{\mathbf{C}}^2.$$

We will also let

$$\rho_2 := \inf \left\{ \sum_{i,j=1}^{\infty} \left(\sum_{\ell=1}^d \langle \omega(e_i, e_j), f_\ell \rangle_{\mathbf{C}} x_\ell \right)^2 : \sum_{\ell=1}^d x_\ell^2 = 1 \right\} \text{ and}$$

$$\rho_{2,P} := \inf \left\{ \sum_{i,j=1}^n \left(\sum_{\ell=1}^d \langle \omega(e_i, e_j), f_\ell \rangle_{\mathbf{C}} x_\ell \right)^2 : \sum_{\ell=1}^d x_\ell^2 = 1 \right\}.$$

It is clear that

$$0 < \rho_2 \leq \|\omega\|_2^2 \quad \text{and} \quad 0 < \rho_{2,P} \leq \|\omega\|_{2,P}^2.$$

First we need the following computational lemmas.

Lemma 4.1. *For any smooth cylinder function f ,*

$$\rho_2 \Gamma^Z(f) \leq \sum_{i,j=1}^{\infty} \left((0, \widetilde{\omega(e_i, e_j)}) f \right)^2 \leq \|\omega\|_2^2 \Gamma^Z(f).$$

Similarly, for any $f \in C^\infty(G_P)$,

$$\rho_{2,P} \Gamma^Z(f) \leq \sum_{i,j=1}^n \left((0, \widetilde{\omega(e_i, e_j)}) f \right)^2 \leq \|\omega\|_{2,P}^2 \Gamma^Z(f).$$

Proof. We will prove only the first set of inequalities, as the second proof is obviously parallel. The upper bound follows from the Cauchy-Schwarz inequality, since

$$\begin{aligned} \sum_{i,j=1}^{\infty} \left((0, \widetilde{\omega(e_i, e_j)}) f \right)^2 &= \sum_{i,j=1}^{\infty} \left(\sum_{\ell=1}^d \langle \omega(e_i, e_j), f_\ell \rangle_{\mathbf{C}} \widetilde{(0, f_\ell)} f \right)^2 \\ &\leq \sum_{i,j=1}^{\infty} \left(\sum_{\ell=1}^d \langle \omega(e_i, e_j), f_\ell \rangle_{\mathbf{C}}^2 \right) \left(\sum_{\ell=1}^d \left(\widetilde{(0, f_\ell)} f \right)^2 \right) = \|\omega\|_2^2 \Gamma^Z(f). \end{aligned}$$

To see the lower bound, simply note that

$$\sum_{i,j=1}^{\infty} \left((0, \widetilde{\omega(e_i, e_j)}) f \right)^2 = \sum_{i,j=1}^{\infty} \left(\sum_{\ell=1}^d \langle \omega(e_i, e_j), f_\ell \rangle_{\mathbf{C}} \widetilde{(0, f_\ell)} f \right)^2 \geq \rho_2 \Gamma^Z(f).$$

□

Lemma 4.2. *For any smooth cylinder function f ,*

$$\Gamma_2^Z(f) = \sum_{j=1}^{\infty} \sum_{\ell=1}^d \left(\widetilde{(e_j, 0)} \widetilde{(0, f_\ell)} f \right)^2 \text{ and}$$

$$\sum_{j=1}^{\infty} \left(\sum_{i=1}^{\infty} \widetilde{(e_i, 0)} \widetilde{(0, \omega(e_i, e_j))} f \right)^2 \leq \|\omega\|_2^2 \Gamma_2^Z(f).$$

Similarly, for any $f \in C^\infty(G_P)$,

$$\Gamma_{2,P}^Z(f) = \sum_{j=1}^n \sum_{\ell=1}^d \left(\widetilde{(e_j, 0)} \widetilde{(0, f_\ell)} f \right)^2 \text{ and}$$

$$\sum_{j=1}^n \left(\sum_{i=1}^n \widetilde{(e_i, 0)} \widetilde{(0, \omega(e_i, e_j))} f \right)^2 \leq \|\omega\|_{2,P}^2 \Gamma_{2,P}^Z(f).$$

Proof. We find $\Gamma_2^Z(f)$ by a straightforward calculation:

$$\begin{aligned} \Gamma_2^Z(f) &= \frac{1}{2} \sum_{j=1}^{\infty} \widetilde{(e_j, 0)}^2 \sum_{\ell=1}^d \left(\widetilde{(0, f_\ell)} f \right)^2 - \sum_{\ell=1}^d \sum_{j=1}^{\infty} \left(\widetilde{(0, f_\ell)} f \right) \left(\widetilde{(0, f_\ell)} \widetilde{(e_j, 0)}^2 f \right) \\ &= \sum_{j=1}^{\infty} \sum_{\ell=1}^d \left\{ \left(\widetilde{(e_j, 0)} \widetilde{(0, f_\ell)} f \right)^2 + \left(\widetilde{(0, f_\ell)} f \right) \left(\widetilde{(e_j, 0)}^2 \widetilde{(0, f_\ell)} f \right) \right\} \\ &\quad - \sum_{\ell=1}^d \sum_{j=1}^{\infty} \left(\widetilde{(0, f_\ell)} f \right) \left(\widetilde{(0, f_\ell)} \widetilde{(e_j, 0)}^2 f \right) \\ &= \sum_{j=1}^{\infty} \sum_{\ell=1}^d \left(\widetilde{(e_j, 0)} \widetilde{(0, f_\ell)} f \right)^2, \end{aligned}$$

where we have used that $\widetilde{(0, f_\ell)}$ and $\widetilde{(e_j, 0)}$ commute by equation (3.9). The Cauchy-Schwarz inequality then implies that

$$\begin{aligned} &\sum_{j=1}^{\infty} \left(\sum_{i=1}^{\infty} \widetilde{(e_i, 0)} \widetilde{(0, \omega(e_i, e_j))} f \right)^2 \\ &= \sum_{j=1}^{\infty} \left(\sum_{i=1}^{\infty} \sum_{\ell=1}^d \langle \omega(e_i, e_j), f_\ell \rangle_{\mathbb{C}} \widetilde{(e_i, 0)} \widetilde{(0, f_\ell)} f \right)^2 \\ &\leq \sum_{j=1}^{\infty} \left(\sum_{i=1}^{\infty} \sum_{\ell=1}^d \langle \omega(e_i, e_j), f_\ell \rangle_{\mathbb{C}}^2 \right) \left(\sum_{i=1}^{\infty} \sum_{\ell=1}^d \left(\widetilde{(e_i, 0)} \widetilde{(0, f_\ell)} f \right)^2 \right) \\ &= \|\omega\|_2^2 \Gamma_2^Z(f). \end{aligned}$$

□

Lemmas 4.1 and 4.2 combine to give us the desired curvature-dimension bound.

Proposition 4.3. *For any $\nu > 0$ and smooth cylinder function f ,*

$$(4.1) \quad \Gamma_2(f) + \nu \Gamma_2^Z(f) \geq \rho_2 \Gamma^Z(f) - \frac{\|\omega\|_2^2}{\nu} \Gamma(f).$$

Similarly, for any $\nu > 0$ and $f \in C^\infty(G_P)$,

$$(4.2) \quad \Gamma_{2,P}(f) + \nu \Gamma_{2,P}^Z(f) \geq \rho_{2,P} \Gamma_P^Z(f) - \frac{\|\omega\|_{2,P}^2}{\nu} \Gamma_P(f).$$

Proof. Again, we prove only the first inequality. To do this, we first use the commutation relation in (3.4) and the estimates in Lemma 4.1 to estimate $\sum_{i,j=1}^\infty \left(\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} f \right)^2$. In what follows, we also use the antisymmetry of the form ω .

$$\begin{aligned} & \sum_{i,j=1}^\infty \left(\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} f \right)^2 \\ &= \sum_{i,j=1}^\infty \left(\frac{\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} + \widetilde{(e_j, 0)} \widetilde{(e_i, 0)}}{2} f + (0, \widetilde{\omega(e_i, e_j)}) f \right)^2 \\ &= \sum_{i,j=1}^\infty \left(\frac{\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} + \widetilde{(e_j, 0)} \widetilde{(e_i, 0)}}{2} f \right)^2 + \sum_{i,j=1}^\infty \left((0, \widetilde{\omega(e_i, e_j)}) f \right)^2 \\ &+ \sum_{i,j=1}^\infty \left(\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} f + \widetilde{(e_j, 0)} \widetilde{(e_i, 0)} f \right) \left((0, \widetilde{\omega(e_i, e_j)}) f \right) \\ &= \sum_{i,j=1}^\infty \left(\frac{\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} + \widetilde{(e_j, 0)} \widetilde{(e_i, 0)}}{2} f \right)^2 + \sum_{i,j=1}^\infty \left((0, \widetilde{\omega(e_i, e_j)}) f \right)^2 \\ &= \|\nabla_H^2 f\|^2 + \sum_{i,j=1}^\infty \left((0, \widetilde{\omega(e_i, e_j)}) f \right)^2. \end{aligned}$$

where

$$\nabla_H^2 f := \sum_{i,j=1}^\infty \frac{\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} f + \widetilde{(e_j, 0)} \widetilde{(e_i, 0)} f}{2}$$

denotes the symmetrized Hessian. Thus by Lemma 4.1

$$(4.3) \quad \|\nabla_H^2 f\|^2 + \rho_2 \Gamma^Z(f) \leq \sum_{i,j=1}^\infty \left(\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} f \right)^2 \leq \|\nabla_H^2 f\|^2 + \|\omega\|_2^2 \Gamma^Z(f).$$

Now we want to compute $\Gamma_2(f)$. The first term is simply

$$\frac{1}{2} L\Gamma(f) = \sum_{i,j=1}^\infty \left(\widetilde{(e_i, 0)} \widetilde{(e_i, 0)} \widetilde{(e_j, 0)} f \right) \widetilde{(e_j, 0)} f + \left(\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} f \right)^2.$$

The second term may be expanded by applying (3.4) twice as follows

$$\begin{aligned}\Gamma(f, Lf) &= \sum_{i,j=1}^{\infty} \left(\widetilde{(e_j, 0)} \widetilde{(e_i, 0)} \widetilde{(e_i, 0)} f \right) \widetilde{(e_j, 0)} f \\ &= \sum_{i,j=1}^{\infty} \left(\widetilde{(e_i, 0)} \widetilde{(e_i, 0)} \widetilde{(e_j, 0)} f \right) \widetilde{(e_j, 0)} f - 2 \left(\widetilde{(e_i, 0)} (0, \omega \widetilde{(e_i, e_j)}) f \right) \widetilde{(e_j, 0)} f.\end{aligned}$$

Thus, by the upper bound in (4.3) we have that

$$\begin{aligned}\Gamma_2(f) &= \frac{1}{2} (L\Gamma(f) - 2\Gamma(f, Lf)) \\ &= \sum_{i,j=1}^{\infty} \left(\widetilde{(e_i, 0)} \widetilde{(e_j, 0)} f \right)^2 + 2 \sum_{i,j=1}^{\infty} \left(\widetilde{(e_i, 0)} (0, \omega \widetilde{(e_i, e_j)}) f \right) \widetilde{(e_j, 0)} f \\ (4.4) \quad &\leq \|\nabla_H^2 f\|^2 + \|\omega\|_2^2 \Gamma^Z(f) + 2 \sum_{i,j=1}^{\infty} \left(\widetilde{(e_i, 0)} (0, \omega \widetilde{(e_i, e_j)}) f \right) \widetilde{(e_j, 0)} f.\end{aligned}$$

Similarly, the lower bound in (4.3) implies that

$$(4.5) \quad \Gamma_2(f) \geq \|\nabla_H^2 f\|^2 + \rho_2 \Gamma^Z(f) + 2 \sum_{i,j=1}^{\infty} \left(\widetilde{(e_i, 0)} (0, \omega \widetilde{(e_i, e_j)}) f \right) \widetilde{(e_j, 0)} f.$$

The Cauchy-Schwarz inequality now implies that for any $\nu > 0$

$$\begin{aligned}& 2 \left| \sum_{i,j=1}^{\infty} \left(\widetilde{(e_i, 0)} (0, \omega \widetilde{(e_i, e_j)}) f \right) \widetilde{(e_j, 0)} f \right| \\ & \leq \nu \sum_{j=1}^{\infty} \left(\sum_{i=1}^{\infty} \widetilde{(e_i, 0)} (0, \omega \widetilde{(e_i, e_j)}) f \right)^2 + \frac{1}{\nu} \sum_{j=1}^{\infty} \left(\widetilde{(e_j, 0)} f \right)^2 \\ (4.6) \quad & \leq \nu \|\omega\|_2^2 \Gamma_2^Z(f) + \frac{1}{\nu} \Gamma(f),\end{aligned}$$

where the last inequality follows from Lemma 4.2. Combining (4.4), (4.5), and (4.6) then gives

$$\begin{aligned}(4.7) \quad \Gamma_2(f) + \nu \|\omega\|_2^2 \Gamma_2^Z(f) &\geq \|\nabla_H^2 f\|^2 + \rho_2 \Gamma^Z(f) - \frac{1}{\nu} \Gamma(f) \\ &\geq \rho_2 \Gamma^Z(f) - \frac{1}{\nu} \Gamma(f).\end{aligned}$$

Finally, taking ν to be $\frac{\nu}{\|\omega\|_2^2}$ yields (4.1). \square

Remark 4.4. In the finite-dimensional case, one could use the Cauchy-Schwarz inequality in (4.7) to give a lower bound on $\|\nabla_H^2 f\|^2$ by $(Lf)^2$ with a coefficient depending on $\dim(H)$. Such an estimate standardly leads, for example, to Li-Yau type Harnack inequalities and bounds on logarithmic derivatives of the heat kernel. See for example [2, 4, 21].

We now prove that the desired commutation formula holds trivially on G and G_P .

Lemma 4.5. *For any smooth cylinder function f on G ,*

$$\Gamma(f, \Gamma^Z(f)) = \Gamma^Z(f, \Gamma(f)).$$

Similarly, for any $f \in C^\infty(G_P)$,

$$\Gamma_P(f, \Gamma_P^Z(f)) = \Gamma_P^Z(f, \Gamma_P(f)).$$

Proof. The proof is a straightforward computation.

$$\begin{aligned} \Gamma(f, \Gamma^Z(f)) &= \sum_{j=1}^{\infty} \left(\widetilde{(e_j, 0)} f \right) \left(\widetilde{(e_j, 0)} \Gamma^Z(f) \right) \\ &= \sum_{j=1}^{\infty} \sum_{\ell=1}^d \left(\widetilde{(e_j, 0)} f \right) \left(\widetilde{(e_j, 0)} \left(\widetilde{(0, f_\ell)} f \right)^2 \right) \\ &= 2 \sum_{j=1}^{\infty} \sum_{\ell=1}^d \left(\widetilde{(e_j, 0)} f \right) \left(\widetilde{(0, f_\ell)} f \right) \left(\widetilde{(e_j, 0)} \widetilde{(0, f_\ell)} f \right) \\ &= 2 \sum_{j=1}^{\infty} \sum_{\ell=1}^d \left(\widetilde{(0, f_\ell)} f \right) \left(\widetilde{(e_j, 0)} f \right) \left(\widetilde{(f_\ell, 0)} \left(\widetilde{(0, e_j)} f \right) \right) \\ &= \Gamma^Z(f, \Gamma(f)), \end{aligned}$$

where we have again used the commutativity of $\widetilde{(e_j, 0)}$ and $\widetilde{(0, f_\ell)}$ in the penultimate equality. The computation for the second equality is completely analogous. \square

4.2. Functional inequalities on G_P . Again note that for any $P \in \text{Proj}(W)$, $G_P = PH \times \mathbf{C}$ is a finite-dimensional step 2 stratified Lie group. If $\{e_j\}_{j=1}^n$ is an orthonormal basis of PH , then $\{\widetilde{(e_j, 0)}\}_{j=1}^n$ is a Hörmander set of vector fields on G_P . Thus we may apply the results of Section 2 to G_P . In particular, by Proposition 2.5, the curvature bound for G_P found in Proposition 4.3 implies the following reverse Poincaré inequality holds on all G_P . Here, we let $L_P = \sum_{j=1}^n \widetilde{(e_j, 0)}^2$ and $\{P_t^P\}_{t>0}$ denote the associated semi-group. Also, define the function classes \mathcal{C}_P and \mathcal{C}_P^+ analogously for G_P as was done in Notation 2.3.

Proposition 4.6 (Reverse Poincaré inequality). *For any $P \in \text{Proj}(W)$, $T > 0$, and $f \in \mathcal{C}_P$,*

$$\Gamma_P(P_T^P f) + \rho_{2,P} T \Gamma_P^Z(P_T^P f) \leq \frac{1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}}}{T} (P_T^P(f^2) - (P_T^P f)^2).$$

In particular, the following Reverse Poincaré inequality holds

$$\Gamma_P(P_T^P f) \leq \frac{1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}}}{T} (P_T^P(f^2) - (P_T^P f)^2).$$

Similarly, Theorem 2.7 implies that Proposition 4.3 coupled with the commutation relation of Lemma 4.5 give the following reverse log Sobolev inequality on all G_P .

Theorem 4.7 (Reverse log Sobolev inequality). *For any $P \in \text{Proj}(W)$, $T > 0$, and $f \in \mathcal{C}_P^+$,*

$$\Gamma_P(\ln P_T^P f) + \rho_2 T \Gamma_P^Z(\ln P_T^P f) \leq \frac{1 + \frac{2\|\omega\|_2^2}{\rho_2}}{T} \left(\frac{P_T^P(f \ln f)}{P_T^P f} - \ln P_T^P f \right)$$

In particular, the following reverse log Sobolev inequality holds

$$\Gamma_P(\ln P_T^P f) \leq \frac{1 + \frac{2\|\omega\|_2^2}{\rho_2}}{T} \left(\frac{P_T^P(f \ln f)}{P_T^P f} - \ln P_T^P f \right).$$

The combination of Proposition 2.10 with the reverse log Sobolev inequality found in Theorem 4.7 implies that the following Harnack type inequalities hold on each G_P .

Proposition 4.8 (Wang type Harnack inequality). *Let $P \in \text{Proj}(W)$. Then, for all $T > 0$, $x, y \in G_P$, $f \in L^\infty(G_P)$ with $f \geq 0$, and $p \in (1, \infty)$,*

$$(P_T^P f)^p(x) \leq P_T^P f^p(y) \exp \left(\left(1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}} \right) \frac{d_P^2(x, y)}{4(p-1)T} \right).$$

5. HEAT KERNEL MEASURE ON G AND A QUASI-INVARIANCE THEOREM

In this section, we show how the Wang type Harnack inequalities on G_P obtained in the previous section lead to the quasi-invariance of the subelliptic heat kernel measure on G . First, we must of course define the heat kernel measure on G , which we define as the end point distribution of a Brownian motion.

5.1. Brownian motion on G . We define a “subelliptic” Brownian motion $\{g_t\}_{t \geq 0}$ on G and collect various of its properties. The primary references for this section are Sections 4 of [11] and [13] and Section 2.5 of [15]. Any statements made here without proof are proved in these references.

Let $\{B_t\}_{t \geq 0}$ be a Brownian motion on W with variance determined by

$$\mathbb{E}[\langle B_s, h \rangle_H \langle B_t, k \rangle_H] = \langle h, k \rangle_H \min(s, t),$$

for all $s, t \geq 0$ and $h, k \in H_*$. The following is Proposition 4.1 of [11] and this result implicitly relies on the fact that Lemma 3.7 implies that the bilinear form ω is a Hilbert-Schmidt.

Proposition 5.1. *For $P \in \text{Proj}(W)$, let M_t^P denote the continuous L^2 -martingale on \mathbf{C} defined by*

$$M_t^P = \int_0^t \omega(PB_s, dPB_s).$$

In particular, if $\{P_n\}_{n=1}^\infty \subset \text{Proj}(W)$ is an increasing sequence of projections and $M_t^n := M_t^{P_n}$, then there exists an L^2 -martingale $\{M_t\}_{t \geq 0}$ in \mathbf{C} such that, for all $p \in [1, \infty)$ and $t > 0$,

$$\lim_{n \rightarrow \infty} \mathbb{E} \left[\sup_{\tau \leq t} \|M_\tau^n - M_\tau\|_{\mathbf{C}}^p \right] = 0,$$

and M_t is independent of the sequence of projections.

As M_t is independent of the defining sequence of projections, we will denote the limiting process by

$$M_t = \int_0^t \omega(B_s, dB_s).$$

Definition 5.2. The continuous G -valued process given by

$$g_t = \left(B_t, \frac{1}{2} M_t \right) = \left(B_t, \frac{1}{2} \int_0^t \omega(B_s, dB_s) \right).$$

is a *Brownian motion* on G . For $t > 0$, let $\nu_t = \text{Law}(g_t)$ denote the *heat kernel measure at time t* on G .

We include the following proposition (see [15, Proposition 2.30]) which states that, as the name suggests, the Cameron-Martin subgroup is a subspace of heat kernel measure 0.

Proposition 5.3. *For all $t > 0$, $\nu_t(G_{CM}) = 0$.*

Proposition 5.1 along with the fact that, for all $p \in [1, \infty)$ and $t > 0$,

$$\lim_{n \rightarrow \infty} \mathbb{E} \left[\sup_{\tau \leq t} \|B_\tau - P_n B_\tau\|_W^p \right] = 0$$

(see for example Proposition 4.6 of [11]) makes the following proposition clear.

Proposition 5.4. *For $P \in \text{Proj}(W)$, let g_t^P be the continuous process on G_P defined by*

$$g_t^P = \left(PB_t, \frac{1}{2} \int_0^t \omega(PB_s, dPB_s) \right).$$

Then g_t^P is a Brownian motion on G_P . In particular, let $\{P_n\}_{n=1}^\infty \subset \text{Proj}(W)$ be increasing projections and $g_t^n := g_t^{P_n}$. Then, for all $p \in [1, \infty)$ and $t > 0$,

$$\lim_{n \rightarrow \infty} \mathbb{E} \left[\sup_{\tau \leq t} \|g_\tau^n - g_\tau\|_{\mathfrak{g}}^p \right] = 0.$$

Notation 5.5. For all $P \in \text{Proj}(W)$ and $t > 0$, let $\nu_t^P := \text{Law}(g_t^P)$, and for all $n \in \mathbb{N}$ let $\nu_t^n := \text{Law}(g_t^n) = \text{Law}(g_t^{P^n})$.

Proposition 5.6. Let L be as defined in Proposition 3.17. Then we will call L the subelliptic Laplacian, and $\frac{1}{2}L$ is the generator for $\{g_t\}_{t \geq 0}$, so that, for any smooth cylinder function $f : G \rightarrow \mathbb{R}$,

$$f(g_t) - \frac{1}{2} \int_0^t Lf(g_s) ds$$

is a local martingale.

Corollary 5.7. Let $f = F \circ \pi_P$ be a cylinder function on G such that $F \in C^2(G_P)$ and there exist $K > 0$ and $p < \infty$ such that

$$|F(h, c)| + \|F'(h, c)\| + \|F''(h, c)\| \leq K(1 + \|h\|_{PH} + \|c\|_{\mathbf{C}})^p,$$

for all $(h, c) \in G_P$. Then

$$\mathbb{E}[f(g_t)] = f(e) + \frac{1}{2} \int_0^t \mathbb{E}[(Lf)(g_s)] ds.$$

That is,

$$\nu_t(f) := \int_0^t f d\nu_s = f(e) + \frac{1}{2} \int_0^t \nu_s(Lf) ds$$

is a weak solution to the heat equation

$$\partial_t \nu_t = \frac{1}{2} L \nu_t, \quad \text{with } \lim_{t \downarrow 0} \nu_t = \delta_e.$$

For all projections satisfying Hörmander's condition, the Brownian motions on G_P are subelliptic diffusions and thus their laws are absolutely continuous with respect to the finite-dimensional reference measure and their transition kernels are smooth. The following is Lemma 2.27 of [15].

Lemma 5.8. For all $P \in \text{Proj}(W)$ and $t > 0$, we have $\nu_t^P(dx) = p_t^P(x)dx$, where dx is the Riemannian volume measure (equal to Haar measure) and $p_t^P(x)$ is the heat kernel on G_P .

5.2. Quasi-invariance and Radon-Nikodym derivative estimates. For now, let us fix $P \in \text{Proj}(W)$, and recall that by Proposition 4.8, for all $T > 0$, $x, y \in G_P$, $f \geq 0$, and $p \in (1, \infty)$,

$$(P_T^P f)^p(x) \leq P_T^P f^p(y) \exp \left(\left(1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}} \right) \frac{d_P^2(x, y)}{4(p-1)T} \right).$$

Then Lemma 2.11 implies that this estimate is equivalent to

$$\left(\int_G \left[\frac{p_T^P(y, z)}{p_T^P(x, z)} \right]^{1/(p-1)} p_T^P(x, z) dz \right)^{p-1} \leq \exp \left(\left(1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}} \right) \frac{d_P^2(x, y)}{4(p-1)T} \right)$$

where p_T^P is the heat kernel on G_P . In particular, for $p \in (1, 2)$ and $q = 1/(p-1) \in (1, \infty)$, this implies that

$$(5.1) \quad \left(\int_G \left[\frac{p_T^P(y, z)}{p_T^P(x, z)} \right]^q p_T^P(x, z) dz \right)^{1/q} \leq \exp \left(\left(1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}} \right) \frac{qd_P^2(x, y)}{4T} \right)$$

Using the properties of heat kernels on finite-dimensional groups given in Lemma 2.1, we have that

$$\begin{aligned} \int_G \left[\frac{p_T^P(y, z)}{p_T^P(x, z)} \right]^q p_T^P(x, z) dz &= \int_G \left[\frac{p_T^P(yz^{-1})}{p_T^P(xz^{-1})} \right]^q p_T^P(xz^{-1}) dz \\ &= \int_G \left[\frac{p_T^P(zy^{-1})}{p_T^P(zx^{-1})} \right]^q p_T^P(zx^{-1}) dz. \end{aligned}$$

Then, for $x = e$, we may rewrite inequality (5.1) as

$$(5.2) \quad \left(\int_G \left[\frac{p_T^P(zy^{-1})}{p_T^P(z)} \right]^q p_T^P(z) dz \right)^{1/q} \leq \exp \left(\left(1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}} \right) \frac{qd_P^2(e, y)}{4T} \right).$$

Now, for $y \in G_P$, again we let $R_y : G_P \rightarrow G_P$ denote right translation. Then $\nu_T^P \circ R_y^{-1}$ is the push forward of ν_T under R_y . For fixed $T > 0$, let J_P^r denote the Radon-Nikodym derivative of $\nu_T^P \circ R_y^{-1}$ with respect to ν_T^P . Then (5.2) is equivalent to

$$(5.3) \quad \|J_P^r\|_{L^q(G_P, \nu_T^P)} \leq \exp \left(\left(1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}} \right) \frac{qd_P^2(e, y)}{4T} \right).$$

Alternatively, again using the properties of p_T^P described in Lemma 2.1 and the translation invariance of Haar measure, we could write

$$\begin{aligned} \int_G \left[\frac{p_T^P(y, z)}{p_T^P(x, z)} \right]^q p_T^P(x, z) dz &= \int_G \left[\frac{p_T^P(y^{-1}z)}{p_T^P(x^{-1}z)} \right]^q p_T^P(x^{-1}z) dz \\ &= \int_G \left[\frac{p_T^P(y^{-1}xz)}{p_T^P(z)} \right]^q p_T^P(z) dz. \end{aligned}$$

Then taking $x = e$ and combining this with the inequality (5.1) gives

$$\left(\int_G \left[\frac{p_T^P(y^{-1}z)}{p_T^P(z)} \right]^q p_T^P(z) dz \right)^{1/q} \leq \exp \left(\left(1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}} \right) \frac{qd_P^2(e, y)}{4T} \right).$$

which is equivalent to the left translation analogue

$$(5.4) \quad \|J_P^l\|_{L^q(G, \nu_T)} \leq \exp \left(\left(1 + \frac{2\|\omega\|_{2,P}^2}{\rho_{2,P}} \right) \frac{qd_P^2(e, y)}{4T} \right),$$

where $L_y : G_P \rightarrow G_P$ is left translation, $\nu_T^P \circ L_y^{-1}$ is the push forward of ν_T^P under L_y , and J_P^l denotes the Radon-Nikodym derivative of $\nu_T^P \circ L_y^{-1}$ with respect to ν_T .

Such estimates on the finite-dimensional projection groups G_P may be used to prove a quasi-invariance theorem on the infinite-dimensional group G . The following proof is analogous to the proofs of Theorem 7.2 and 7.3 in [12]. Similar methods were also used for the elliptic setting in loop groups in [9], in infinite-dimensional Heisenberg-like groups in [11], and in semi-infinite Lie groups in [24].

Theorem 5.9 (Quasi-invariance of ν_t). *For all $y \in G_{CM}$ and $T > 0$, ν_T is quasi-invariant under left and right translations by y . Moreover, for all $q \in (1, \infty)$,*

$$(5.5) \quad \left\| \frac{d(\nu_T \circ R_y^{-1})}{d\nu_T} \right\|_{L^q(G, \nu_T)} \leq \exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{qd^2(e, y)}{4T} \right)$$

and

$$\left\| \frac{d(\nu_T \circ L_y^{-1})}{d\nu_T} \right\|_{L^q(G, \nu_T)} \leq \exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{qd^2(e, y)}{4T} \right).$$

Proof. Fix $T > 0$ and $P_0 \in \text{Proj}(W)$. Let $y \in G_0$ and $\{P_n\}_{n=1}^\infty$ be an increasing sequence of projections such that $P_0H \subset P_nH$ for all n and $P_n|_H \uparrow I_H$. Let $J_n^r := J_{P_n}^r$ denote the Radon-Nikodym derivative of $\nu_T^n \circ R_y^{-1}$ with respect to ν_T^n . Then by the previous discussion and (5.3), we have

$$\|J_n^r\|_{L^q(G_n, \nu_T^n)} \leq \exp \left(\left(1 + \frac{2\|\omega\|_{2,n}^2}{\rho_{2,n}} \right) \frac{qd_n^2(e, y)}{4T} \right),$$

where we let $\|\omega\|_{2,n} := \|\omega\|_{2, P_n}$ and $\rho_{2,n} = \rho_{2, P_n}$.

Now let σ be an arbitrary horizontal path in G_{CM} such that $\sigma(0) = e$ and $\sigma(1) = y = (w, c) \in G_0 \subset G_n \subset G_{CM}$, and recall the projected horizontal paths $\sigma_n : [0, 1] \rightarrow G_n$ introduced in Lemma 3.23 given by

$$\sigma_n(t) = (A_n(t), a_n(t)) := \left(P_n A(t), \frac{1}{2} \int_0^t \omega(P_n A(s), P_n \dot{A}(s)) ds \right).$$

with $A_n(1) = P_n A(1) = P_n w = w$, and

$$\varepsilon_n := c - a_n(1) = c - \frac{1}{2} \int_0^1 \omega(P_n A(s), P_n \dot{A}(s)) ds \in \mathbf{C}.$$

It was proved in (3.15) that

$$d_n(e, y) \leq \ell(\sigma_n) + C\sqrt{\|\varepsilon_n\|_{\mathbf{C}}},$$

for a constant C independent of n , and thus

$$\|J_n^r\|_{L^q(G_n, \nu_T^n)} \leq \exp \left(\left(1 + \frac{2\|\omega\|_{2,n}^2}{\rho_{2,n}} \right) \frac{q(\ell(\sigma_n) + C\sqrt{\|\varepsilon_n\|_{\mathbf{C}}})^2}{4T} \right).$$

By Proposition 5.4, we have that for any bounded continuous f on G ,

$$(5.6) \quad \int_G f d\nu_t = \lim_{n \rightarrow \infty} \int_{G_n} f \circ i_n d\nu_t^n,$$

where $i_n : G_n \rightarrow G$ denotes the inclusion map. Note that

$$\begin{aligned} \int_{G_n} |(f \circ i_n)(xy)| d\nu_T^n(x) &= \int_{G_n} J_n^r(x) |(f \circ i_n)(x)| d\nu_T^n(x) \\ &\leq \|f \circ i_n\|_{L^{q'}(G_n, \nu_T^n)} \exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_{2,n}} \right) \frac{q(\ell(\sigma_n) + C\sqrt{\|\varepsilon_n\|_{\mathbf{C}}})^2}{4T} \right), \end{aligned}$$

where q' is the conjugate exponent to q . As in Lemma 3.23, we have via dominated convergence that $\ell(\sigma_n) \rightarrow \ell(\sigma)$ and $\|\varepsilon_n\|_{\mathbf{C}} \rightarrow 0$ as $n \rightarrow \infty$. Thus, allowing $n \rightarrow \infty$ in this last inequality yields

$$\int_G |f(xy)| d\nu_T(x) \leq \|f\|_{L^{q'}(G, \nu_T)} \exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{q\ell(\sigma)^2}{4T} \right)$$

by (5.6). Now optimizing this inequality over all horizontal paths σ in G_{CM} connecting e and y gives

$$(5.7) \quad \int_G |f(xy)| d\nu_T(x) \leq \|f\|_{L^{q'}(G, \nu_T)} \exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{qd(e, y)^2}{4T} \right).$$

Thus, we have proved that (5.7) holds for $f \in BC(G)$ and $y \in \cup_{P \in \text{Proj}(W)} G_P$. As this union is dense in G by Proposition 3.21, dominated convergence along with the continuity of $d(e, y)$ in y implies that (5.7) holds for all $y \in G_{CM}$.

Since the bounded continuous functions are dense in $L^{q'}(G, \nu_T)$ (see for example Theorem A.1 of [19]), the inequality in (5.7) implies that the linear functional $\varphi_y : BC(G) \rightarrow \mathbb{R}$ defined by

$$\varphi_y(f) = \int_G f(xy) d\nu_T(x)$$

has a unique extension to an element, still denoted by φ_y , of $L^{q'}(G, \nu_T)^*$ which satisfies the bound

$$|\varphi_y(f)| \leq \|f\|_{L^{q'}(G, \nu_T)} \exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{qd(e, y)^2}{4T} \right)$$

for all $f \in L^{q'}(G, \nu_T)$. Since $L^{q'}(G, \nu_T)^* \cong L^q(G, \nu_T)$, there then exists a function $J_y^r \in L^q(G, \nu_T)$ such that

$$(5.8) \quad \varphi_y(f) = \int_G f(x) J_y^r(x) d\nu_T(x),$$

for all $f \in L^{q'}(G, \nu_T)$, and

$$\|J_y^r\|_{L^q(G, \nu_T)} \leq \exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{qd(e, y)^2}{4T} \right).$$

Now restricting (5.8) to $f \in BC(G)$, we may rewrite this equation as

$$(5.9) \quad \int_G f(x) d\nu_T(xy^{-1}) = \int_G f(x) J_y^r(x) d\nu_T(x).$$

Then a monotone class argument (again use Theorem A.1 of [19]) shows that (5.9) is valid for all bounded measurable functions f on G . Thus, $d(\nu_T \circ R_y^{-1})/d\nu_T$ exists and is given by J_y^r , which is in L^q for all $q \in (1, \infty)$ and satisfies the bound (5.5).

A parallel argument employing the estimate in (5.4) gives the analogous result for $d(\nu_T \circ L_y^{-1})/d\nu_T$. Alternatively, one could use the right translation invariance just proved along with the facts that ν_T inherits invariance under the inversion map $y \mapsto y^{-1}$ from its finite-dimensional projections and that $d(e, y^{-1}) = d(e, y)$. \square

We may now observe that, by the proof of Lemma 2.11 or by the same limiting arguments as above, we have the following Wang type Harnack inequality on G .

Corollary 5.10 (Wang type Harnack inequality on G). *For all $T > 0$, $x, y \in G_{CM}$, $f \in L^\infty(G, \nu_T)$ with $f \geq 0$, and $p \in (1, \infty)$*

$$(P_T f)^p(x) \leq P_T f^p(y) \exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{d^2(x, y)}{4(p-1)T} \right).$$

Our final result is the following corollary which concerns two points. First, we are interested in the smoothing properties of the semi-group P_T in the absence of a reference measure. One way to approach this is to use a method similar to the proof of [28, Theorem 1.1]. A significant difference is that in [28] and some related work, the proof that the semi-group is strong Feller uses not only Harnack type inequalities, but also a version of Girsanov's theorem for the solution of the stochastic differential equation they consider. Note that we do not presently have path space quasi-invariance available in our degenerate case. The second point concerns the fact that the semi-group has smoothing properties, but only on the Cameron-Martin subgroup which has the heat kernel measure 0 as has been shown in [15, Proposition 2.30]. This is an infinite-dimensional phenomenon. In the flat abstract Wiener space setting such a phenomenon has been observed in [1]. In the context of holomorphic functions on complex abstract Wiener space it has been proved in [25, 26] and on complex infinite-dimensional groups in [7, 11, 14]. The importance of the strong Feller property for probabilistic potential theory in infinite dimensions has been discussed in [16]. Also, in [17] those authors explore the implications of a weaker version of the strong Feller property in a hypoelliptic setting.

Corollary 5.11 (Strong Feller skeleton). *For all $T > 0$ and $f \in L^\infty(G, \nu_T)$ such that $f \geq 0$, $(P_T f)(y) \rightarrow (P_T f)(x)$ as $d(x, y) \rightarrow 0$ for $x, y \in G_{CM}$.*

Proof. Fix $T > 0$ and $P_0 \in \text{Proj}(W)$. Let $y \in G_0$ and $\{P_n\}_{n=1}^\infty$ be an increasing sequence of projections such that $P_0 H \subset P_n H$ for all n and

$P_n|_H \uparrow I_H$. Then

$$\begin{aligned} |(P_T^n f)(y) - (P_T^n f)(x)| &\leq \int_{G_n} f(z) |p_T^n(x, z) - p_T^n(y, z)| dz \\ &\leq \|f\|_{L^\infty(G_n, \nu_T^n)} \int_{G_n} |p_T^n(x, z) - p_T^n(y, z)| dz \\ &\leq \|f\|_{L^\infty(G, \nu_T)} \int_{G_n} |p_T^n(x, z) - p_T^n(y, z)| dz. \end{aligned}$$

We also have that

$$\begin{aligned} \left(\int_{G_n} |p_T^n(x, z) - p_T^n(y, z)| dz \right)^2 &= \left(\int_{G_n} \left| \frac{p_T^n(x, z)}{p_T^n(y, z)} - 1 \right| p_T^n(y, z) dz \right)^2 \\ &\leq \int_{G_n} \left(\frac{p_T^n(x, z)}{p_T^n(y, z)} - 1 \right)^2 p_T^n(y, z) dz = \int_{G_n} \left(\frac{p_T^n(x, z)}{p_T^n(y, z)} \right)^2 p_T^n(y, z) dz - 1 \\ &\leq \exp \left(\left(1 + \frac{2\|\omega\|_{2,n}^2}{\rho_{2,n}} \right) \frac{d_n^2(x, y)}{2T} \right) - 1 \end{aligned}$$

by (5.1) with $q = 2$. Now by the same arguments as in the proof of Theorem 5.9, we may show that for any $x, y \in G_{CM}$

$$\begin{aligned} |(P_T f)(x) - (P_T f)(y)|^2 &\leq \|f\|_{L^\infty(G, \nu_T)}^2 \left(\exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{d^2(x, y)}{2T} \right) - 1 \right) \\ &\leq \|f\|_{L^\infty(G, \nu_T)}^2 \left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{d^2(x, y)}{2T} \exp \left(\left(1 + \frac{2\|\omega\|_2^2}{\rho_2} \right) \frac{d^2(x, y)}{2T} \right), \end{aligned}$$

since $e^x - 1 \leq xe^x$ for any $x > 0$. It is clear that the right-hand side tends to 0 when $d(x, y) \rightarrow 0$. \square

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